#### **CosPA 2007 (2007/11/14)**

# A Moment Method for Measuring the Higher Order Weak Lensing Effects

**References:** 

Okura, Umetsu, Futamase 2007, ApJ, 660, 995

Okura, Umetsu, Futamase 2007 (astro-ph/0710.2262)

(Umetsu & Broadhurst 2007)

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# **Gravitational Lensing**

The images of background sources carry the imprint of the gravitational potential of the intervening matter.

Gravitational Lensing is essentially a coordinate mapping from a foreground (source) to the background (image) plane:



#### **Light-Ray Bending and Image Deformation**



For an infinitesimal light source: 
$$d^2 \vec{\beta} = \mathbf{A} d^2 \vec{\theta}$$
 with  $\mathbf{A} \equiv \mathbf{1} - \nabla \nabla \phi \neq \mathbf{1}$  A: Jacobian matrix of the lens equation the lens equation the lens equation Spinoreal:  
Lens mapping Eq. in general:  $\delta \beta_i = D_{ij}^{(1)} \delta \theta_j + \frac{1}{2} D_{ijk}^{(2)} \delta \theta_j \delta \theta_k + \dots$ 

# Ordinal (1<sup>st</sup> order) Weak Lensing

**2<sup>nd</sup> derivatives of** φ: **Area distortion** (magnification) and **Quadrupole shape deformation** 

Jacobian matrix of lens eq.

$$\mathbf{A}_{ij}(\mathbf{\theta}) = \delta_{ij} - \nabla_i \nabla_j \phi$$

$$\mathbf{A}(\mathbf{\theta}) = (1 - \kappa) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \begin{pmatrix} \gamma_1 & \gamma_2 \\ \gamma_2 & -\gamma_1 \end{pmatrix}$$

 $\partial := \partial_1 + i \partial_2$ 

complex differential operator

2

#### 2D Poission Eq.

$$\kappa = \frac{1}{2} \partial \partial^* \phi = \frac{\Sigma_m}{\Sigma_{\text{crit}}}$$
$$\gamma = \frac{1}{2} \partial \partial \phi \equiv \gamma_1 + i\gamma$$

 $(\partial \partial \phi, \partial \partial^* \phi)$ dof=2+1\_\_\_\_

#### Lensing convergence ::

Dimensionless surface mass density of the lens (spin-0)

**Complex Gravitational Shear**:: spin-2 quadrupole image distoriton

## 2<sup>nd</sup> Order Weak Lensing: Flexion

Flexion is 3<sup>rd</sup> order derivatives of the potential:

$$\mathbf{D}_{ijk}^{(2)} = \mathbf{A}_{ij,k} = -\nabla_i \nabla_j \nabla_k \phi \left[ \begin{array}{c} \frac{\mathrm{dof}=2+2}{(\partial \partial \phi, \partial \partial \partial^* \phi)} \end{array} \right]$$

#### **Complex First flexion**

$$F = \frac{1}{2} \partial^* \partial \partial \phi = \partial \kappa$$
 Spin

**Complex Second flexion** 

$$G = \frac{1}{2} \partial \partial \partial \phi = \partial \gamma$$

Spin-3

(Goldberg & Natarajan 2002, Bacon et al. 2005)

#### Effects of Convergence, Shear, Flexion



#### **Flexion and Shear in Shapelets**



#### Where is Flexion useful?

An intermediate regime between WEAK and STRONG lensing can be well described by shearing and flexing effects: Arclets = lensed images with slight curvatures



#### **Shear vs. Flexion**

**Resolution limit in ordinal (=Spin-2) weak lensing with ground based telescopes (Subaru, CFHT, etc.):** 

FWHM = 0.'9 
$$\left(\frac{S/N}{5}\right) \left(\frac{\kappa}{0.2}\right)^{-1} \left(\frac{\sigma_{\gamma}}{0.4}\right)^{-1} \left(\frac{n_g}{30 \text{ arcmin}^{-2}}\right)^{1/2}$$

Ordinal WL is sensitive to structures of  $\underline{1'-10'}$ , which is dominated by clusters of galaxies

Flexion measures the gradient of shear; so is relatively sensitive to small-scale structures (e.g., galaxies, groups of galaxies)

$$\frac{S(F)}{S(\gamma)} \sim \frac{L\phi/r^3}{\phi/r^2} \sim \frac{L}{r}$$

L: image size

r: distance from the lens

Even though the higher-order effect is small, at small scales (r), for large images (L), Flexion signal might dominate over Shear signal

#### How to Measure Flexion? Higher Order Lensing Image Characteristics (HOLICs)

Observable brightness-weighted shape moments of galaxy images:

$$Q_{ijk} = \langle x_i x_j x_k \rangle$$
$$Q_{ijkl} = \langle x_i x_j x_k x_l \rangle$$

etsu.

