

Stellar Physics: Midterm Exam - Apr 25, 2014

Name:

Student ID:

Please work all of the following 5 problems. No notes, books, internet resources to be used. Some constants (in cgs units) that you may find useful are: $G = 6.67 \times 10^{-8}$, $h = 6.626 \times 10^{-27}$, $e = 4.8 \times 10^{-10}$, $m_u = 1.66 \times 10^{-24}$, $k = 1.38 \times 10^{-16}$, $\text{eV} = 1.6 \times 10^{-12}$, $\sigma = 5.67 \times 10^{-5}$, $M_\odot = 1.99 \times 10^{33}$, $R_\odot = 6.96 \times 10^{10}$, $L_\odot = 3.9 \times 10^{33}$.

1. (20 pts) In the first homework set, you were asked to solve the interior structure of a star of total mass M and radius R for a constant-density model (i.e. $\rho = \text{constant}$) and for a linear-density model (i.e. $\rho(r) \propto 1 - r/R$), subject to the zero outer boundary condition, i.e. $P(R) = T(R) = 0$. Assume that the star is composed of an ideal monatomic gas. Determine the range in radius for the convectively unstable region in these two stellar models.

Solution: Thermal convection occurs in an ideal monatomic gas when $(d \ln \rho / d \ln P)_{\text{ambient}} < (d \ln \rho / d \ln P)_{\text{blob}} = 1/\gamma = 3/5$, where the density change of the blob has been assumed adiabatic. In a constant-density model,

$$(d \ln \rho / d \ln P)_{\text{ambient}} = 0 < 3/5 \implies \text{the entire star is convective.} \quad (1)$$

In a linear-density model, let $x = r/R$ and then $\rho(x) = \rho_c(1 - x)$. Furthermore, $dP/dr = -\rho GM(r)/r^2$, $M_r = \int_0^r 4\pi r'^2 \rho_c(1 - r'/R) dr'$, and $P(x = 1) = 0$ give $P(x) = P_c(1 - 24x^2/5 + 28x^3/5 - 9x^4/5) = P_c[1 + 2x - (9/5)x^2](1 - x)^2$. Hence, thermal convection occurs when

$$\begin{aligned} (d \ln \rho / d \ln P)_{\text{ambient}} &= \frac{d \ln(1 - x)}{dx} \cdot \frac{dx}{d \ln[(1 + 2x - \frac{9}{5}x^2)(1 - x)^2]} \\ &= \frac{5 + 10x - 9x^2}{12x(4 - 3x)} < \frac{3}{5} \\ \implies &0.346 < x < 1.145 \\ \implies &\text{convection zone lies in the region } r \geq 0.346R. \end{aligned} \quad (2)$$

2. (20 pts)

2.1 Use the de Broglie's relation for the particle-wave duality, the Coulomb potential barrier, and the Maxwell-Boltzmann velocity distribution. Derive the thermonuclear reaction rate $\propto \int_0^\infty \exp(-E/kT - \bar{\eta}/\sqrt{E}) dE$, i.e. the area below the so-called "Gamow peak". In the equation, E is the relative kinetic energy between

two nuclei, T is the temperature, and $\bar{\eta} \approx Z_1 Z_2 e^2 \sqrt{(m_u/2)[A_1 A_2 / (A_1 + A_2)]} / h$ with Z being the atomic number and A being the mass number of each nucleus.

Solution: refer to the class notes.

- 2.2 Find the value of the kinetic energy $E = E_0$ that gives rise to the maximum value of the Gamow peak. Consider the pp chain reaction at $T = 10^7 K$. Calculate the energy ratio E_0/kT . What does the ratio physically mean?

Solution: The maximum value of the Gamow peak is obtained from

$$d(-E/kT - \bar{\eta}/\sqrt{E})/dE = 0 \implies E_0 = [(1/2)\bar{\eta}kT]^{2/3}. \text{ Hence,}$$

$$\frac{E_0}{kT} = \frac{[(1/2)\bar{\eta}kT]^{2/3}}{kT} \approx 0.45. \quad (3)$$

The ratio $\lesssim 1$ implies that the quantum tunneling does not enhance the pp-chain reaction rate for thermal protons in the high energy tail of the Maxwellian energy distribution against the Coulomb barrier. As a result, the reaction rate is less sensitive to temperature than it ought to be.

