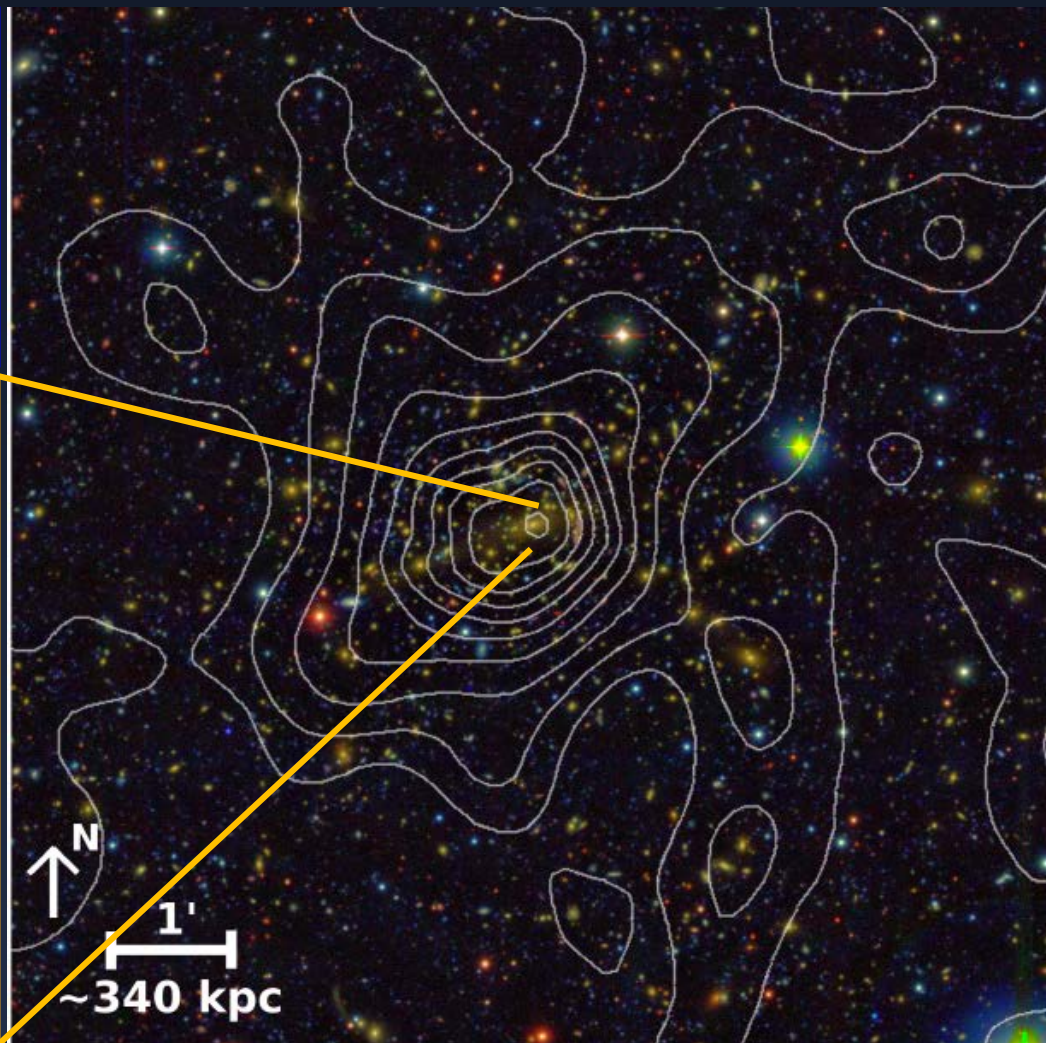


# The Full Strength of Cluster Gravitational Lensing: *Matter Distribution in and around Cosmic Giants from the CLASH Survey*

Cluster **L**ensing **A**nd **S**upernova survey with **H**ubble



Keiichi Umetsu (ASIAA, Taiwan) with the CLASH team

# Contents

- 1. Introduction: Galaxy Clusters**
- 2. Cluster Weak Lensing: Shear & Magnification**
- 3. Cluster Lensing Results from CLASH**
- 4. Intracluster Dark Matter Equation of State**
- 5. Summary**

# 1. Introduction

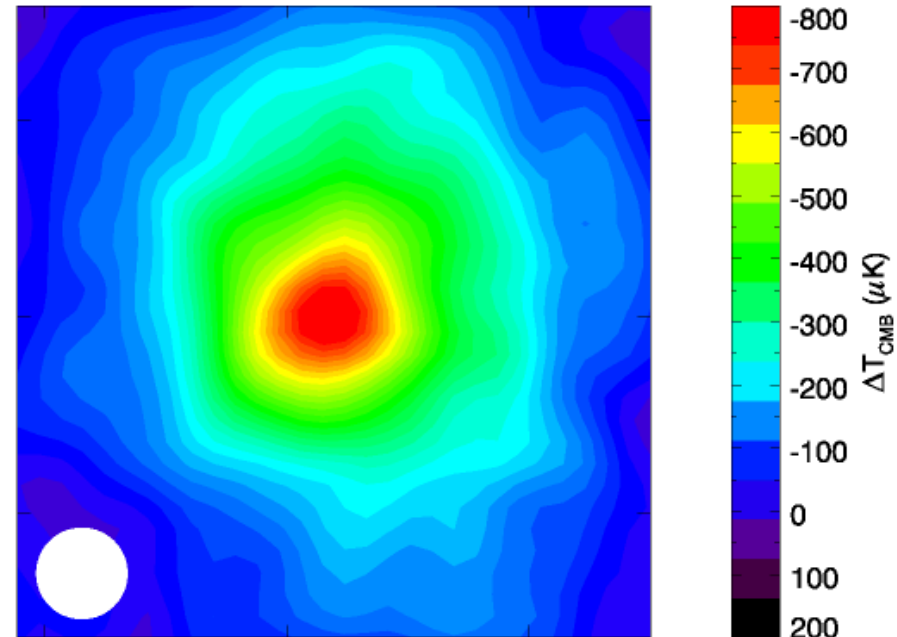
## **Galaxy Clusters as Cosmological Probe**



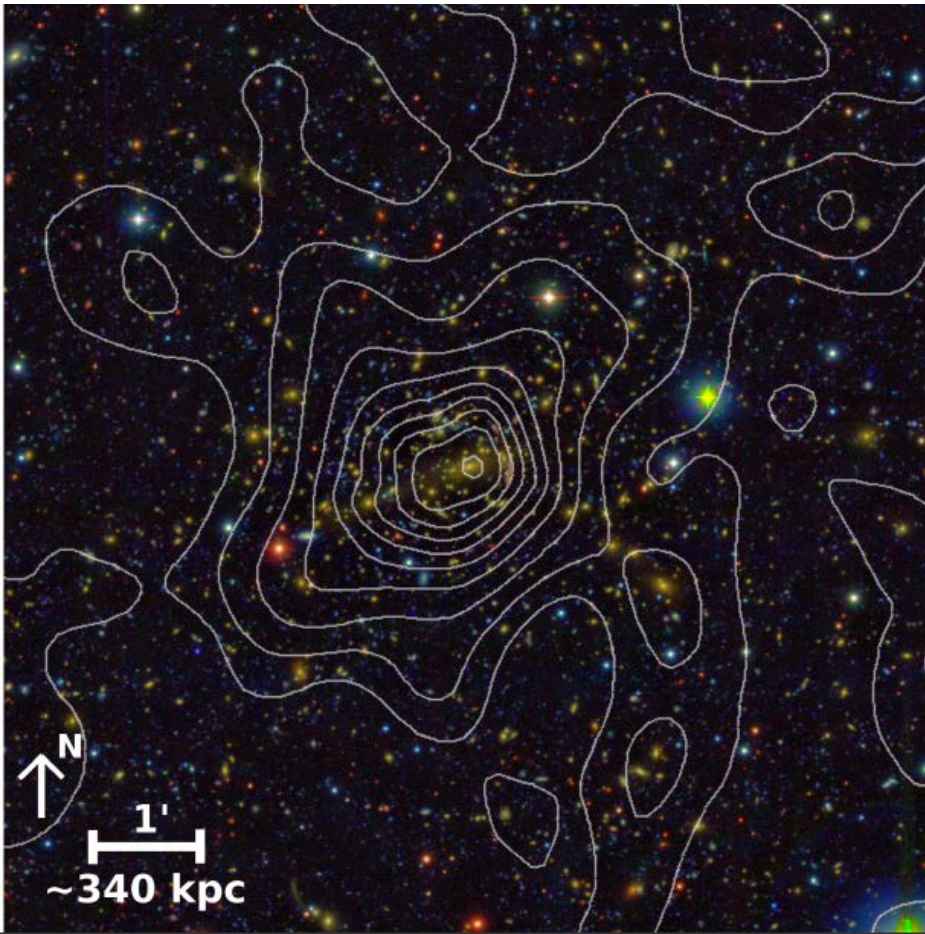
# Clusters of Galaxies

**Clusters** = composed of 100-1000 galaxies and filled with hot, diffuse intracluster plasmas of  $k_B T = 3-10 \text{ keV}$ , corresponding to  $M_{\text{grav}} = 10^{14-15} M_{\text{sun}}$

## Sunyaev-Zel'dovich Effect (SZE)



12<sup>h</sup> 6<sup>m</sup> 30<sup>s</sup>      6<sup>m</sup> 15<sup>s</sup>      6<sup>m</sup> 0<sup>s</sup>  
RA (J2000)

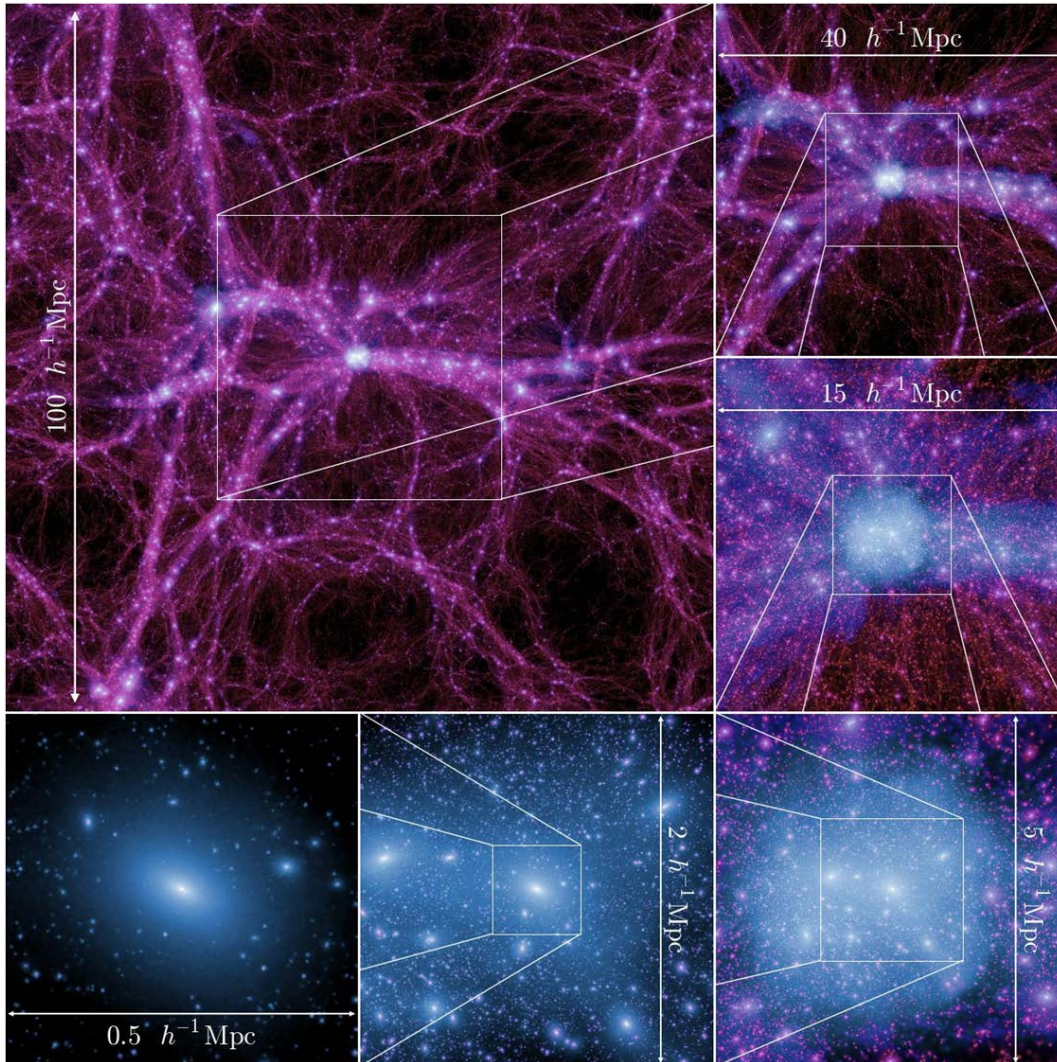


MACS1206 cluster at  $z=0.44$   
(Umetsu et al. 2012, *ApJ*, 755, 56)



# Clusters: the largest/youngest class of DM halos

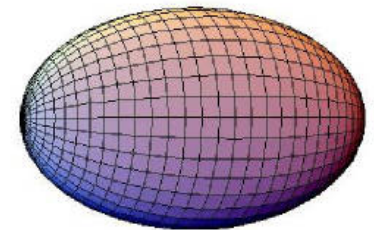
**Halos = gravitationally-bound objects**  $\frac{1}{2}\ddot{I} = 2K + U - E^{(S)} \sim 0$



Clusters formed at the intersection of filaments and sheets

Typical formation epoch:  
 $z_f = 0.5 - 0.7$

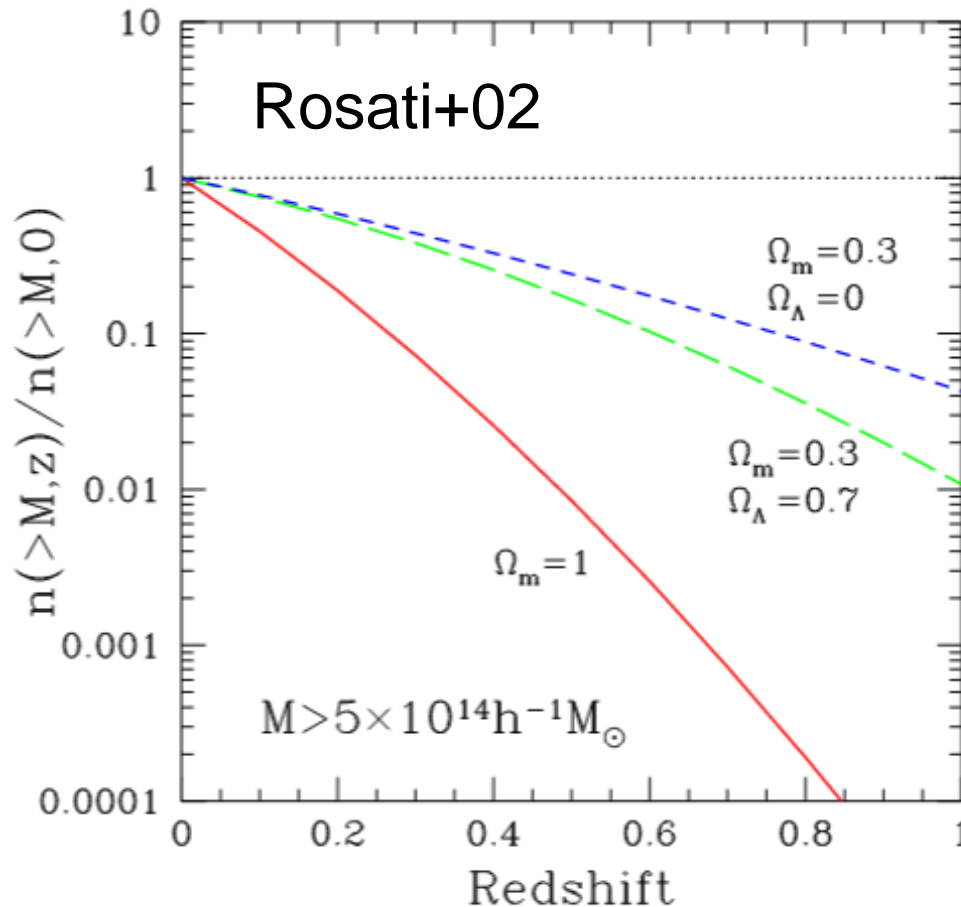
Young halos are prolate  
(collisionless nature)



Boylan-Kolchin+09

# Clusters as Cosmological Probe

Cluster counts  $\frac{dN(> M_{\text{lim}}, z)}{d\Omega dz} = \int_{M_{\text{lim}}}^{\infty} dM \frac{dV(z)}{d\Omega dz} \frac{d^2 n}{dV dM}(M, z)$



## Comoving volume element

$$\frac{d^2 V}{dz d\Omega} = \frac{cr^2[\chi(z)]}{H(z)}, \quad \chi(z) = \int_0^z \frac{dz'}{H(z')}$$

## Halo mass function

$$\frac{d^2 n}{dV dM}(M, z) \propto \exp\left[-\frac{\nu^2}{2}\right]$$

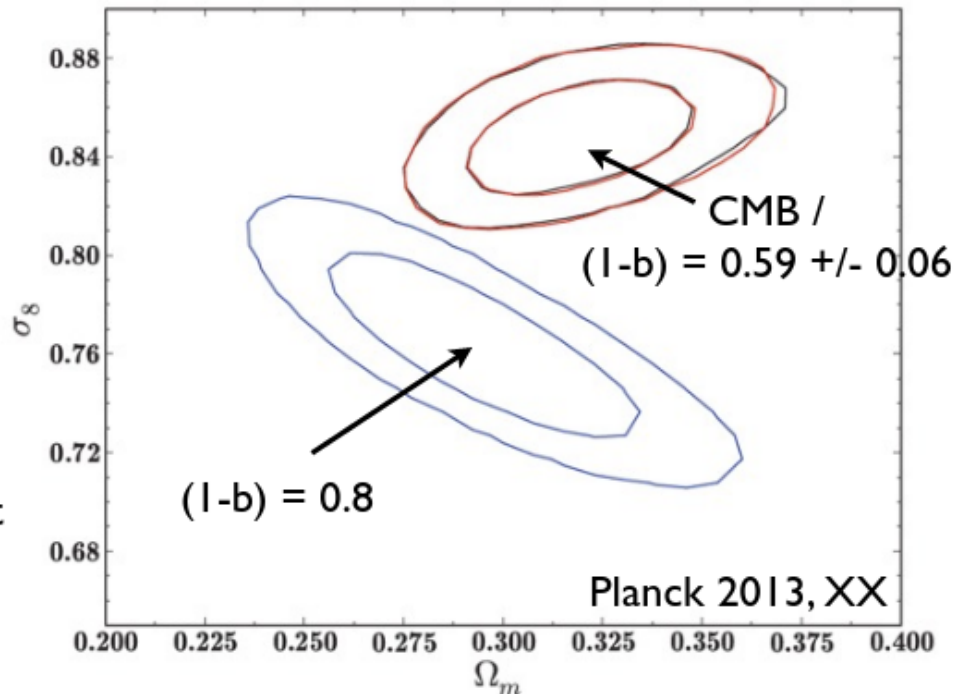
$$\nu \equiv \frac{\delta_c(z)}{\sigma(M, z)} \approx \frac{1.69}{D_+(z)\sigma(M)} \sim 3 \text{ for clusters}$$