Okura, Umetsu, Futamase (2007a) find natural combinations of higher-order moments to define HOLICs, and have shown explicit relations between HOLICs and Flexion:

Spin-1 HOLICs  

$$\begin{aligned} \varsigma \coloneqq \frac{Q_{111} + Q_{122} + i(Q_{112} + Q_{222})}{Q_{1111} + Q_{2222} + 2Q_{1122}} \sim \frac{e^{i\varphi}}{L} \\
\end{aligned}$$
Spin-3 HOLICs  

$$\begin{aligned} \delta \coloneqq \frac{Q_{111} - 3Q_{122} + i(3Q_{112} - Q_{222})}{Q_{1111} + Q_{2222} + 2Q_{1122}} \sim \frac{e^{3i\varphi}}{L} \\
\zeta^{(s)} \approx \zeta - \frac{9}{4}F, \quad \delta^{(s)} \approx \delta - \frac{3}{4}G \end{aligned}$$

$$\begin{aligned} e \approx e^{(s)} + 2\gamma \\
\end{aligned}$$
Okura, Um  
Futamase

#### Flexion to Surface Mass Density $\kappa = \Sigma / \Sigma_{crit}$

Observable Flexion measures the 1<sup>st</sup> derivative of shear

Non-parametric inversion from flexion to mass density

0.00

0.01

$$F = \frac{1}{2} \partial \partial \partial^* \phi = \partial \kappa$$
$$G = \frac{1}{2} \partial \partial \partial \phi = \partial \gamma$$



$$\kappa = \Delta^{-1}(\partial^* F)$$
$$\kappa = \Delta^{-1}(\partial^* \partial^* \partial^* G)$$

Bacon et al. 2005





0.02

0.03

# **Application to a Real Cluster?**



Negative peak at cD location!!

Substructure detected by 1<sup>st</sup> Flexion

No significant detection using a large # of galaxies: ng ~ 70 galaxies/arcmin^2

#### A1689 HST/ACS data



Core regon of A1689 (z=0.18)

#### **Problems and Motivations**

Shapelets:: Leonard et al. (2007) found HOLICs technique is less sensitive than Shapelets to contamination by bright wings of cluster (=lens) galaxies and other neighboring sources. So HOLICs technique seems to be suitable for cluster lensing.

HOLICS+ :: Goldberg & Leonard (2007) extended our HOLICs method to include isotropic PSF corrections but ignoring (1) spin-1 and spin-3 PSF anisotropy corrections, (2) additional correction terms for Gaussian-weighted moment calculations, etc..

Fully generalize "KSB formalism" for ordinal spin-2 weak lensing, by including higher-order Gaussian-weighted moment calculations (HOLICs) and relevant corrections for spin-1 and spin-3 PSF anisotropies Okura, Umetsu, Futamase 2007b

#### **HOLICs Moment Method**

KSB formalism (Kaiser, Squares, Broadhurst 1995): Adds in Lensing and PSF induced response to intrinsic ellipticity that is randomly oriented

$$e_{\alpha} = e_{\alpha}^{(s)} + P_{\alpha\beta}^{\gamma} \gamma_{\beta} + P_{\alpha\beta}^{sm} q_{\beta}$$

Higher-order generalization of KSB moment method by Okura, Umetsu, Futamase 2007b

Unbiased estimator for Flexion

$$F \approx C_F^{-1} \left\langle \varsigma - C^q \varsigma_q \right\rangle \quad G \approx D$$

$$G \approx D_F^{-1} \left\langle \delta - D^q \delta_q \right\rangle$$

## **Subaru Observations of A1689**



#### Spin-1 PSF Anisotropy Correction: Application to Subaru A1689 data



#### Mass Map of A1689 from Spin-1 Flexion

Mass reconstruction in the 4'x4' core region of A1689 (z=0.18)E-mode (lensing)B-mode (noise)



# Mass and Light in A1689 (Subaru)



# Results

- First successful cluster mass reconstruction only by Flexion incorporating full higherorder PSF anisotropy corrections
- Significant detections of cD/group-scale substructures at 5-6 sigma levels
- Main density peak of κ~2.5<sup>+0.45</sup>-0.45 is consistent with the ACS/Subaru mass model (Broadhurst, Takada, Umetsu et al. 05).

#### **Conclusions and Future Work**

Weak lensing fiexion has been proposed in the past years, and is just becoming a usable practical tool for observational cosmology - not just a theoretical tool.

A comparison with high-resolution HST/ACS data shows that indeed higher-order PSF anisotropy corrections are necessary for a high-precision measurement of the higher-order WL signal.

Although the current HOLICs-based flexion analysis pipeline works well to locate mass substructures, an accurate mass measurement will require systematic tests and feasible "calibrations" as people have done for quadrupole WL (e.g., the STEP project).

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