Comment: in reality, E_0/kT should be larger than 1 as we have learned in class. The error comes from the simplified expression for $\bar{\eta}$ we use, which underestimates $\bar{\eta}$. The more accurate expression is given by $\bar{\eta} = \pi(2m_{reduced})^{1/2}Z_1Z_2e^2/\hbar$, leading to $E_0/kT \approx (4\pi^2)^{2/3}(0.45) = 5.2 > 1$.

3. (20 pts) Consider a pre-main-sequence star of one solar mass on the Hayashi track. Its luminosity is $10L_\odot$ and its effective temperature is $T_{eff} = 4000 K$.

- 3.1 Ignore any thermonuclear heating and any mass accretion onto the star. Estimate the contraction timescale of the star due to radiative cooling from the stellar surface.

Solution:

The radius of the star is given by $R_* = \sqrt{L_*/4\pi\sigma T_{eff}^4} = \sqrt{10L_\odot/4\pi\sigma(4000K)^4} \approx 4.62 \times 10^{11} \text{ cm}$. The contraction time can then be estimated:

$$t_{KH} = \frac{GM_*^2}{R_*L_*} \approx \frac{GM_\odot^2}{(4.62 \times 10^{11} \text{ cm})(10L_\odot)} \approx 0.46 \text{ Myrs.} \quad (4)$$

- 3.2 Calculate the mean molecular weight and the number density n at the photosphere of the star. Take the Rosseland mean opacity $\kappa_R = 0.1 \text{ cm}^2/\text{g}$, $X = 0.74$, $Y = 0.25$. Assume that the gases are ideal, monatomic, and all neutrals. Ignore the contributions from elements heavier than He.

Solution: The mean molecular weight for an ideal, monatomic, electrically neutral gas consisting of H and He is given by

$$\mu = \left(\frac{X}{\mu_H} + \frac{Y}{\mu_{He}} \right)^{-1} = \left(\frac{0.74}{1} + \frac{0.25}{4} \right)^{-1} \approx 1.246. \quad (5)$$

The gas pressure at the photosphere is given by $P = (2/3)g/\kappa = (2/3)GM_*/R_*^2\kappa \approx 4.15 \times 10^3$ dyne/cm². The number density is obtained from the ideal gas law:

$$n = \frac{P}{kT_{eff}} \approx 7.52 \times 10^{15} \text{ cm}^{-3}. \quad (6)$$

- 3.3 Calculate the thermal ionization fraction $\equiv n_e/n$ contributed from elements K, Na, and H at the photosphere. Assume that the ions are far less than the neutrals for each elements. The ionization potentials of K and Na are $\chi_K = 4.34$ eV and $\chi_{Na} = 5.14$ eV. Take their abundances by number $n_K/n = 10^{-7}$ and $n_{Na}/n = 10^{-5.5}$.

Solution: For a weakly ionized gas, $n_j^+ \ll n_j$ and the Saha equation is reduced to $n_j^+/n_j \approx (m_e kT/2\pi\hbar^2)^{3/2} \exp(-\chi_j/kT)$. Summing over all species, we can obtain the thermal ionization fraction $n_e/n \approx [\sum_j (f_j/n) (m_e kT/2\pi\hbar^2)^{3/2} \exp(-\chi_j/kT)]^{1/2} \equiv [\sum_j (n_e/n)_j^2]^{1/2}$, where $f_j \equiv n_j/n$. Given that $f_K = 10^{-7}$, $f_{Na} = 10^{-5.5}$, and $f_H = (\rho_H/\mu_H)/(\rho/\mu) = (\mu/\mu_H)X \approx 0.922$, the thermal ionization fraction is given by

$$\begin{aligned} \frac{n_e}{n} &\approx \left[\left(\frac{n_e}{n}\right)_K^2 + \left(\frac{n_e}{n}\right)_{Na}^2 + \left(\frac{n_e}{n}\right)_H^2 \right]^{1/2} \\ &\approx (2.76 \times 10^{-8} + 8.59 \times 10^{-8} + 5.61 \times 10^{-13})^{1/2} \approx 3 \times 10^{-4}. \end{aligned} \quad (7)$$

Comment: the results show that the assumption for weak ionization of the alkaline elements K and Na is extremely poor. One should numerically solve the full version of the Saha equation for the thermal ionization fraction.

- 3.4 Assume that the interior can be modeled as an $n = 3/2$ polytrope. Calculate how much energy needs to be radiated away for the star to contract along the Hayashi track down to $L = L_\odot$ and $T_{eff} = 4500$ K.