Cluster counts are exponentially sensitive to *cosmology* and **cluster mass calibration!**

# Planck13 CMB vs. Cluster Cosmology

$b=0.2?? - 0.4??$

- Planck:  $3\sigma$  tension between SZ cluster counts and CMB cosmology
- assumes  $M_{\text{Planck}} / M_{\text{true}} = (1-b) = 0.8$
- calibrated with XMM hydrostatic masses (Arnaud et al. 2010) + simulations



suggested explanations:

- **mass bias underestimated** (and no accounting for uncertainties)
- $2.9\sigma$  detection of neutrino masses:  $\Sigma m_\nu = (0.58 \pm 0.20) \text{ eV}$   
(Planck+WMAPpol+ACT+BAO:  $\Sigma m_\nu < 0.23 \text{ eV}$ , 95% CL)

*Slide taken from Anja von der Linden's presentation*

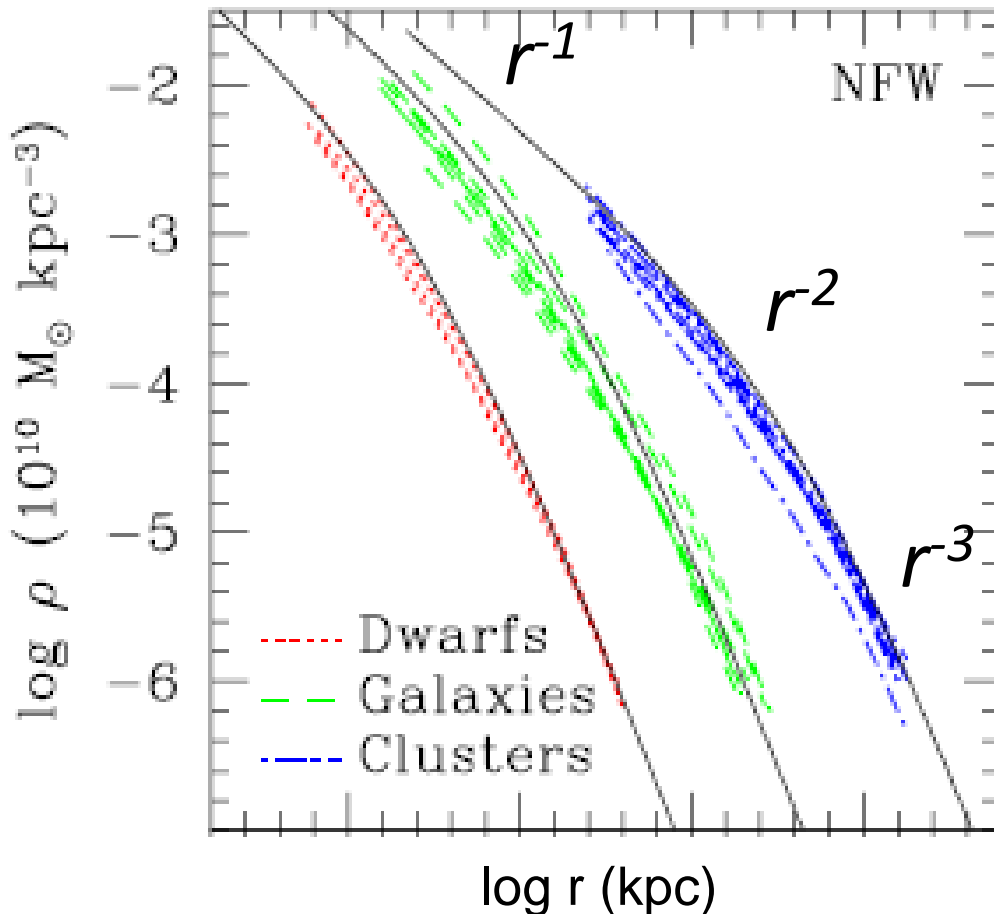


# Key Predictions of nonlinear structure formation models

(1) Quasi self-similar DM-halo density profiles

# Quasi Self-similar Halo Density Profile for collisionless CDM

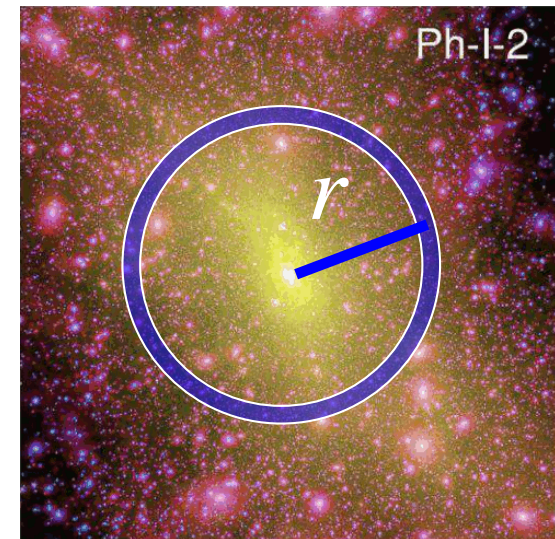
Spherically-averaged DM density profiles  $\rho(r)$  from numerical simulations



Cuspy, outwardly-steepening density profiles

**Empirical fitting formula by Navarro-Frenk-White (NFW)**

$$\begin{aligned}\rho(r) &= \rho_s f(r/r_s) \\ &= \frac{\rho_s}{(r/r_s)(1+r/r_s)^2}\end{aligned}$$



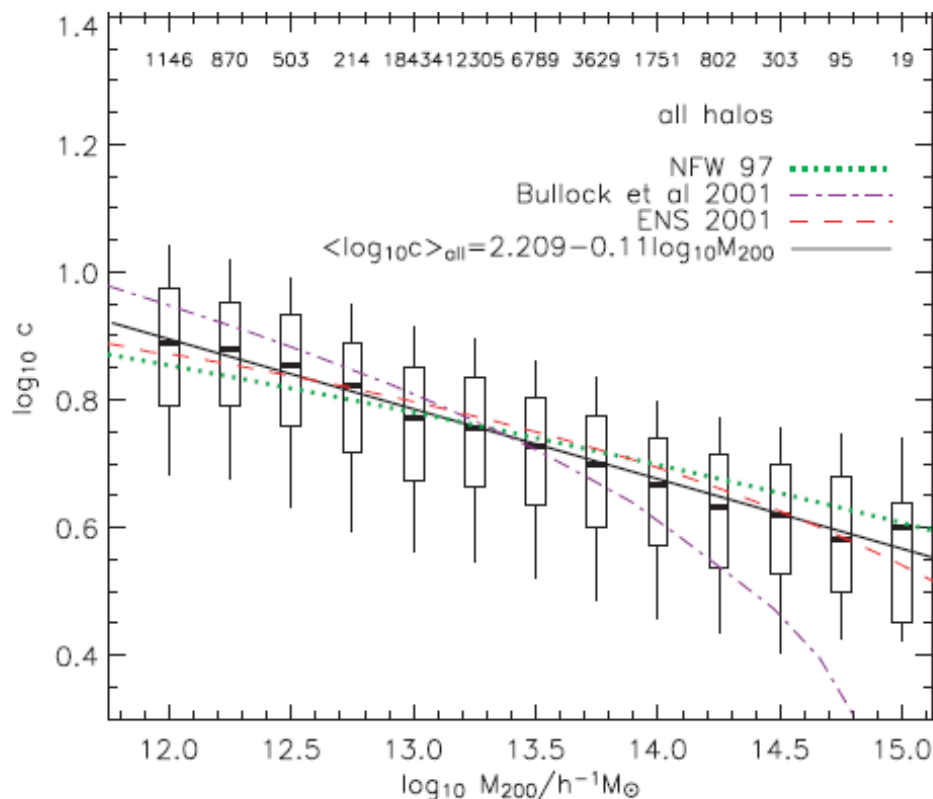
# Key Predictions of nonlinear structure formation models

(2) Halo concentration-mass relation



# Degree of Mass Concentration

$$c_{200} \equiv \frac{r_{200}}{r_s} = \frac{\text{(Virial radius)}}{\text{(Scale radius)}}$$



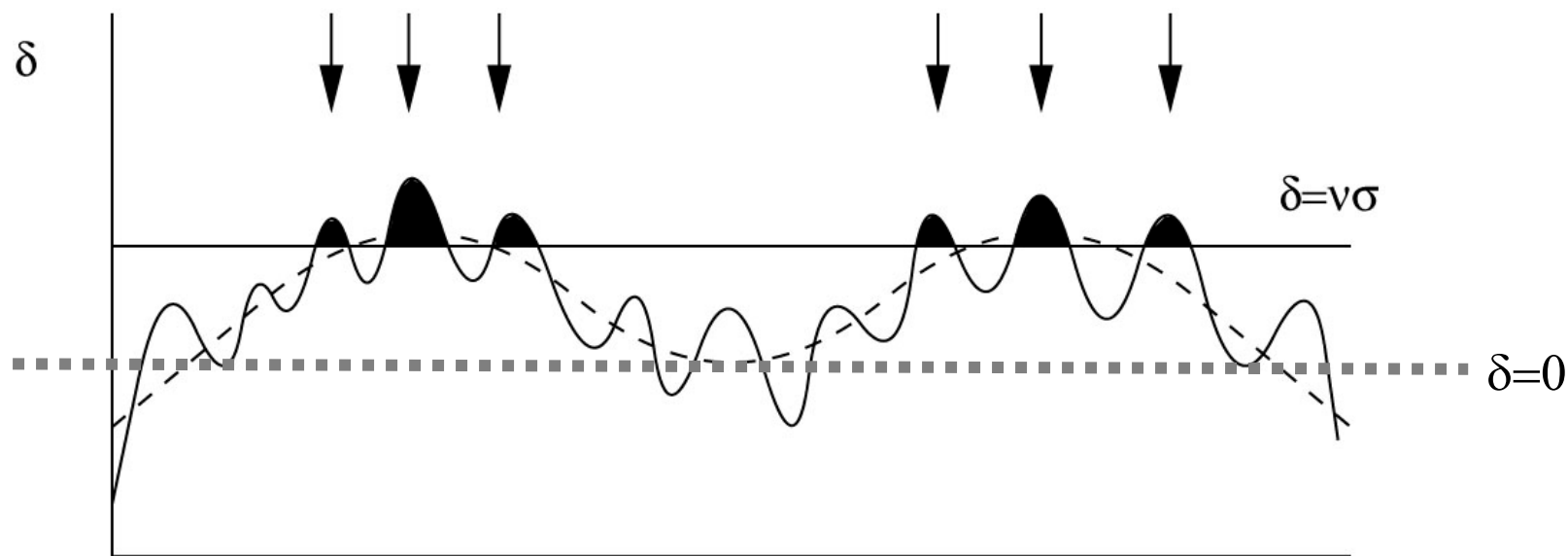
In hierarchical structure formation,  $\langle c \rangle$  is predicted to decrease with increasing  $M$

DM halos that are more massive collapse later on average, when the mean background density of the universe is correspondingly lower (Bullock+01; Neto+07; Duffy+08; Bhattacharya+13)

Clusters (groups) of galaxies are predicted to have  $\langle c \rangle = 3-4$  (5-6)

# Key Predictions of nonlinear structure formation models

## (3) Halo bias: surrounding large-scale structure



$$\delta(\mathbf{x}) := \frac{\rho - \bar{\rho}}{\bar{\rho}} = \int \frac{d^3k}{(2\pi)^3} \tilde{\delta}(\mathbf{k}) e^{i\mathbf{k}\cdot\mathbf{x}} \quad \mathbf{x}$$

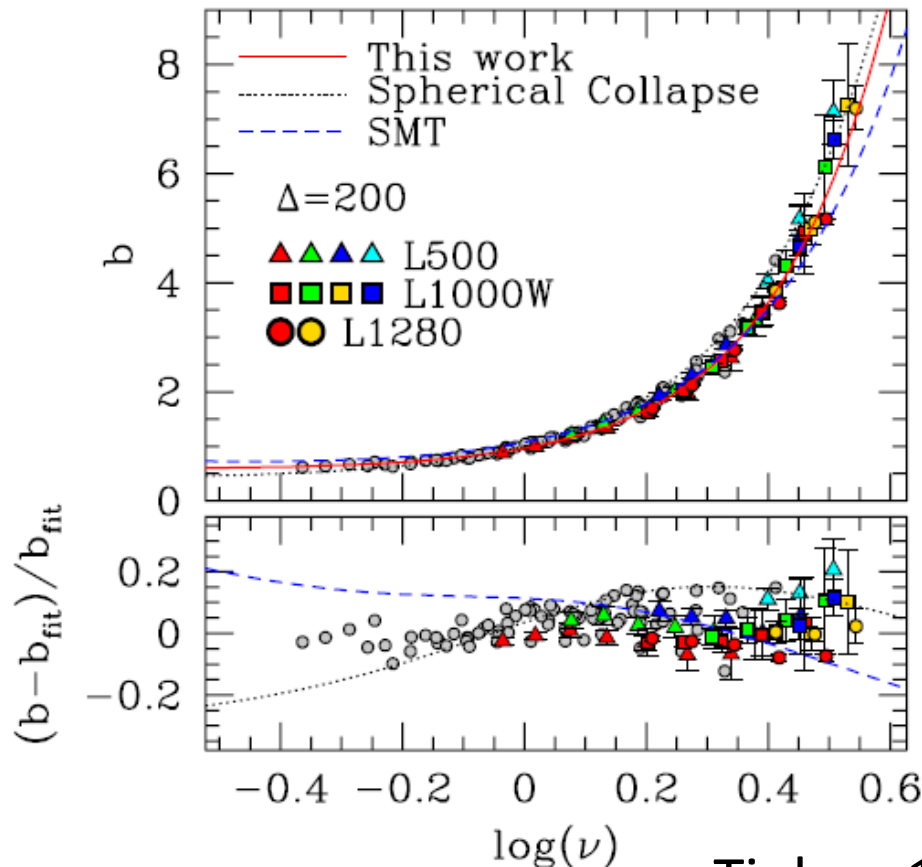
$$\langle \tilde{\delta}(\mathbf{k}) \tilde{\delta}(\mathbf{k}') \rangle = (2\pi)^3 \delta_D^3(\mathbf{k} + \mathbf{k}') P(k)$$

# Halo Bias Factor: $b_h$

Clustering of matter  
around halos with  $M$ :

$$\xi_{\text{hm}}(r|M) \equiv \langle \delta_h(\mathbf{x}|M)\delta_m(\mathbf{x}+\mathbf{r}) \rangle$$

$$= \frac{\langle \rho_{\text{halo}}(r|M) \rangle}{\bar{\rho}} + \underline{b_h(M)\xi_{\text{mm}}(r)} \quad \text{2h term}$$



**Correlated matter distribution (2h term)**

**Matter correlation function:**

$$\xi_{\text{mm}}(\mathbf{r}) \equiv \langle \delta_m(\mathbf{x})\delta_m(\mathbf{x}+\mathbf{r}) \rangle = \int \frac{d^3k}{(2\pi)^3} P(k) e^{i\mathbf{k}\cdot\mathbf{r}}$$

$$\propto \sigma_8^2$$

**Linear halo bias:**

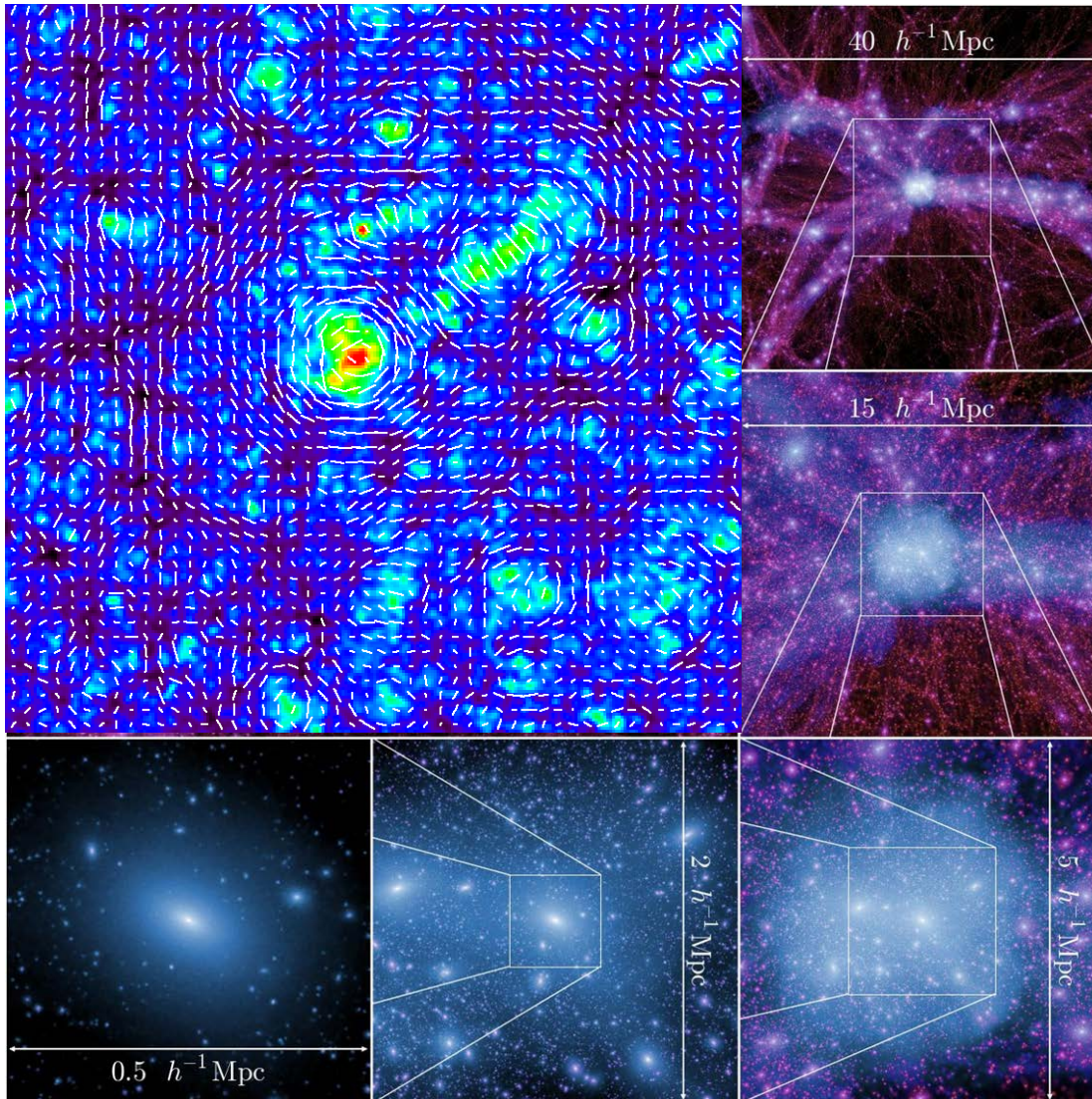
$$b_h(\nu) \approx 1 + \frac{\nu^2 - 1}{\delta_c}$$