Solution: The stellar radius at 4500 K is given by $R_* = \sqrt{L_*/4\pi\sigma T_{eff}^4} = \sqrt{L_\odot/4\pi\sigma(4500K)^4} \approx 1.16 \times 10^{11}$ cm. The energy radiated away is given by the difference of the total energy:

$$W_f - W_i = -\alpha_{n=3/2} \frac{GM_*^2}{2r_f} + \alpha_{n=3/2} \frac{GM_*^2}{2r_i} \approx -8.5 \times 10^{47} (\alpha_{n=3/2}) \text{ erg}. \quad (8)$$

where $\alpha_{n=3/2}$ needs to be solved numerically using the Lane-Emden Equation but it should be on the order of unity.

4. (20 pts) The concept of the Schönberg-Chandrasekhar limit for a collapse of an isothermal inert He core in a star can apply to the condition for star formation in an interstellar cloud.

- 4.1 Derive the virial theorem for a self-gravitating interstellar cloud of total mass M and radius R confined by an external pressure of P_0 . Assuming that the cloud has a constant temperature T .

Solution: Refer to the slide for the “Schönberg-Chandrasekhar limit. Starting from the hydrostatic equilibrium, we have

$$\frac{dP}{dm} = -\frac{Gm}{4\pi r^4} \implies \int_{\text{cloud}} 4\pi r^3 dP = - \int_{\text{cloud}} \frac{Gm}{r} dm = E_g, \quad (9)$$

where E_g is the gravitational energy of the cloud. Integration by parts on the LHS gives

$$4\pi R^3 P_0 - 3 \int_{\text{cloud}} 4\pi P r^2 dr = 4\pi R^3 P_0 - 2 \int_{\text{cloud}} (3/2)nkT4\pi r^2 dr = 4\pi R^3 P_0 - 2E_i, \quad (10)$$

where E_i is the internal energy ($E_i = (3/2)(RT/\mu)M$ for a constant T). Hence, the virial theorem for the cloud confined by an external pressure is given by

$$4\pi R^3 P_0 = 2E_i + E_g. \quad (11)$$

- 4.2 Holding M and T fixed, show that a maximum value, P_{max} , exists for the external pressure for a certain value of the cloud radius, R_{max} . Derive the expressions for R_{max} and P_{max} in terms of the cloud mass M , temperature T , and fundamental constants such as G . In what condition does the cloud gravitationally collapse?

Solution: Following up the above result (or referring to the class notes), we have

$$P_0 = \left(\frac{3RT}{4\pi\mu} \right) \frac{M}{R^3} - \left(\frac{\alpha G}{4\pi} \right) \frac{M^2}{R^4}, \quad (12)$$

where $\alpha \equiv -E_g/(GM^2/R)$ and is an order of unit constant. At R_{max} , $dP_0/dR = 0$, which yields

$$\left(\frac{-9RT}{4\pi\mu} \frac{M}{R^4} + \frac{4\alpha G}{4\pi} \frac{M^2}{R^5} \right)_{R=R_{max}} = 0 \implies R_{max} = \left(\frac{4\alpha G\mu}{9RT} \right) M. \quad (13)$$

Thus,

$$P_{max} = \left(\frac{3^7 R^4}{4^5 \pi \alpha^3 G^3 \mu^4} \right) \frac{T^4}{M^2}. \quad (14)$$

The cloud collapses when $P_0 > P_{max}$.

5. (20 pts) Sketch and explain the evolution of a $1M_{\odot}$ star from a pre-main-sequence star to a cool white dwarf on the H-R diagram. You may draw the pre-main sequence evolution and the main sequence and post-main sequence evolution on separate H-R diagrams to avoid crowding.

Solution: refer to the class notes. There are a number of evolutionary states, but it should be noted that the gravitational core collapse are as important as thermonuclear reactions throughout the stellar evolution: stellar birth line, Hayashi track, Henyey track, ZAMS (pp-chain H burning begins), MS (luminosity gradually increases), Schönberg-Chandrasekhar limit, subgiant (He core contracts, H shell burning), RGB (He core keeps contracting and becomes degenerate), He core flash, horizontal branch (He core burning, H shell burning), AGB (H and He shell burning, C/O contracts and becomes degenerate), AGB winds (He shell flashes), PNe, WD