$$\nu \equiv \frac{\delta_c}{\sigma(M, z)} \sim 3 - 4 \text{ for clusters}$$

Tinker+10 LCDM simulations



# 2. Cluster Weak Gravitational Lensing



## Key Objectives

### Cluster structure (1h)

*Halo mass,  $M_{200}$*

*Halo density profile,  $\rho(r)$*

*c-M relation,  $c(M,z)$*

### Surrounding LSS (2h)

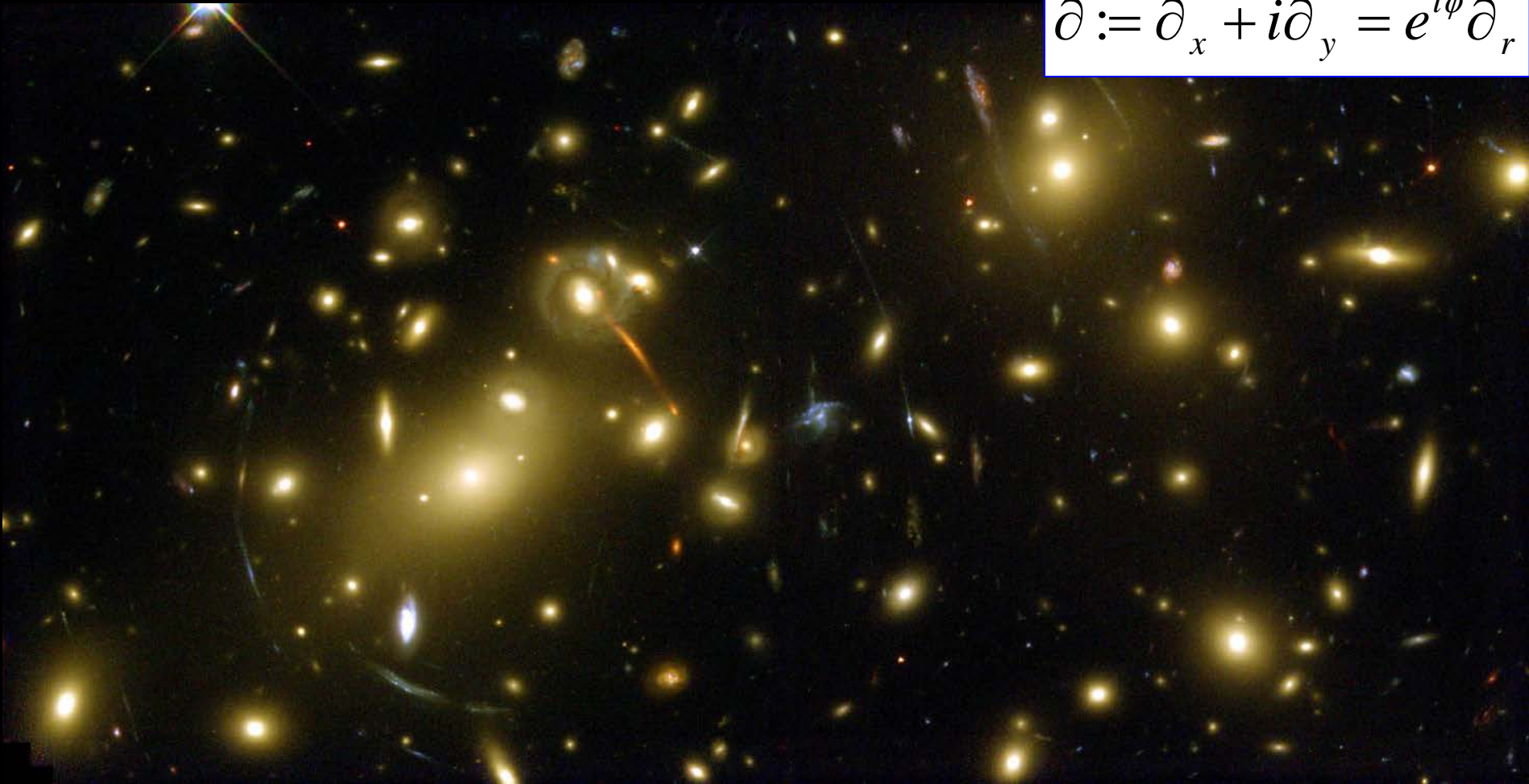
*Halo bias  $b_h(M,z)$*

*Clustering strength  $\sigma_8$*

# Gravitational Shear

$$\gamma = \partial\bar{\partial}\Psi / 2$$

$$\partial := \partial_x + i\partial_y = e^{i\phi}\partial_r$$





# Gravitational Magnification

$$\kappa = \partial\bar{\partial}^*\Psi / 2 = \Delta\Psi / 2$$

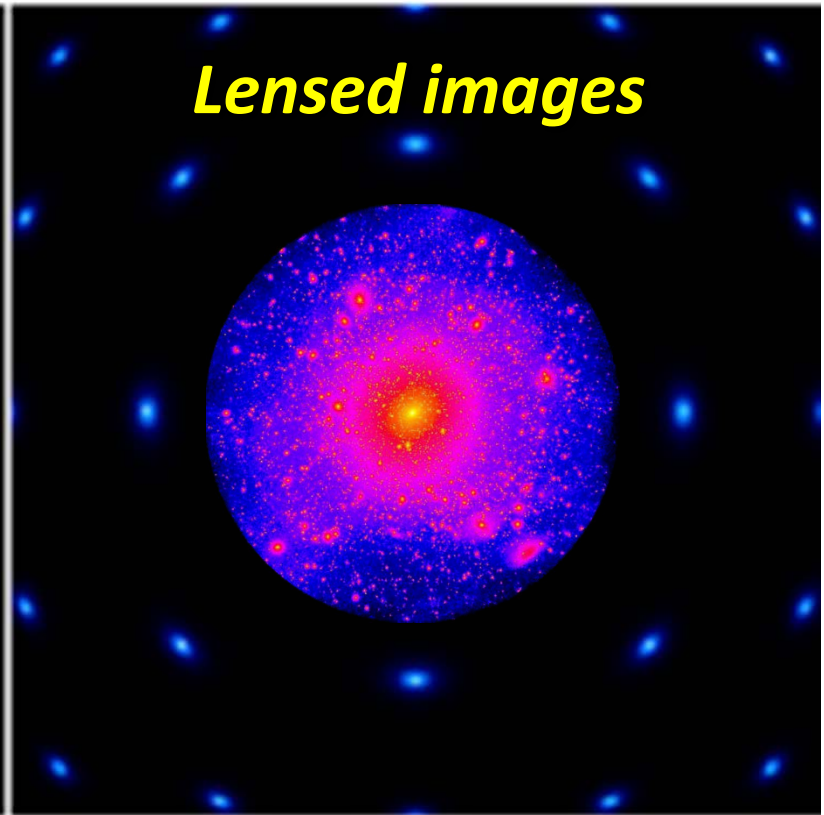
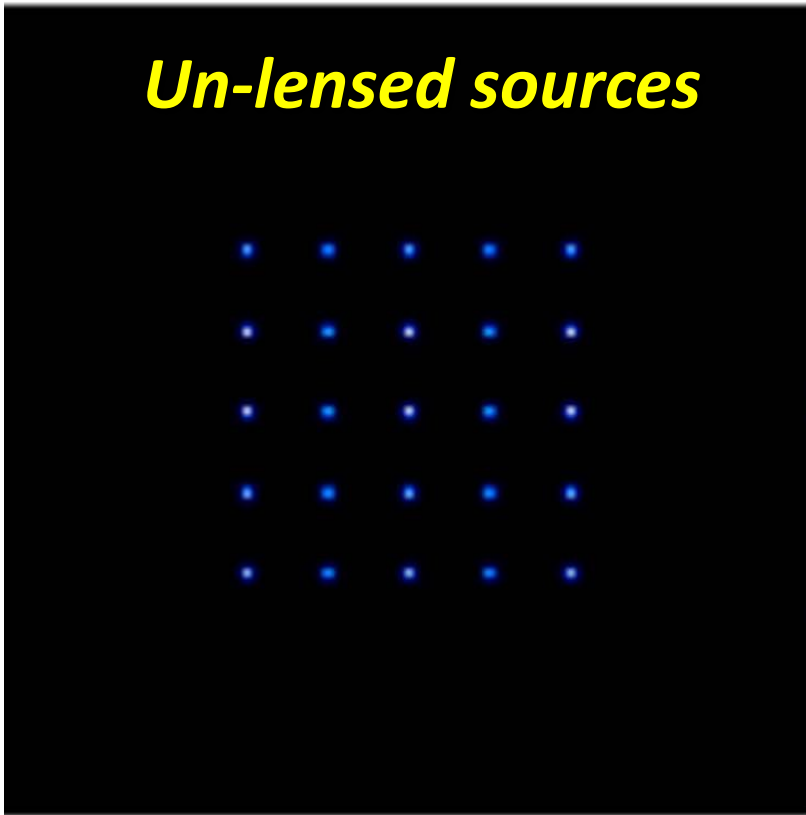
$$\partial := \partial_x + i\partial_y = e^{i\phi}\partial_r$$

MACSJ1149 (z=0.54)

Zheng+CLASH. 2012, *Nature*, 489, 406



# Shear and Magnification Effects



- **Shear**

✓ Shape distortion:  $\delta e_+ \sim \gamma_+$

*Sensitive to “modulated” matter density*

$$\Sigma_{\text{crit}} \gamma_+ = \Delta \Sigma(R) \equiv \Sigma(< R) - \Sigma(R)$$

- **Magnification**

✓ Flux amplification:  $\mu F$

✓ Area distortion:  $\mu \Delta \Omega$

*Sensitive to “total” matter density*

$$\mu \approx 1 + 2\kappa; \quad \Sigma_{\text{crit}} \kappa = \Sigma(R)$$

# Tangential Shear, $\gamma_+$

A measure of azimuthally-averaged tangential coherence of elliptical distortions around a given point (Kaiser 95):

$$\gamma_+(R) = \Delta\Sigma_+(R) / \Sigma_{\text{crit}}$$

$$(\Gamma_+)_{ij} = \left( \delta_i \delta_j - \frac{1}{2} \Delta^{(2)} \delta_{ij} \right) \psi_+$$

$$\gamma_\times(R) = 0$$

$$(\Gamma_\times)_{ij} = (\epsilon_{kj} \partial_i \partial_k - \epsilon_{ki} \partial_j \partial_k) \psi_\times$$

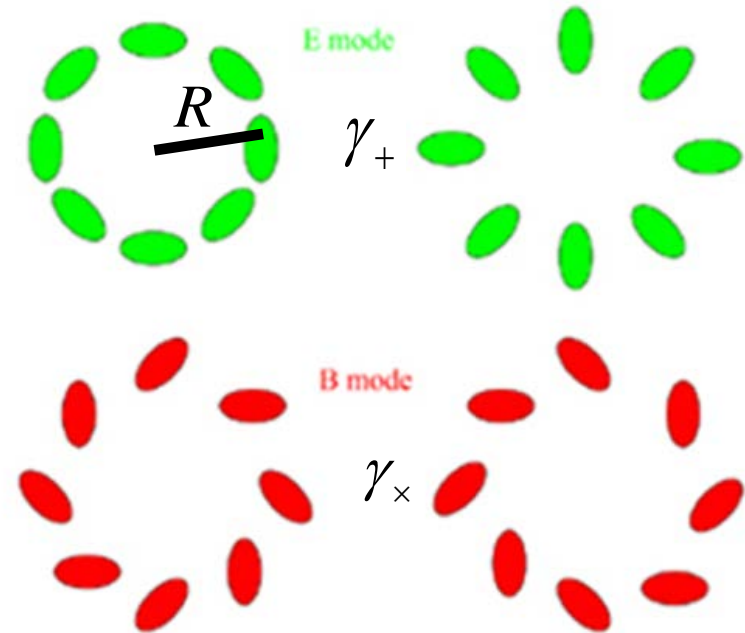
$\Delta\Sigma(R)$  is radially-modulated surface mass density:

$$\Delta\Sigma_+(R) = \Sigma(< R) - \Sigma(R)$$

Sensitive to interior mass

$$\Sigma(\mathbf{R}) = \int \delta\rho_m(\mathbf{r}) dx_{\parallel}$$

$$\Sigma_{\text{crit}} = \frac{c^2}{4\pi G} \frac{D_S}{D_L D_{LS}}$$

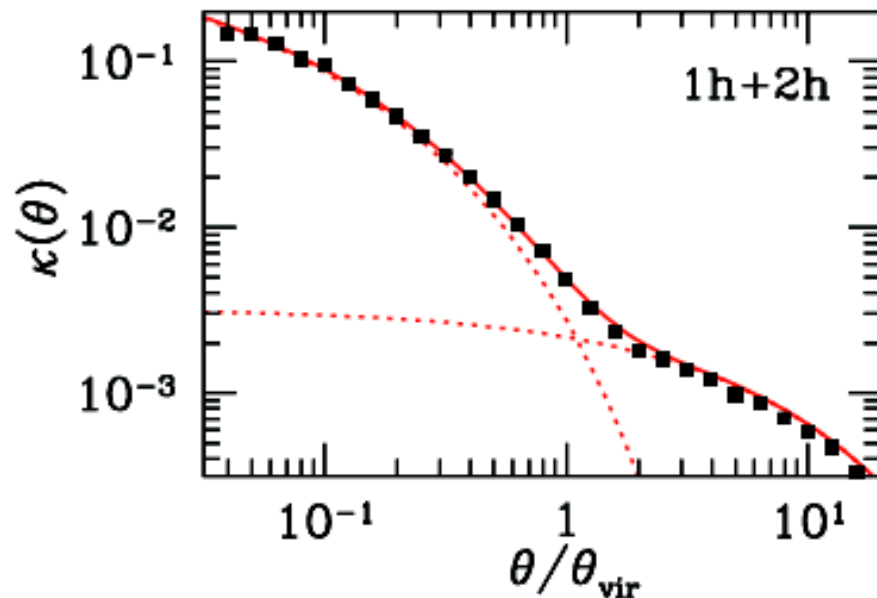


# Shear doesn't see mass sheet

Averaged lensing profiles in/around LCDM halos (Oguri+Hamana 11)

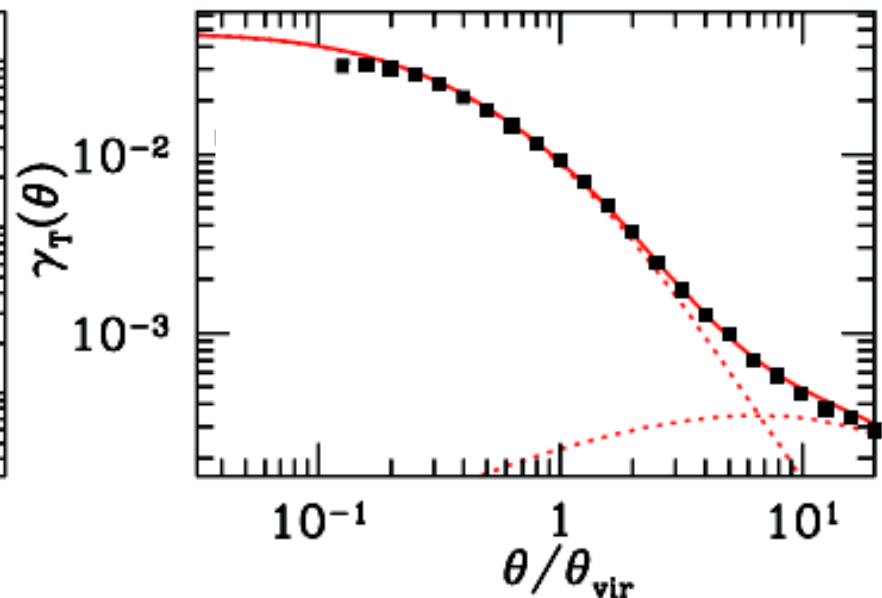
Total

$$\kappa = \Sigma(R) / \Sigma_{\text{crit}}$$



Modulated

$$\gamma_+ = \Delta\Sigma_+(R) / \Sigma_{\text{crit}}$$



- Tangential shear is a powerful probe of **1-halo term**, or **internal halo structure**.
- Shear alone cannot recover absolute mass, known as **mass-sheet degeneracy**:

$\gamma$  remains unchanged by  $\kappa \rightarrow \kappa + \text{const.}$



# Combining Shear and Magnification

**Bayesian joint likelihood approach** (Umetsu+11a; Umetsu 13)

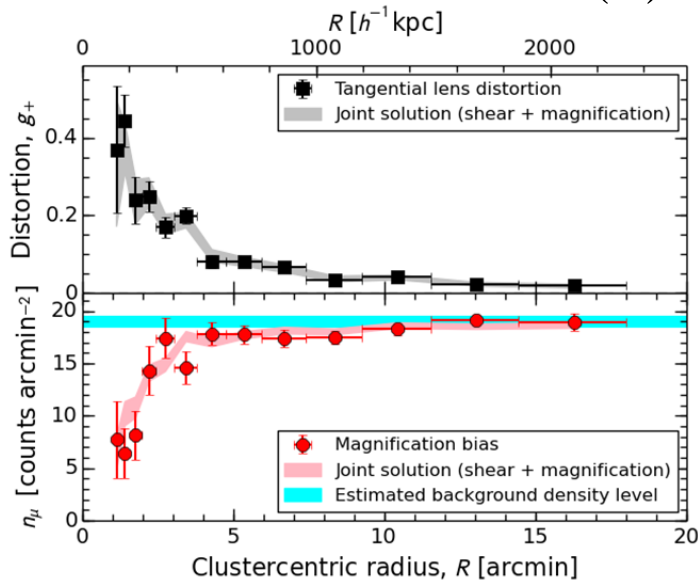
Tangential distortion

Inverse magnification

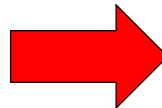
$$g_+(R) = \frac{\kappa(< R) - \kappa(R)}{1 - \kappa(R)},$$

$$\mu^{-1}(R) = [1 - \kappa(R)]^2 - [\kappa(< R) - \kappa(R)]^2$$

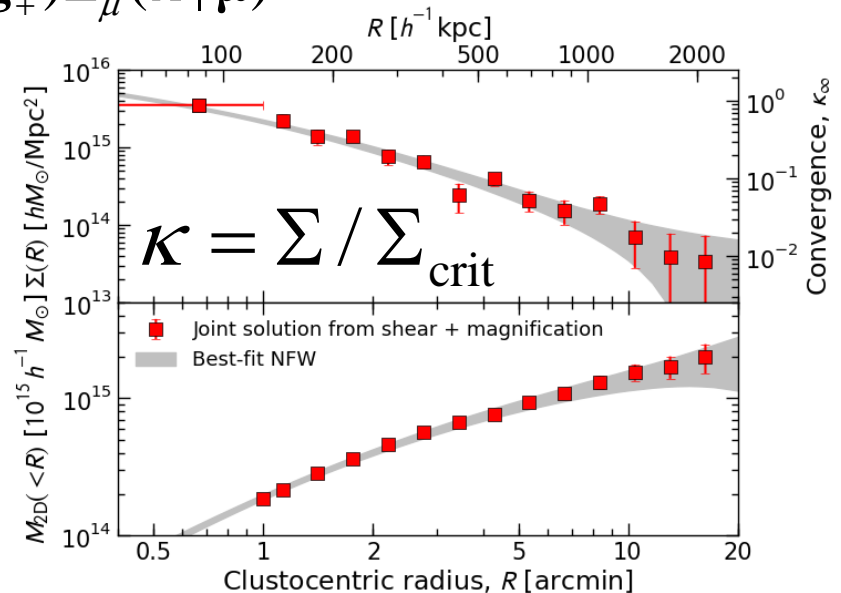
$$L(\boldsymbol{\kappa}) = L_g(\boldsymbol{\kappa} | \mathbf{g}_+) L_\mu(\boldsymbol{\kappa} | \boldsymbol{\mu})$$



$g_+$



$\mu$



- Mass-sheet degeneracy broken
- Total statistical precision improved by  $\sim 20\text{-}30\%$
- Calibration uncertainties marginalized over:  $c = \{ \langle W \rangle_s, f_{W,s}, \langle W \rangle_\mu, \bar{n}_\mu, s_{\text{eff}} \}$ .

# Cluster Lensing And Supernova survey with Hubble

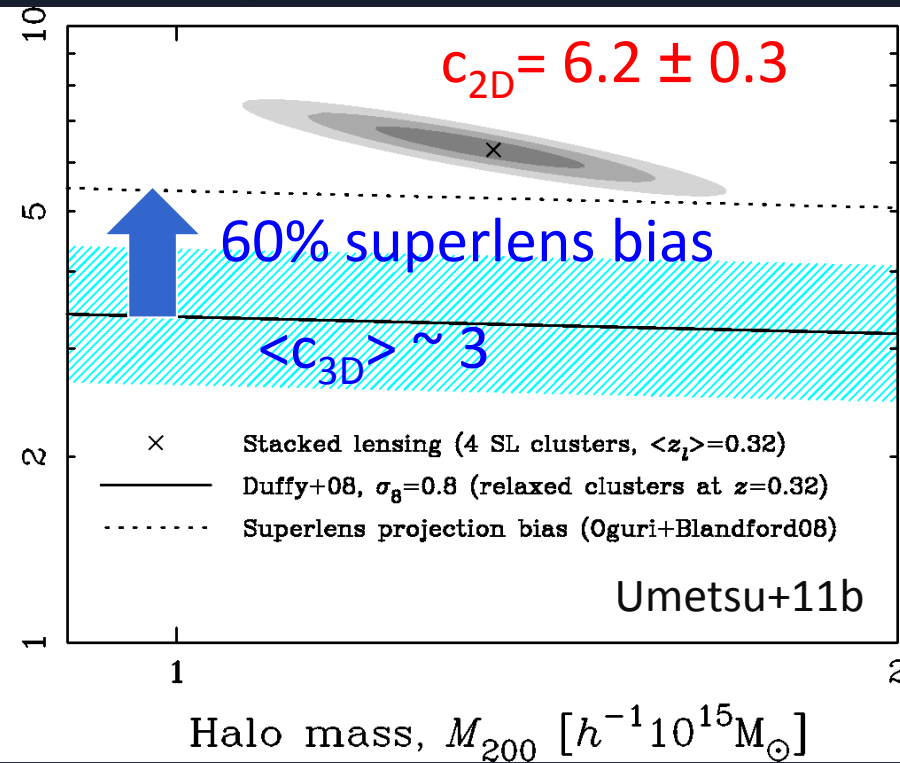
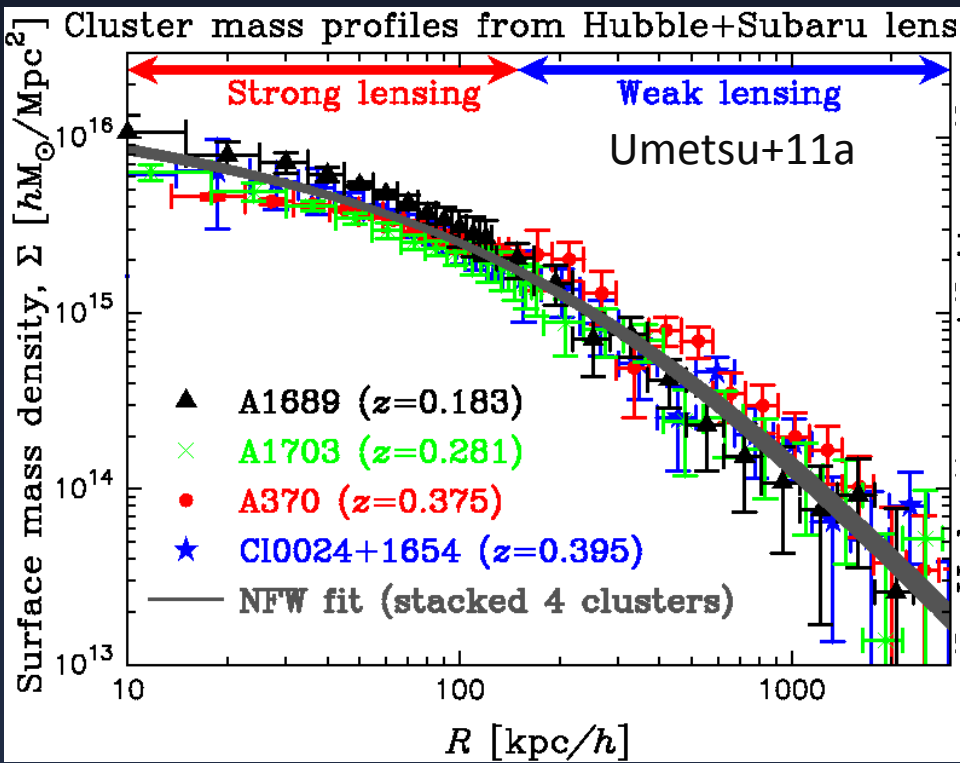


PI. Marc Postman (STScI)

<http://www.stsci.edu/~postman/CLASH/Home.html>

# CLASH Objectives & Motivation

Before CLASH (2010), deep-multicolor Strong (*HST*) + Weak (*Subaru*) lensing data only available for a handful of “super lens” clusters

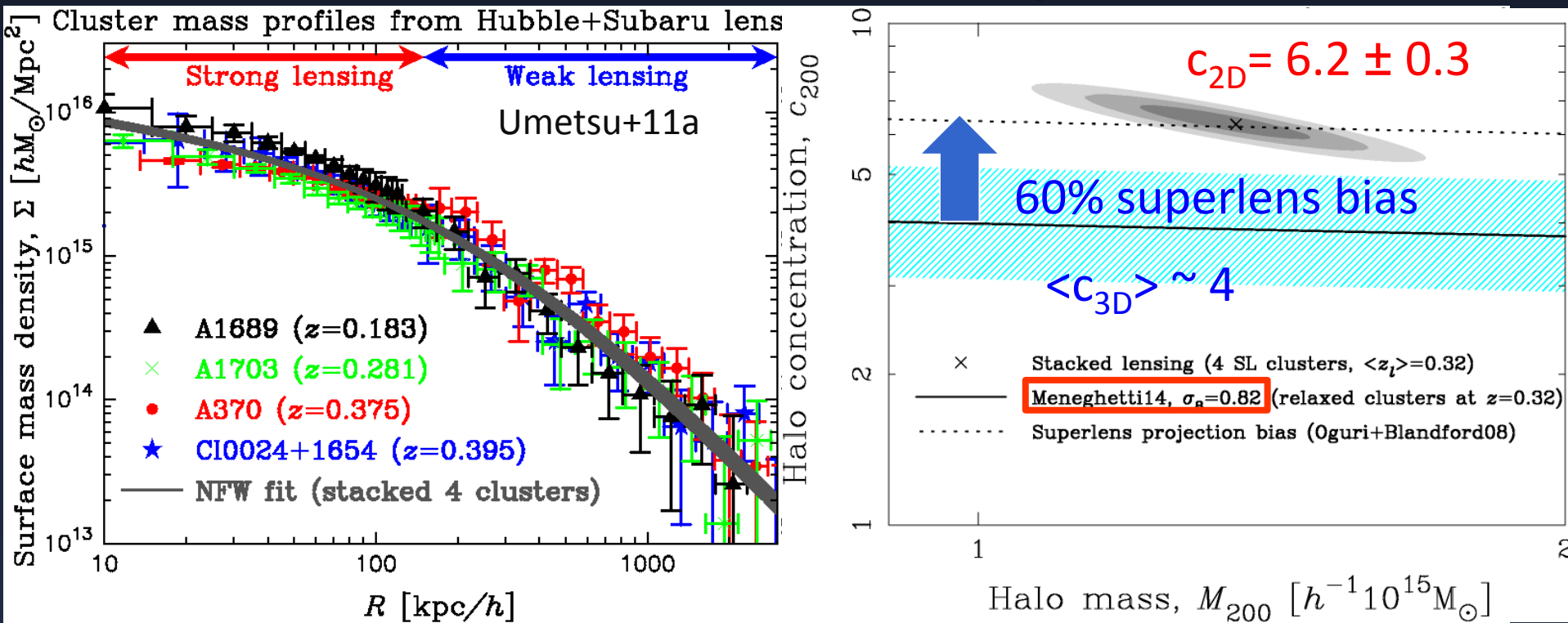


**Total mass profile shape:** consistent w self-similar NFW (cf. Newman+13; Okabe+13)

**Degree of concentration:** predicted superlens correction not enough if  $\langle c_{\text{LCDM}} \rangle \sim 3$ ?

# CLASH Objectives & Motivation

Before CLASH (2010), deep-multicolor Strong (*HST*) + Weak (*Subaru*) lensing data only available for a handful of “super lens” clusters



**Total mass profile shape:** consistent w self-similar NFW (cf. Newman+13; Okabe+13)

**Degree of concentration:** predicted superlens correction is just enough if  $\langle c_{\text{LCDM}} \rangle \sim 4$



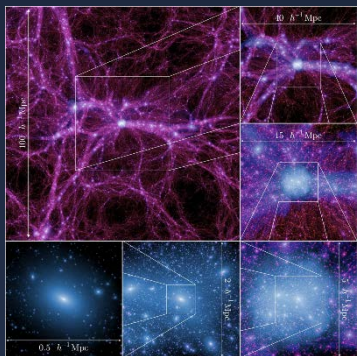


# CLASH: Observational + Theory Efforts

A 524-orbit *HST* Treasury Program to observe 25 clusters in 16 filters (0.23-1.6  $\mu\text{m}$ ) (Postman+CLASH 12)



**Wide-field Subaru imaging** (0.4 - 0.9  $\mu\text{m}$ ) plays a unique role in complementing deep *HST* imaging of cluster cores (Umetsu et al. 2014, *ApJ*, 795, 163)

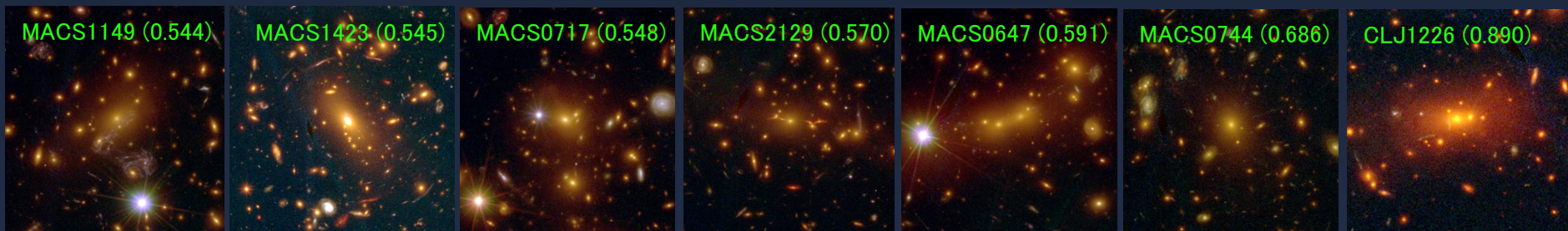
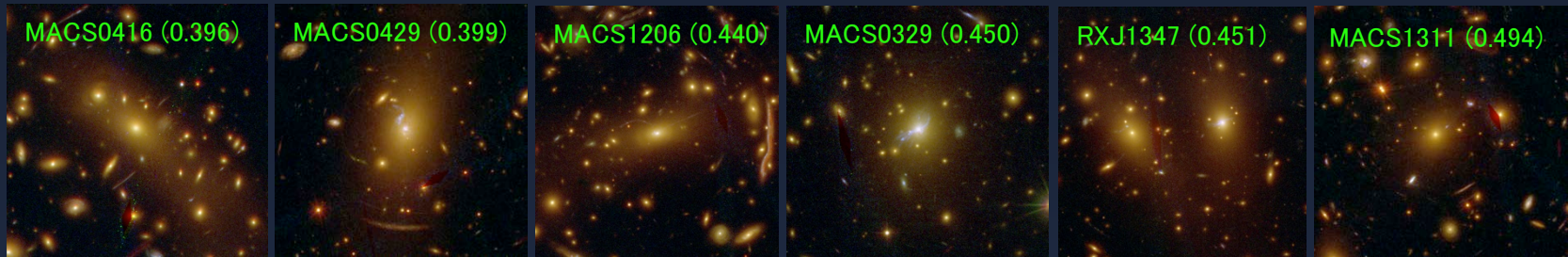
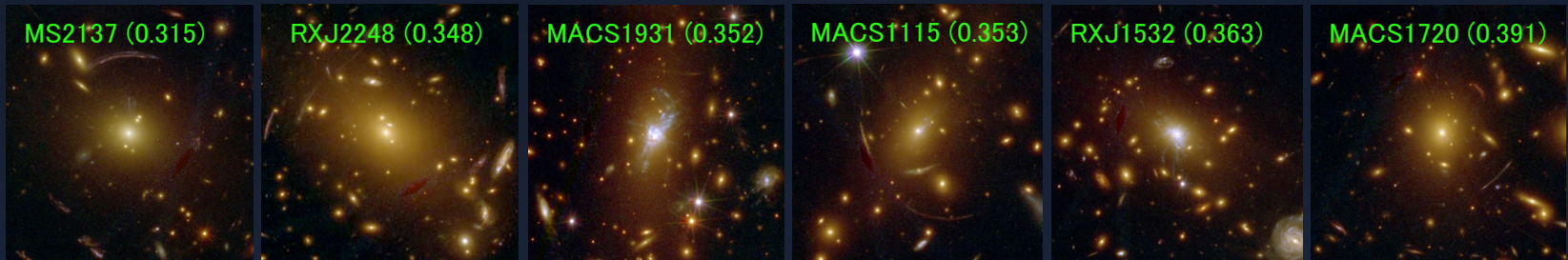


**MUSIC-2** (hydro + N-body re-simulation) provides an accurate characterization of CLASH sample with testable predictions (Meneghetti et al. 2014, *ApJ*, 797, 34)





# The CLASH Gallery (*HST*)

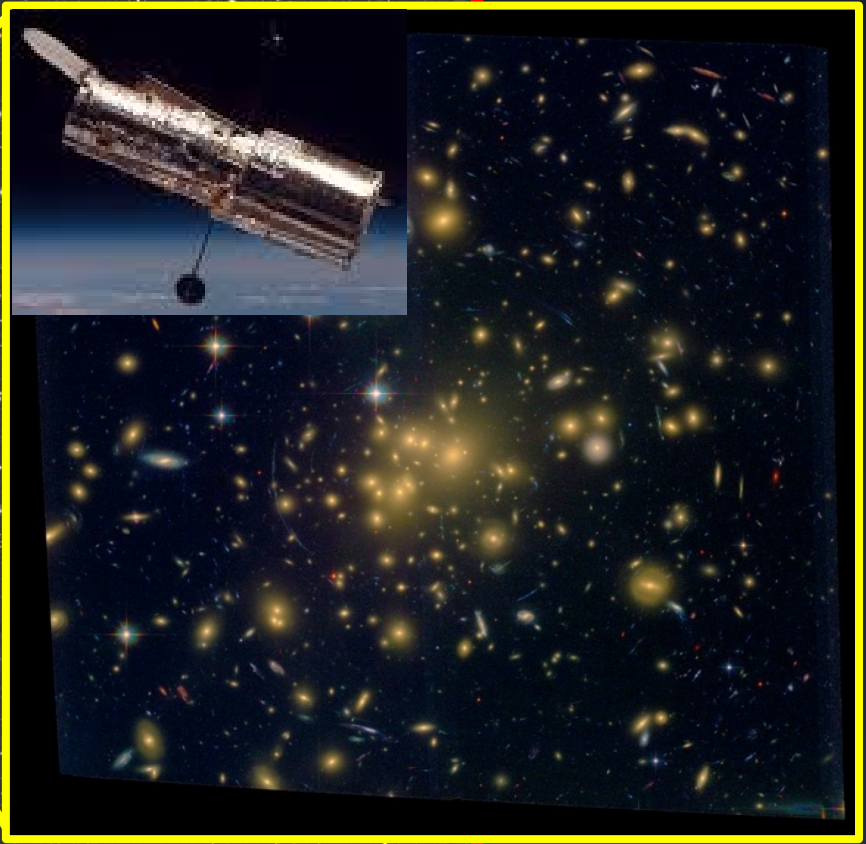


The final HST observation for CLASH was on 9–July–2013 ... 963 days, 15 hrs, 31 min after first obs.



***SUBARU* (S-Cam) multi-color  
imaging for wide-field weak**

**High-resolution space imaging  
with *HST* (ACS/WFC3) for  
strong lensing**



**34 arcmin**



# CLASH X-ray-selected Subsample ( $0.18 < z < 0.9$ )

- **X-ray morphology +  $T_x$  selection**

- $T_x > 5\text{keV}$  ( $M_{200} > 5e14 M_{\text{sun}}/h$ )
- Small BCG to X-ray-peak offset,  $\sigma_{\text{off}} \sim 10\text{kpc}/h$
- Smooth regular X-ray morphology

→ **Optimized for radial-profile analysis ( $R > 2\sigma_{\text{off}} \sim 20\text{kpc}/h$ )**

- **CLASH theoretical predictions (Meneghetti+CLASH 14)**

- Composite relaxed (70%) and unrelaxed (30%) clusters
- Mean  $\langle c_{200} \rangle = 3.9$ ,  $\sigma(c_{200}) = 0.6$ ,  $c_{200} = [3, 6]$
- Negligible orientation bias ( $\sim 2\%$  in  $\langle M_{3D} \rangle$ )
- $> 90\%$  of CLASH clusters to have strong-lensing features



# CLASH Weak-Lensing Results (1)

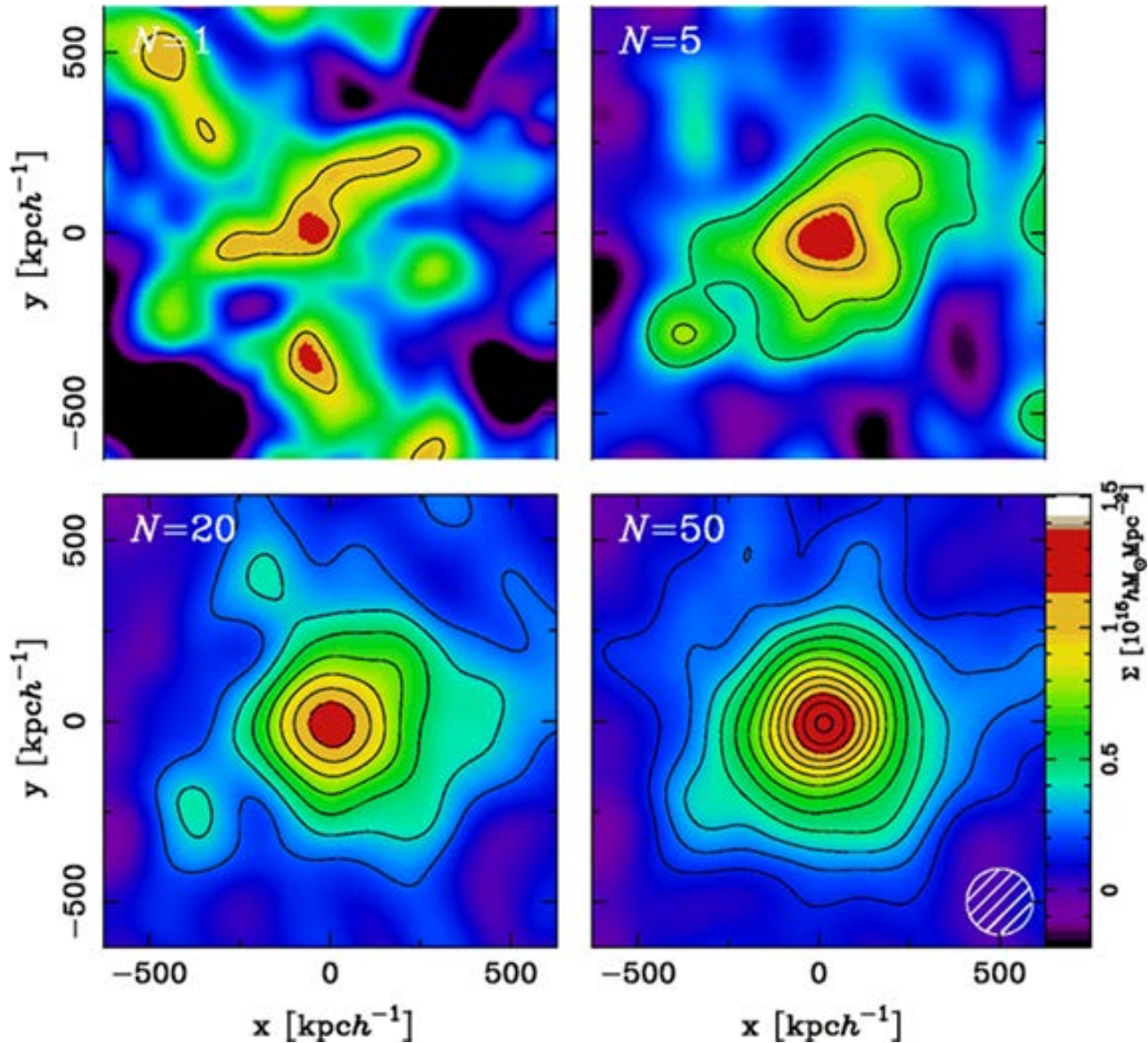
Ensemble-averaged DM halo structure:

- Cluster halo density profile,  $\langle \Delta \Sigma_+(R) \rangle$
- Degree of halo concentration,  $\langle c(M, z) \rangle$

from ***stacked-shear-only WL*** analysis of the X-ray-selected CLASH subsample (16 clusters)



# Power of Stacked WL Analysis





# Averaged Halo Density Profile $\Delta\Sigma_+(R)$

Stacking WL-shear signals of individual clusters by

$$\langle\langle \widehat{\Delta\Sigma_+} \rangle\rangle = \left( \sum_n \mathcal{W}_{+n} \right)^{-1} \left( \sum_n \mathcal{W}_{+n} \widehat{\Delta\Sigma_{+n}} \right),$$

*Summing over clusters ( $n=1, 2, \dots$ )*

with individual “sensitivity” matrix

$$(\mathcal{W}_{+n})_{ij} \equiv \Sigma_{c,n}^{-2} (C_{+n}^{-1})_{ij}$$

defined with total covariance matrix

$$C_+ = C_+^{\text{stat}} + C_+^{\text{sys}} + C_+^{\text{lss}}.$$

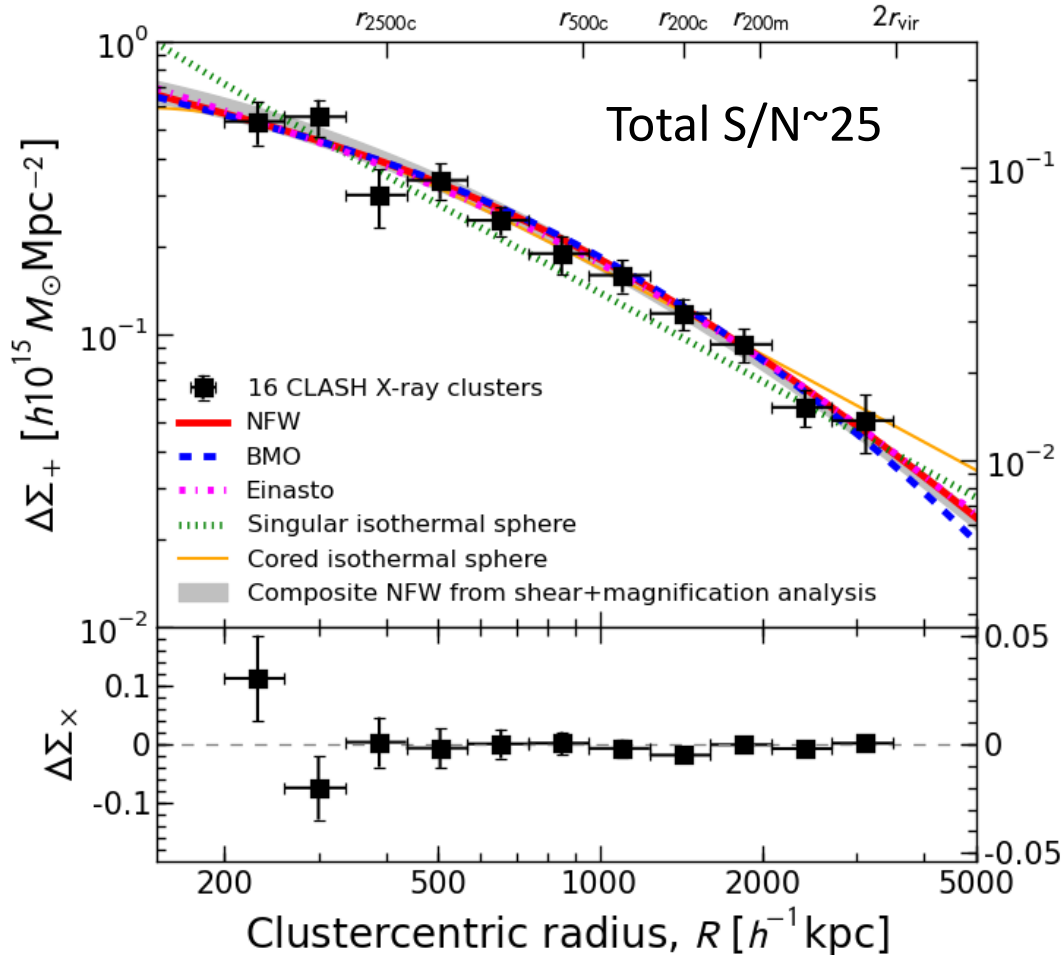
**With “trace-approximation”, averaging (stacking) is interpreted as**

$$\langle\langle \Sigma_c^{-1} \rangle\rangle = \frac{\sum_n \text{tr}(\mathcal{W}_{+n}) \Sigma_{c,n}^{-1}}{\sum_n \text{tr}(\mathcal{W}_{+n})},$$

Umetsu et al. 2014,  
*ApJ*, 795, 163



# CLASH Averaged Halo Density Profile



Stacked-shear-only analysis provides a net 1-halo-only constraint ( $\gamma_{+,2h} < 10^{-3}$ ) at  $\langle z \rangle = 0.35$

**NFW** an excellent fit (PTE = 0.66)

- $M_{200} = (1.3 \pm 0.1) 10^{15} M_{\text{sun}}$
- $c_{200} = 4.01 (+0.35, -0.32)$

**Einasto** (PTE=0.51)

- Einasto shape parameter  $\alpha_E = 0.19 \pm 0.07$ , consistent with NFW curvature ( $\alpha_E \sim 0.18$ )

Umetsu+CLASH 14, *ApJ*, 795, 163

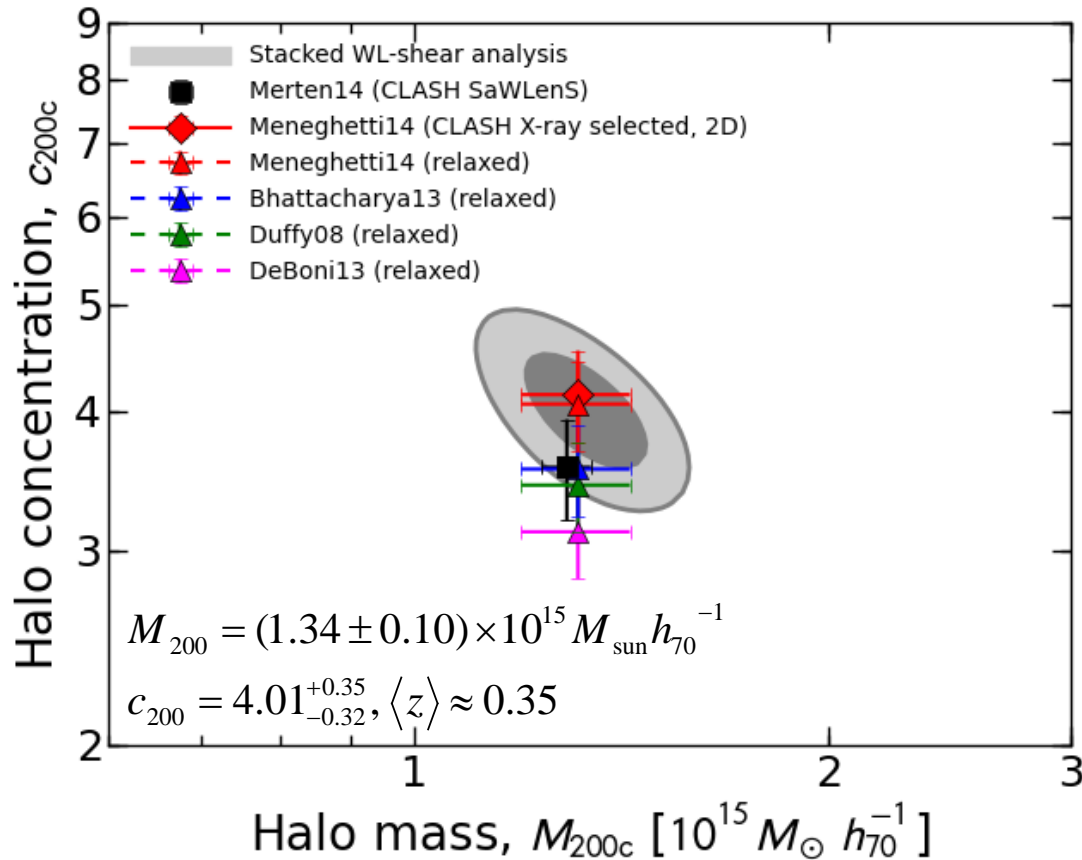
Consistent w a family of density profiles for collisionless DM halos (NFW, truncated variants of NFW, Einasto)





# Integrated Constraints on $c(M, z)$

$$\langle c_{200c} \rangle = \frac{\int dM dz N(M, z) \hat{c}_{200c}(M, z)}{\int dM dz N(M, z)} \approx \frac{\sum_n \text{tr}(\mathcal{W}_n) \hat{c}_{200c}(M_n, z_n)}{\sum_n \text{tr}(\mathcal{W}_n)}$$



Variance in theory due primarily to different cosmology ( $\sigma_8$ )

Meneghetti14:  $\sigma_8=0.82$

Bhattacharya13:  $\sigma_8=0.8$

Duffy08:  $\sigma_8=0.796$

DeBoni13:  $\sigma_8=0.776$

- Excellent agreement with LCDM predictions for CLASH (M14),  $\langle c_{200} \rangle = 3.9$
- Consistent with Bhatt13, Duffy08 LCDM predictions for relaxed halos @  $1\sigma$ ,  $\langle c_{200} \rangle \sim 3.6$



# CLASH Weak-Lensing Results (2)

Individual cluster mass profiles:

- Cluster mass profile  $\Sigma(R)$  reconstruction
- Spherical mass estimates ( $M_{500}$ ,  $M_{200}$ ,  $M_{\text{vir}}$ , ...)

from **joint *shear+magnification* WL analysis** of all CLASH clusters



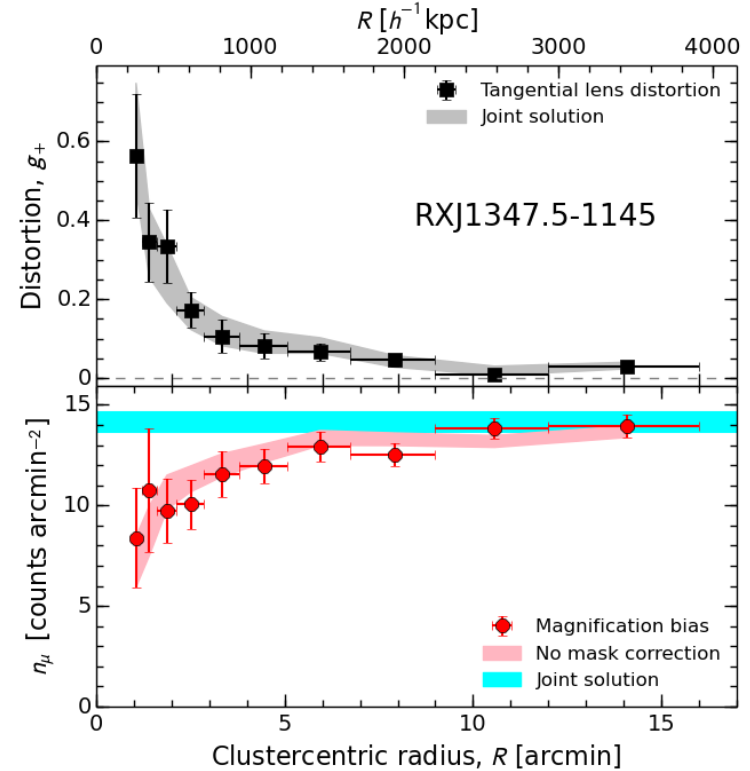
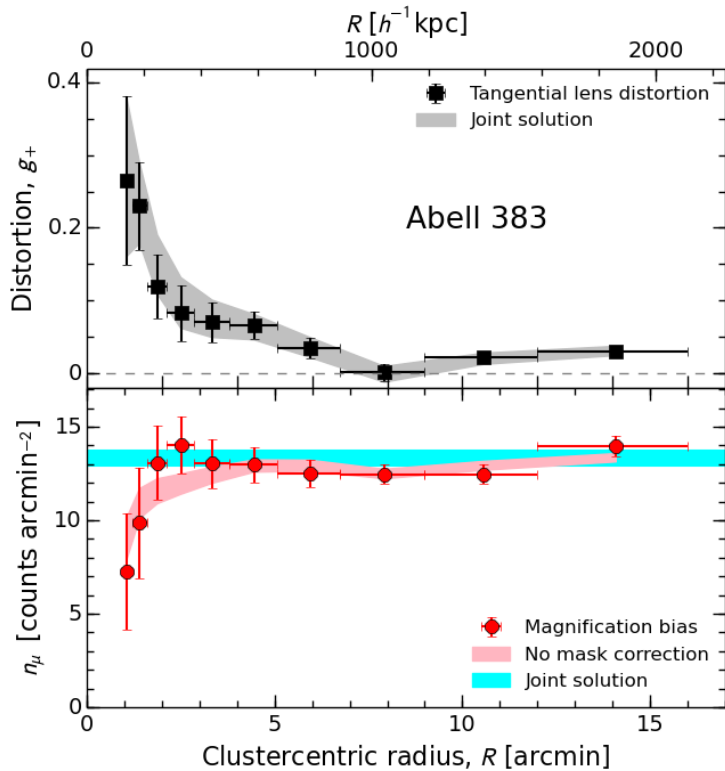
# Joint Shear + Magnification WL Analysis

**CLASH low mass**

$M_{200}=6e14M_{\text{sun}}/h$  ( $z=0.19$ )

**CLASH high mass**

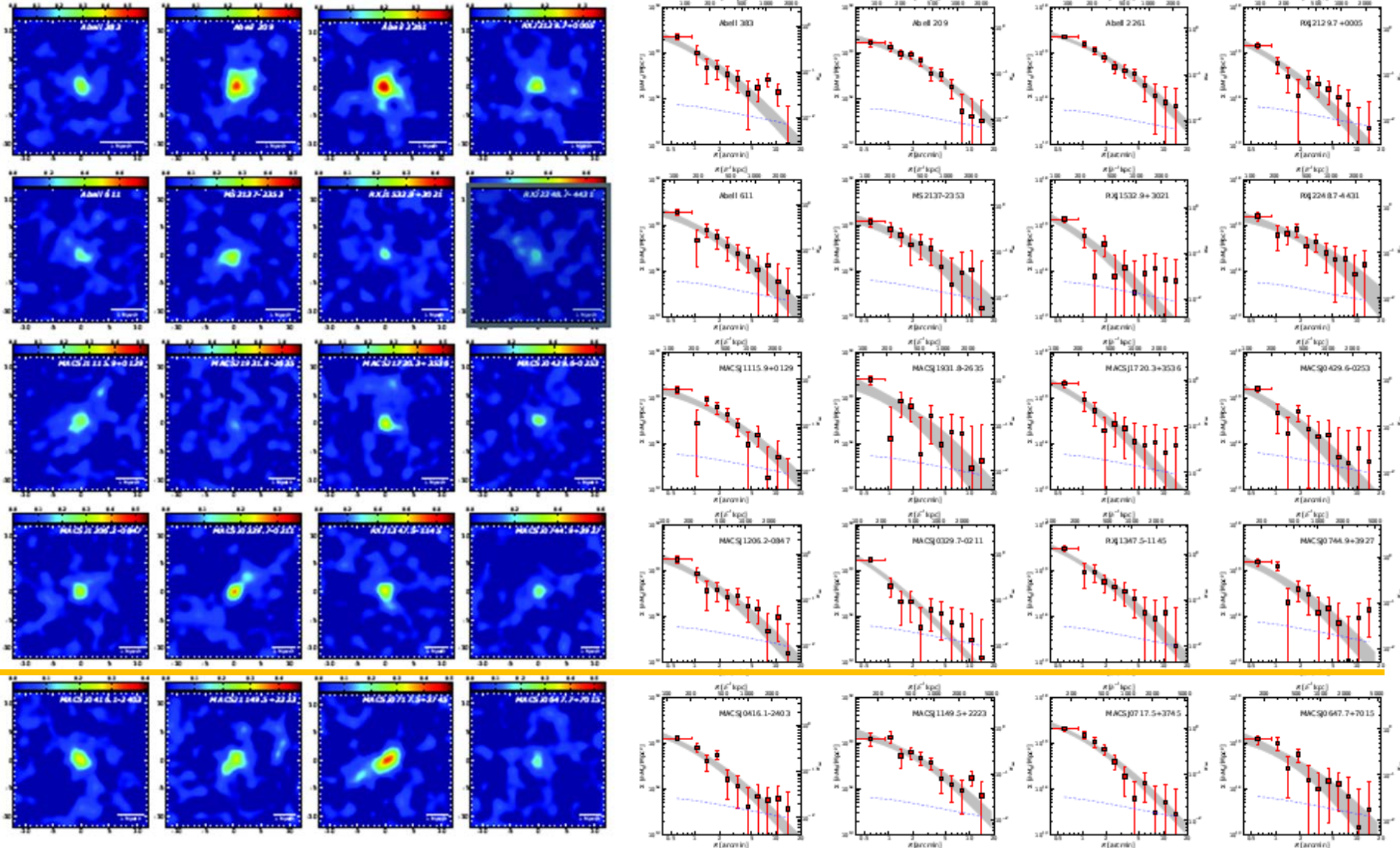
$M_{200}=20e14M_{\text{sun}}/h$  ( $z=0.45$ )



Shear-magnification consistency:  $\langle \chi^2/\text{dof} \rangle = 0.92$  for 20 CLASH clusters



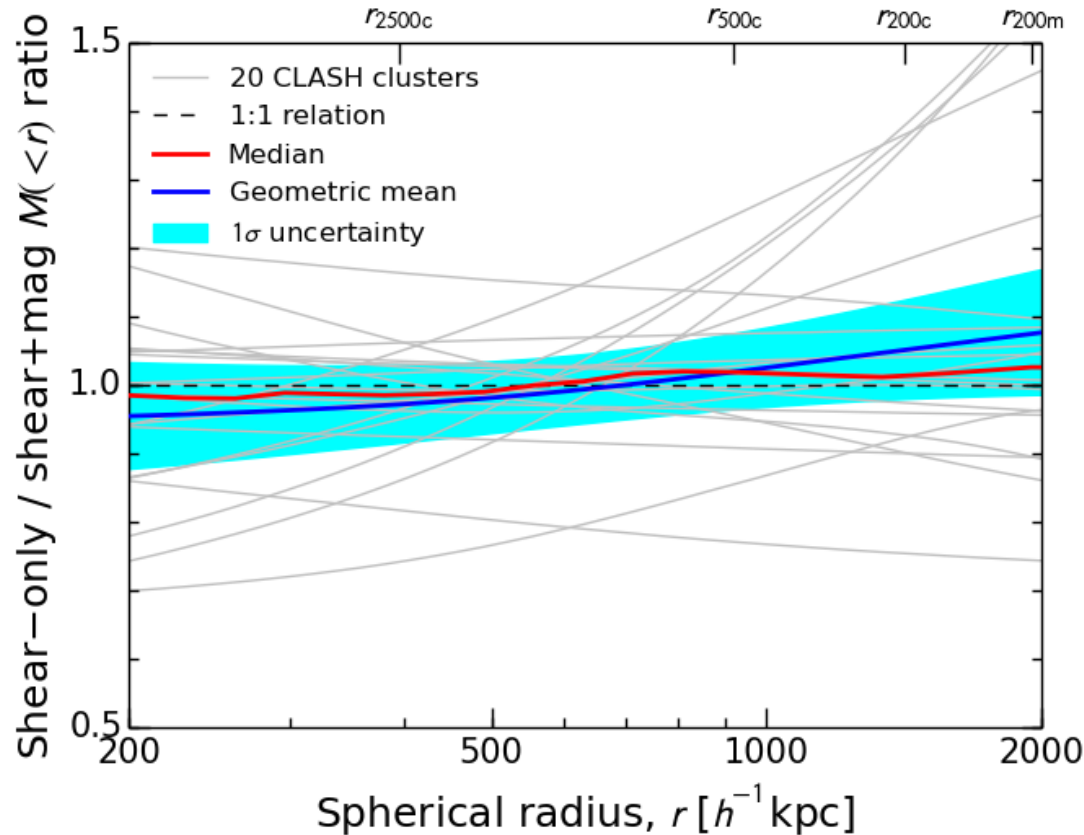
# CLASH Mass Density Profile Dataset





# Shear-Magnification Consistency

$M(<r)$  de-projected assuming spherical NFW density profiles



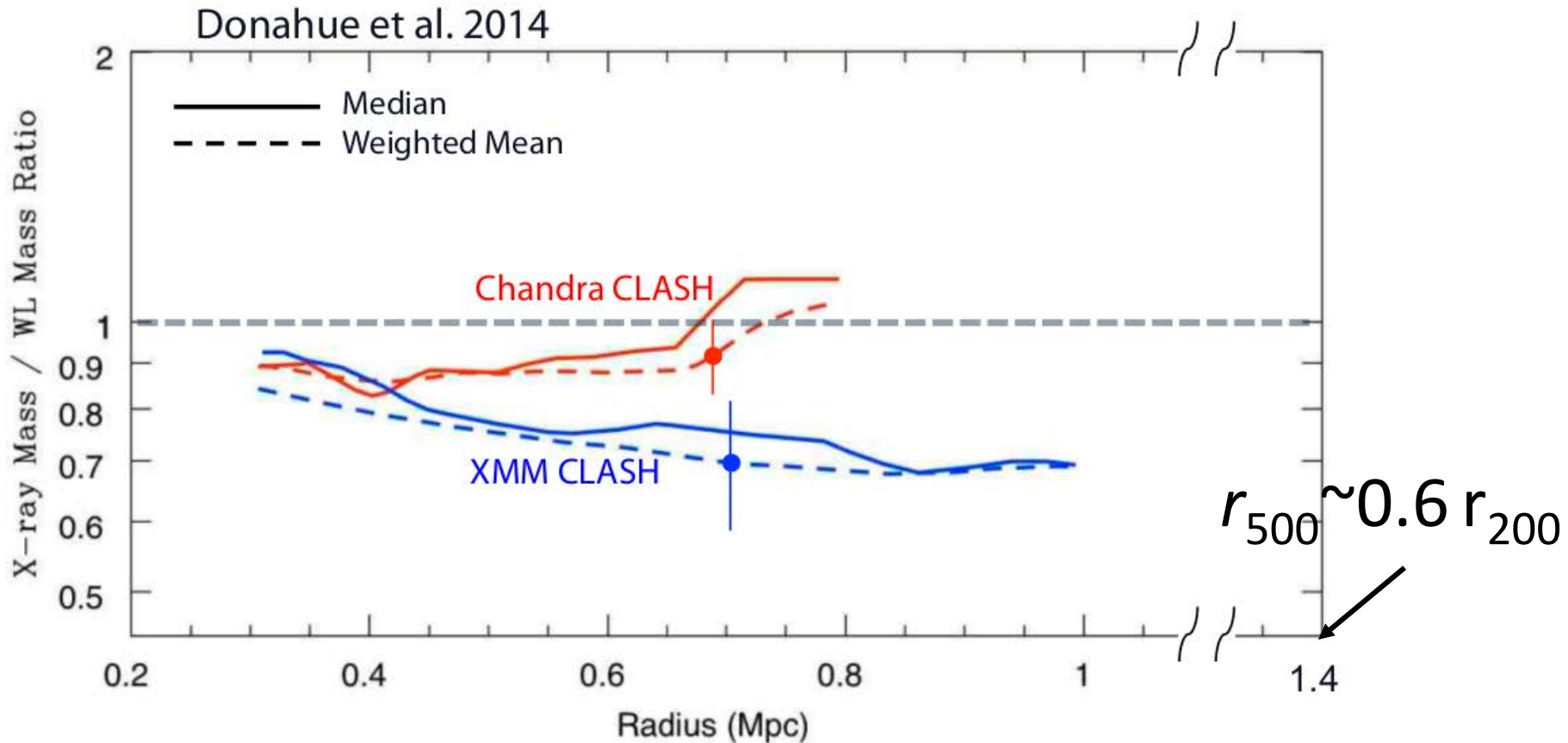
Umetsu et al. 14,  
*ApJ*, 795, 163

Internal systematic uncertainty in the overall mass calibration,  
empirically derived to be about +/- 8%





# CLASH: WL vs. X-ray Mass Comparison



Donahue et al. 2014,  
*ApJ*, 794, 136

## X-ray to WL mass comparison at $r_{500}$

- $b = 1 - \langle M_{\text{Chandra}} / M_{\text{WL}} \rangle = 0.22 \pm 0.10$
- $b = 1 - \langle M_{\text{XMM}} / M_{\text{WL}} \rangle = 0.44 \pm 0.06$



# Full-Lensing Analysis: Strong-lensing, Weak-lensing Shear and Magnification

**Adding “Strong Lensing” to provide tighter constraints on the inner density profile ( $R < 200 \text{kpc}/h$ )**

Multi-probelensing method by Umetsu 2013, *ApJ*, 769, 13

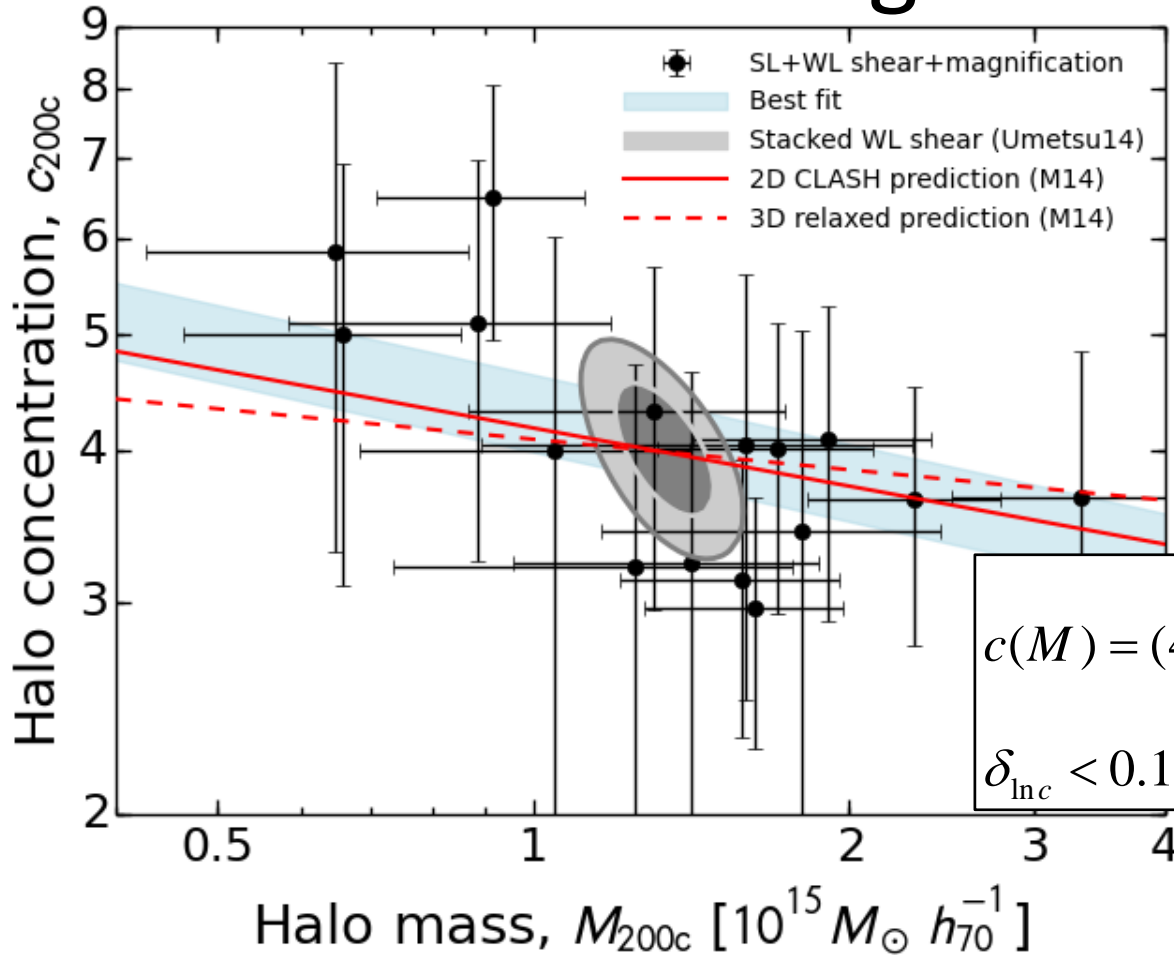
Direct reconstruction of individual mass profiles  $\Sigma(R)$  from **full likelihood analysis of SL, WL shear and magnification** constraints from the CLASH survey

$$L(\boldsymbol{\kappa}) = L_g(\boldsymbol{\kappa} | \mathbf{g}_+) L_\mu(\boldsymbol{\kappa} | \boldsymbol{\mu}) L_{\text{SL}}(\boldsymbol{\kappa} | \mathbf{M}_{\text{proj}})$$

- CLASH-WL shear & magnification constraints from *Subaru* observations (Umetsu+14, *ApJ*, 795, 163)
- CLASH-SL projected mass constraints  $M_{\text{proj}}$  from *HST* observations (Zitrin+14, arXiv:1411.1414)



# CLASH $c$ - $M$ relation from SL, WL shear and magnification



Maximum-likelihood inference

$$\ln \hat{c}(M_i) = A \ln M_i + B$$

$$-2 \ln L = \sum_{i=1}^N \left\{ \frac{[\ln c_i - \ln \hat{c}(M_i)]^2}{\sigma_i^2} + \ln \sigma_i^2 \right\},$$

$$\sigma_i^2 = \sigma_{\ln c_i}^2 + A^2 \sigma_{\ln M_i}^2 - 2 \sigma_{\ln c_i} \sigma_{\ln M_i} + \delta_{\ln c, \text{int}}^2$$

$$c(M) = (4.1 \pm 0.3) \left( \frac{M}{1.3 \times 10^{15} M_{\text{sun}}} \right)^{-0.191 \pm 0.075}$$

$\delta_{\ln c} < 0.1$  (68.3% CL)

Umetsu+2015, in prep

Fully consistent with LCDM predictions + CLASH selection function



# Average Matter Distribution in and around CLASH Clusters

Total matter density profile @  $R=[0.01, 2]r_{\text{vir}}$   
averaged over the X-ray-selected CLASH sample:

$$\Sigma(R | M) = \int \bar{\rho} \xi_{\text{hm}}(\mathbf{r} | M) dx_{\parallel}$$

Clustering of matter  
around halos with  $M$ :

$$\xi_{\text{hm}}(r | M) = \frac{\langle \rho_{\text{halo}}(r | M) \rangle}{\bar{\rho}} + b_{\text{h}}(M) \xi_{\text{mm}}(r)$$

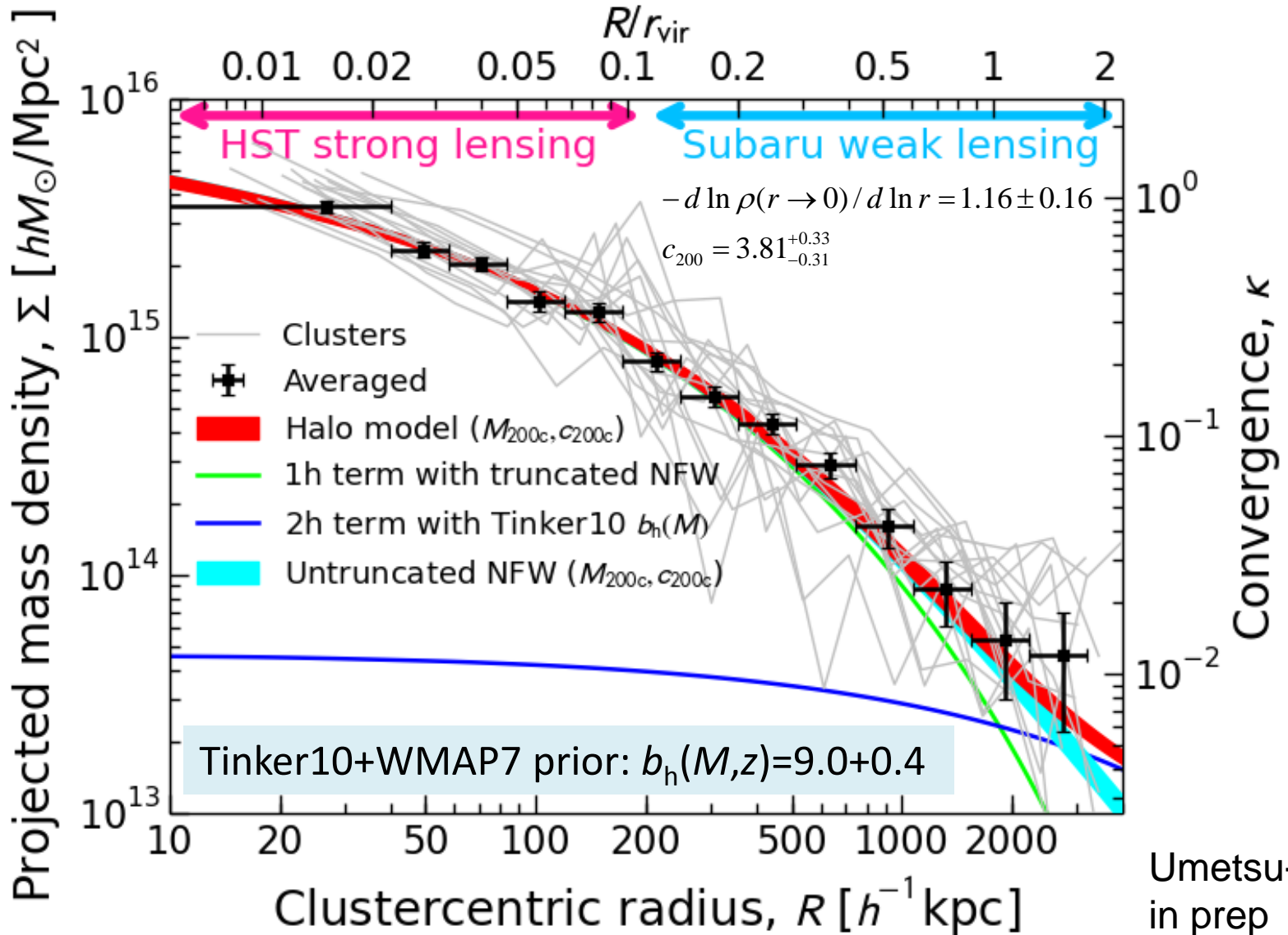
**1h term**

**2h term**





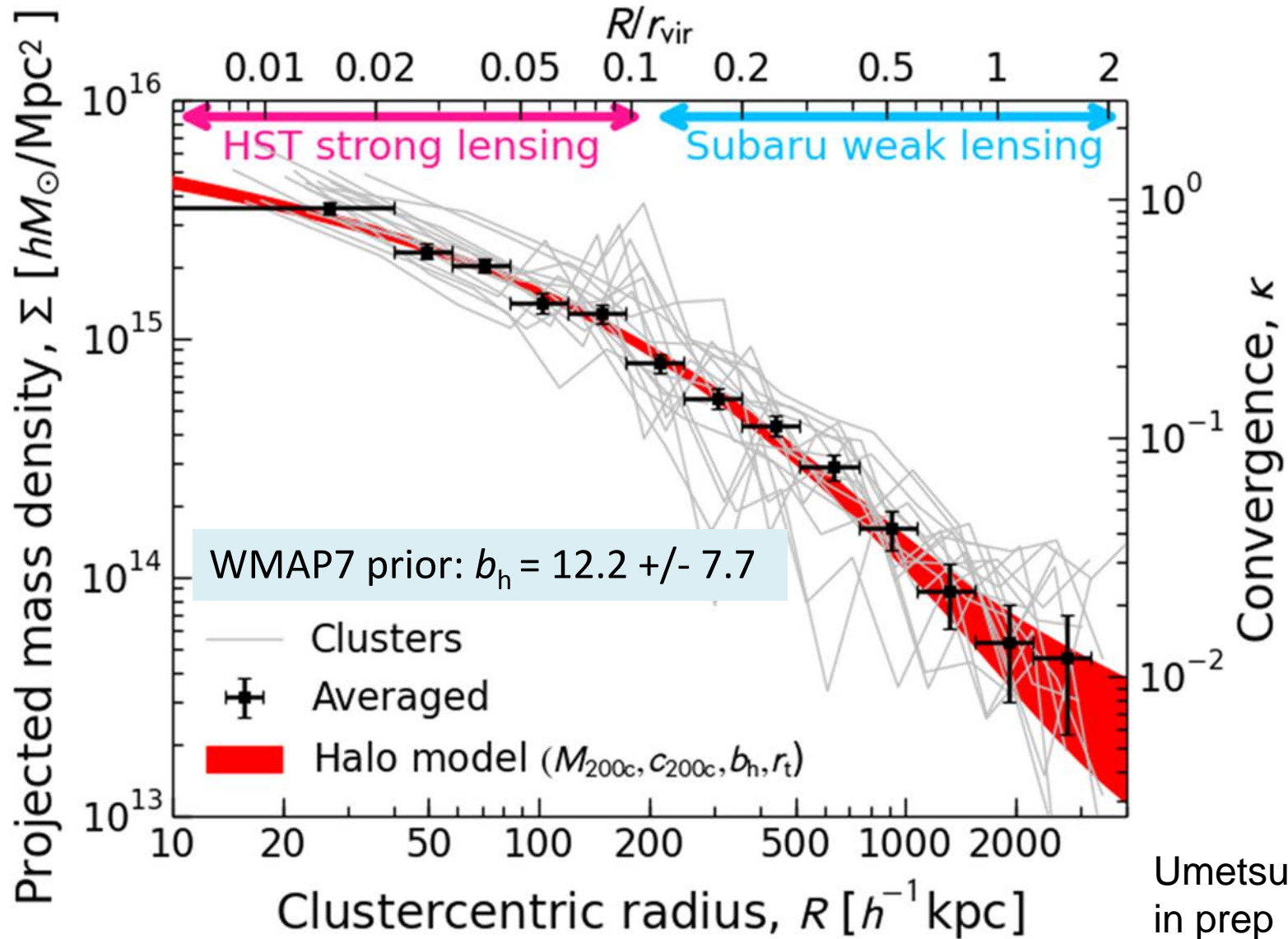
# CLASH Averaged Total Mass Profile vs. LCDM



Umetsu+2015,  
in prep



# CLASH Averaged Total Mass Profile vs. LCDM





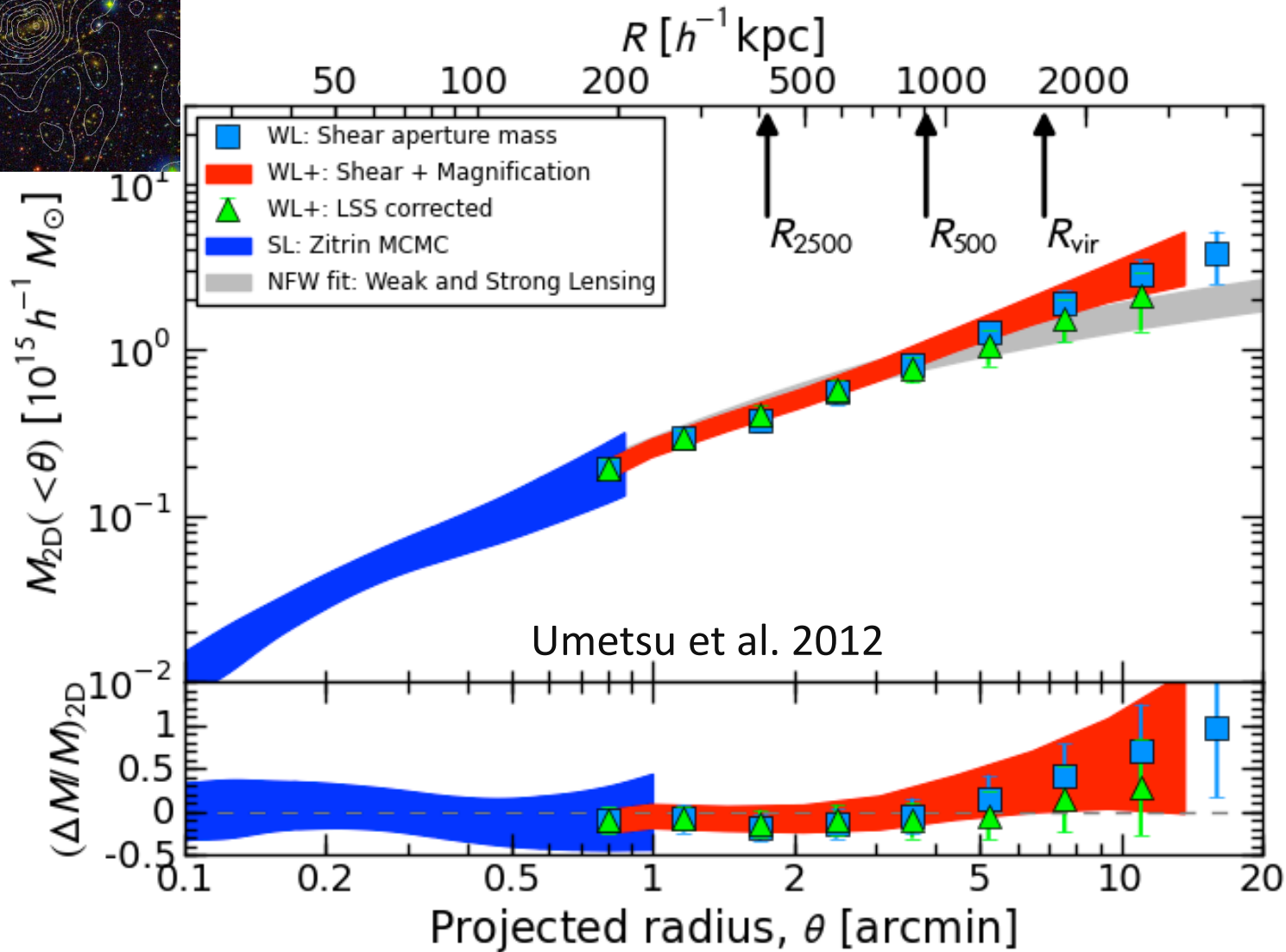
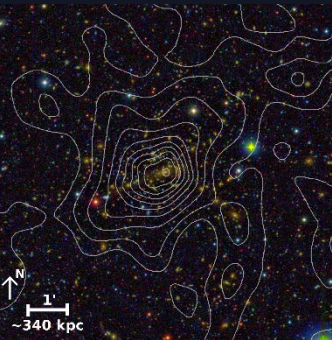
# Constraints on the Intracluster Dark-Matter Equation of State

A Case study from the ongoing CLASH-VLT  
redshift survey (PI: Piero Rosati)

# MACS1206 (z=0.44): A relaxed CLASH cluster



Total mass profiles from completely independent methods agree.

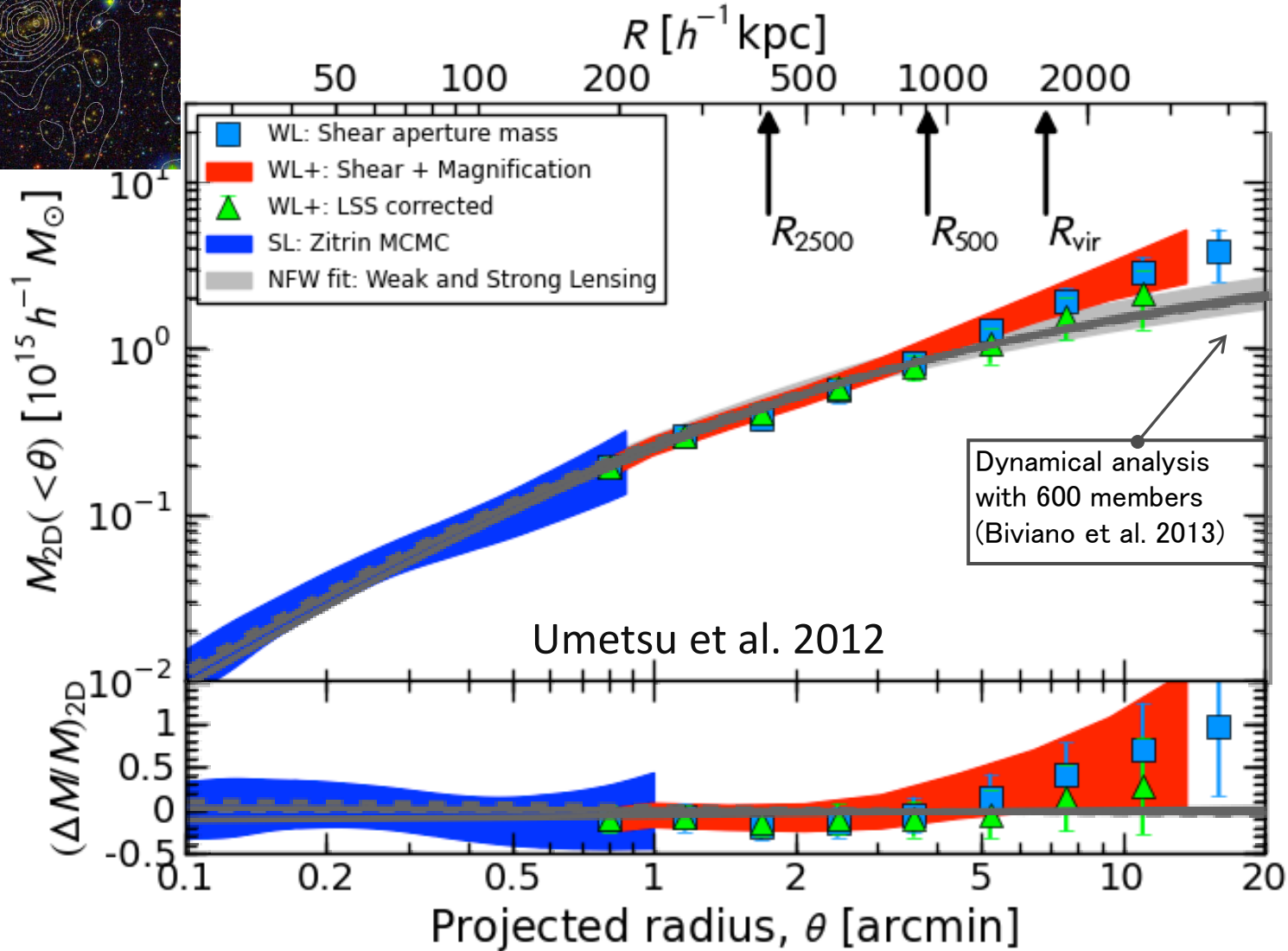
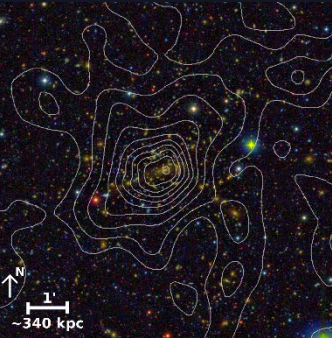




# MACS1206 (z=0.44): A relaxed CLASH cluster



Total mass profiles from completely independent methods agree.

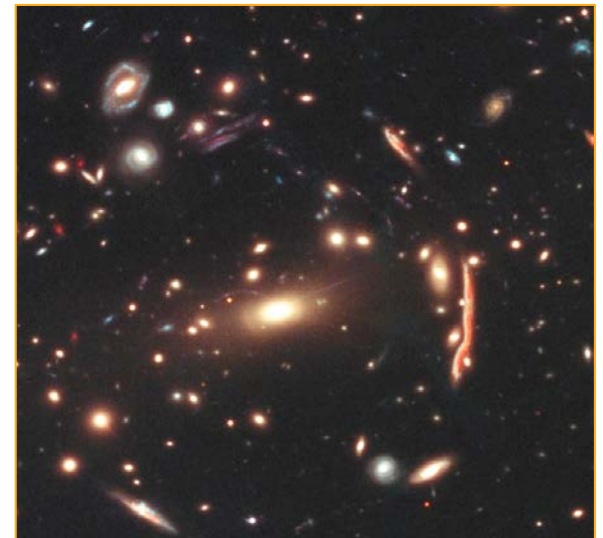




# Constraining DM Equation of State

- By testing whether intracluster DM is pressureless ( $w=0$ ) using cluster mass profiles  $M(<r)$  of MACS1206 determined from 2-independent ways:
  - Gravitational lensing with *HST+Subaru* (Umetsu+2012)
  - Galaxy kinematics with VLT/VIMOS (Biviano+2013)
- Test made possible by our high-quality CLASH data for an equilibrium cluster:

$$w(r) = \frac{p_r(r) + 2p_t(r)}{c^2 3\rho(r)}$$



# Framework

Consider the static, spherically-symmetric metric within a DM halo of the form:

$$ds^2 = -e^{-2\Phi(r)} dt^2 + \left[1 - \frac{2Gm(r)}{r}\right]^{-1} dr^2 + r^2 d\Omega^2.$$



Consider an intracluster DM fluid with anisotropic pressure. In this metric, the Einstein field equations read:

$$\begin{aligned}\rho(r) &= \frac{1}{8\pi G} \frac{m'}{r^2}, \\ p_r(r) &= -\frac{1}{8\pi G} \frac{2}{r^2} \left[ \frac{m}{r} - r\Phi' \left(1 - \frac{2m}{r}\right) \right], \\ p_t(r) &= \frac{1}{8\pi G} \left\{ \left(1 - \frac{2m}{r}\right) \left[ \frac{\Phi'}{r} + \Phi'^2 + \Phi'' \right] - \left(\frac{m}{r}\right)' \left(\frac{1}{r} + \Phi'\right) \right\}.\end{aligned}$$

The equation of state of this DM fluid is defined as

$$w(r) = \frac{p_r + 2p_t}{3\rho}.$$

Consider the weak-field limit,  $|\Phi| \ll 1$ ,  $Gm/r \ll 1$ .

# DM EoS from Kinematics+Lensing

The Jeans equation provides a way to measure the cluster mass profile from **cluster galaxy kinematics**

$$m_K(r) = -\frac{r\sigma_r^2}{G} \left[ \frac{d \ln n_g}{d \ln r} + \frac{d \ln \sigma_r^2}{d \ln r} + 2\beta \right],$$

where galaxies as the probe particles are non-relativistic,  $\sigma_r, \sigma_t \ll 1$  with  $\beta = 1 - \sigma_t^2/(2\sigma_r^2)$ . In our metric, the kinematic mass profile is related to the potential by

$$m_K(r) = \frac{r^2}{G} \Phi' \approx 4\pi \int [1 + 3w(r)] r^2 \rho(r) dr.$$

**Gravitational lensing** is sensitive to  $g_{00}$  and  $g_{rr}$ . Hence, its potential and associated mass profile are defined by

$$2\Phi_l \equiv \Phi + G \int \frac{m(r)}{r^2} dr,$$
$$m_L(r) \equiv \frac{r^2}{G} \Phi'_l = \frac{1}{2} [m_K(r) + m(r)].$$

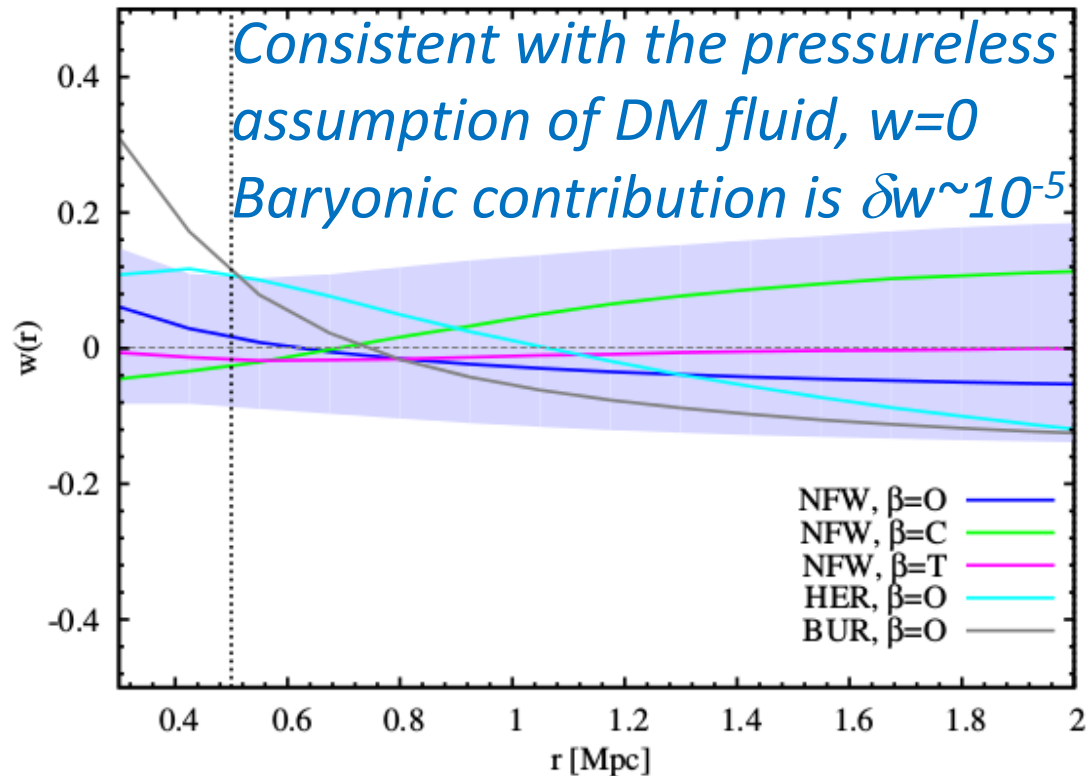
To first order, the DM equation of state is sensitive to the derivatives of the lensing and kinematic mass profiles:

$$w(r) \approx \frac{2 m'_K(r) - m'_L(r)}{3 m'_L(r) - m'_K(r)}.$$





# First application to a relaxed cluster



$$\langle w \rangle = 0.00 \pm 0.15(stat) + 0.08(syst)$$

In CLASH, we have 11 more clusters with VLT redshift measurements to improve the DM EoS constraint



# Summary

- Averaged matter distribution within CLASH clusters is in excellent agreement with standard predictions for collisionless-DM-dominated halos:
  - Outward steepening radial dependence with central cusp slope  $\beta = -d\ln\rho/d\ln r(r \rightarrow 0) = 1.16 \pm 0.16$  (NFW:  $\beta = 1$ )
  - Einasto degree of curvature,  $\alpha_E = 0.190 \pm 0.07$  ( $n_E = 1/\alpha_E \sim 5$ )
  - Average concentration,  $\langle c_{200} \rangle = 4.01 (+0.35, -0.32)$  at  $\langle M_{200} \rangle = (1.3 \pm 0.1) 10^{15} M_{\text{sun}}$ ,  $\langle z \rangle = 0.35$
  - $c$ - $M$  scaling relation with  $d\ln c(M)/d\ln M = -0.191 \pm 0.075$  and intrinsic scatter  $\delta_{\text{inc}} < 0.1$  (68.3%CL)
- Total matter distribution  $\langle \Sigma(R) \rangle$  in/around CLASH clusters  $R = [0.01, 2] r_{\text{vir}}$  is fully consistent with LCDM halo model
  - Total = smoothly-truncated NFW + correlated large-scale structure with  $b_h (\sigma_8/0.81)^2 \sim 9.0$
  - Marginal detection of clustering 2h term ( $\sim 1.6\sigma$ ) within  $2r_{\text{vir}}$



# Summary (contd.)

- Consistent WL shear & magnification measurements allow for accurate cluster mass profile measurements for 20 CLASH clusters
  - ~8% residual mass-calibration uncertainty, comparable to other current best WL efforts (~7% by *Weighing the Giants* project)
  - Crucial for cluster cosmology (cf. 20%-40% mass uncertainty in Planck 2013)
- Our lensing+kinematics study of a single cluster found the DM EoS to be  $\langle w \rangle = 0.00 \pm 0.15 \pm 0.08$  within  $R=0.5-2\text{Mpc}$ , confirming the standard pressureless assumption of DM fluid. A full CLASH-VLT sample of 12 clusters will further tighten the constraint on DM EoS.

# CLASH Products released

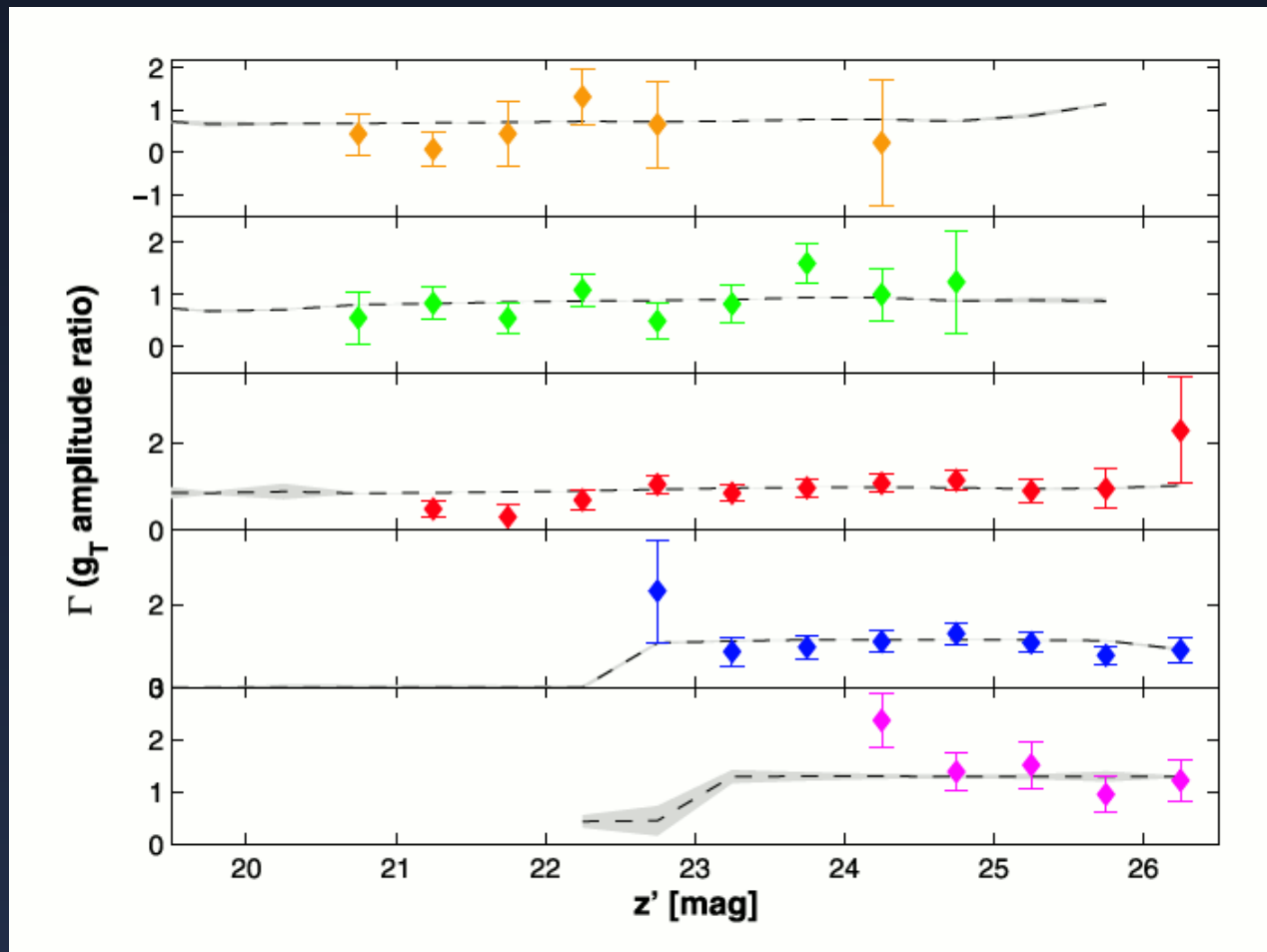
<http://archive.stsci.edu/prepds/clash/>

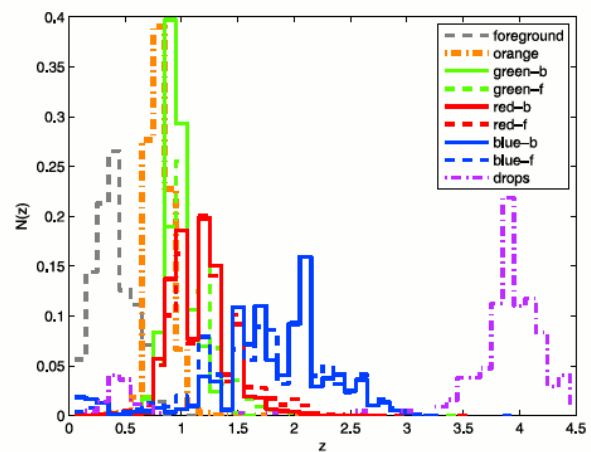
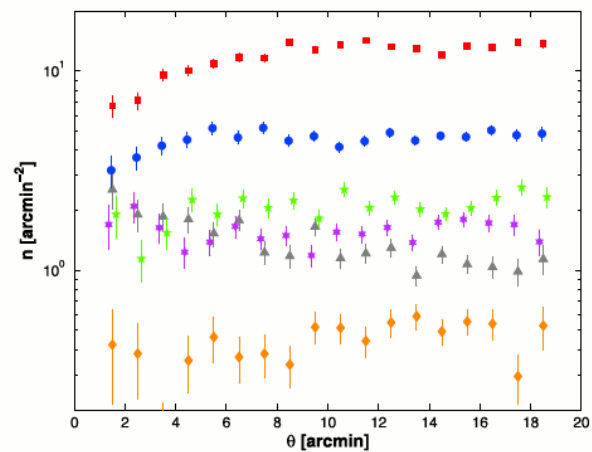
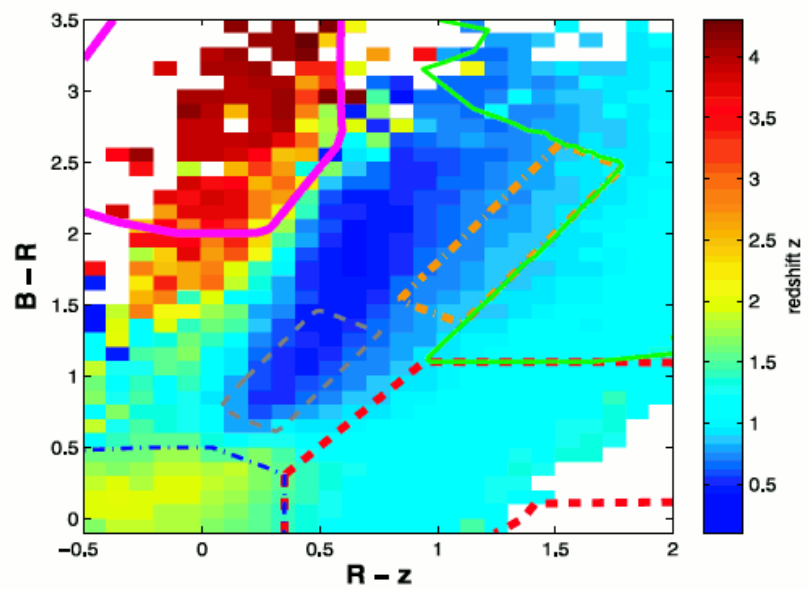
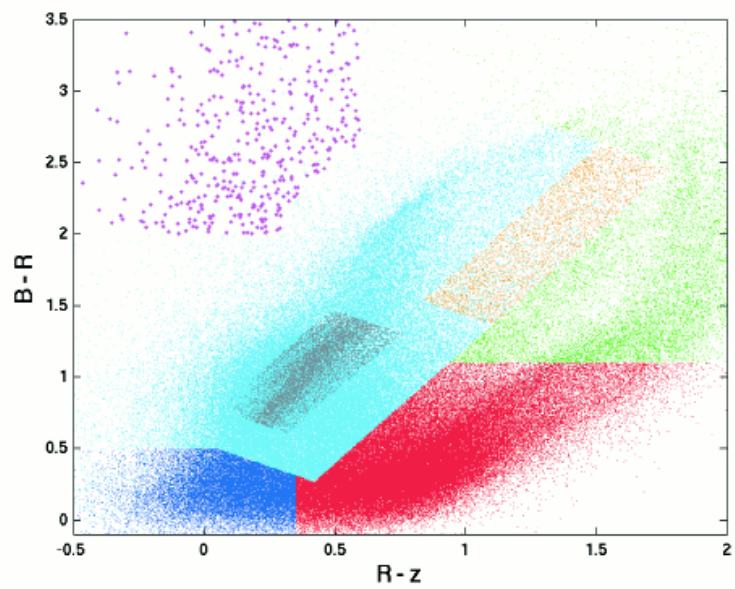
- Calibrated and co-added images [HST, Subaru]
- Object catalogs [HST, Subaru]



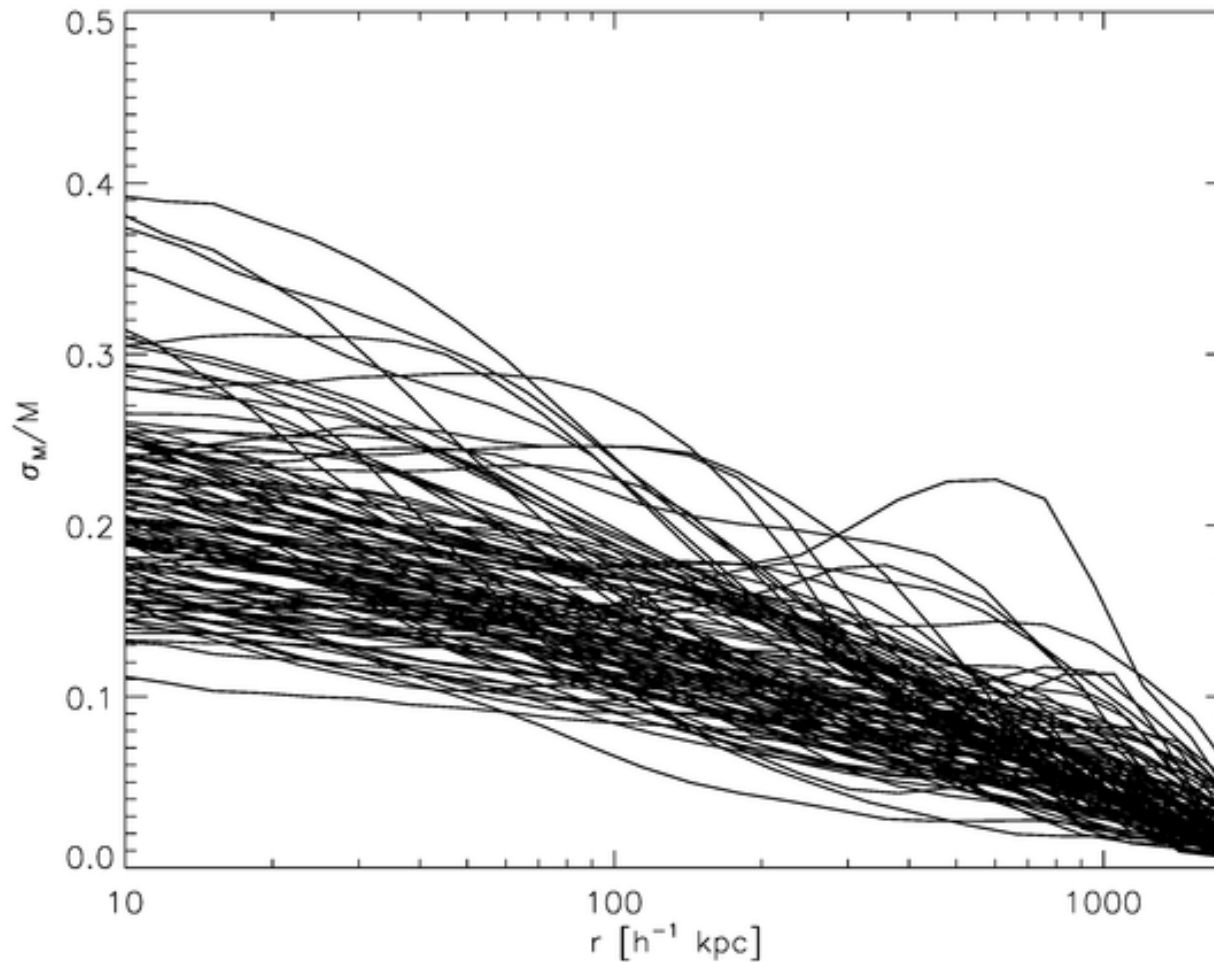
# Supplemental Slides

# SUBARU shear strength as a function of magnitude





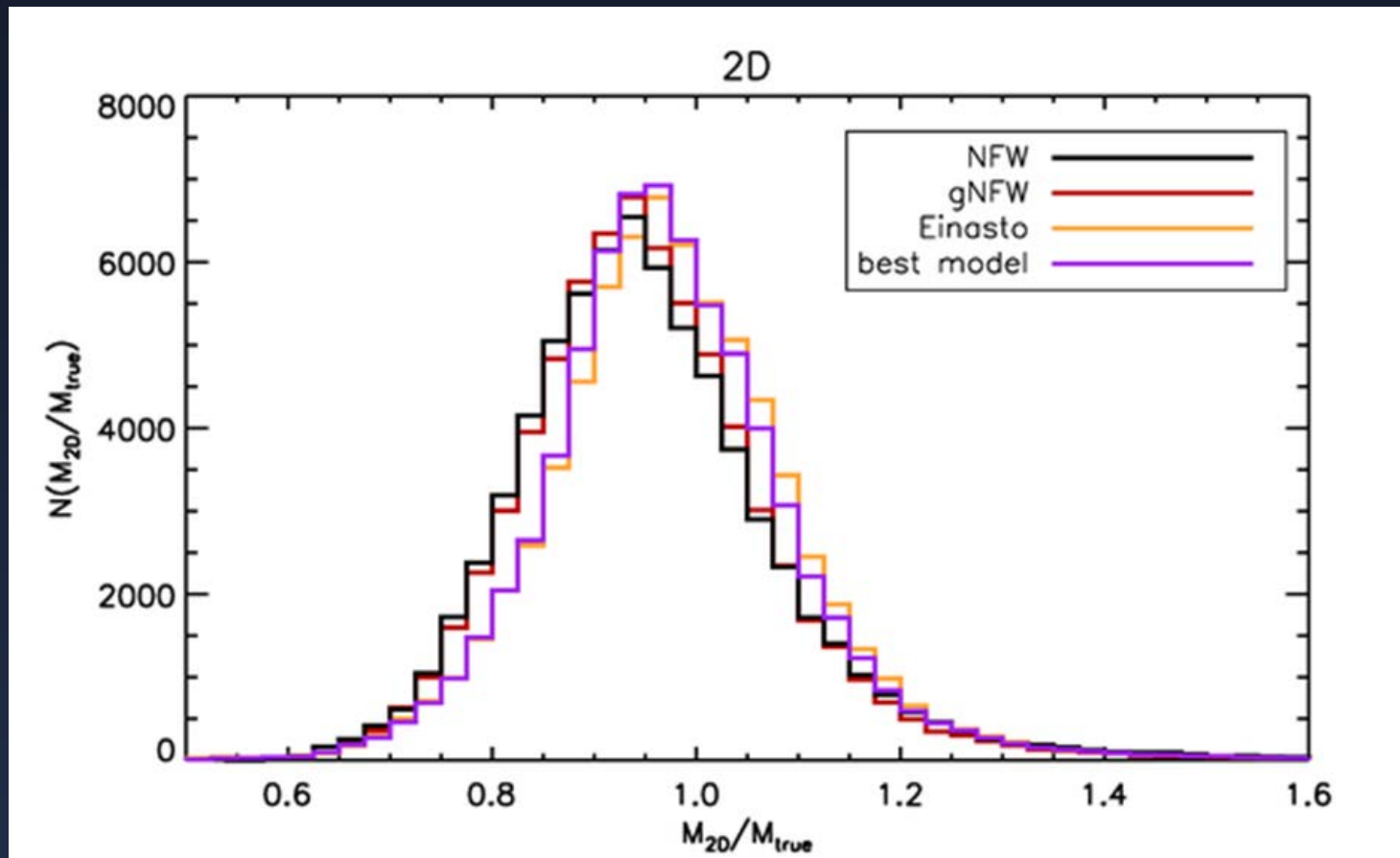
# Scatter in $M_{2D}(R)$ by halo triaxiality



MUSIC-2 simulation by Massimo



# Cluster masses recovered from lensing analysis

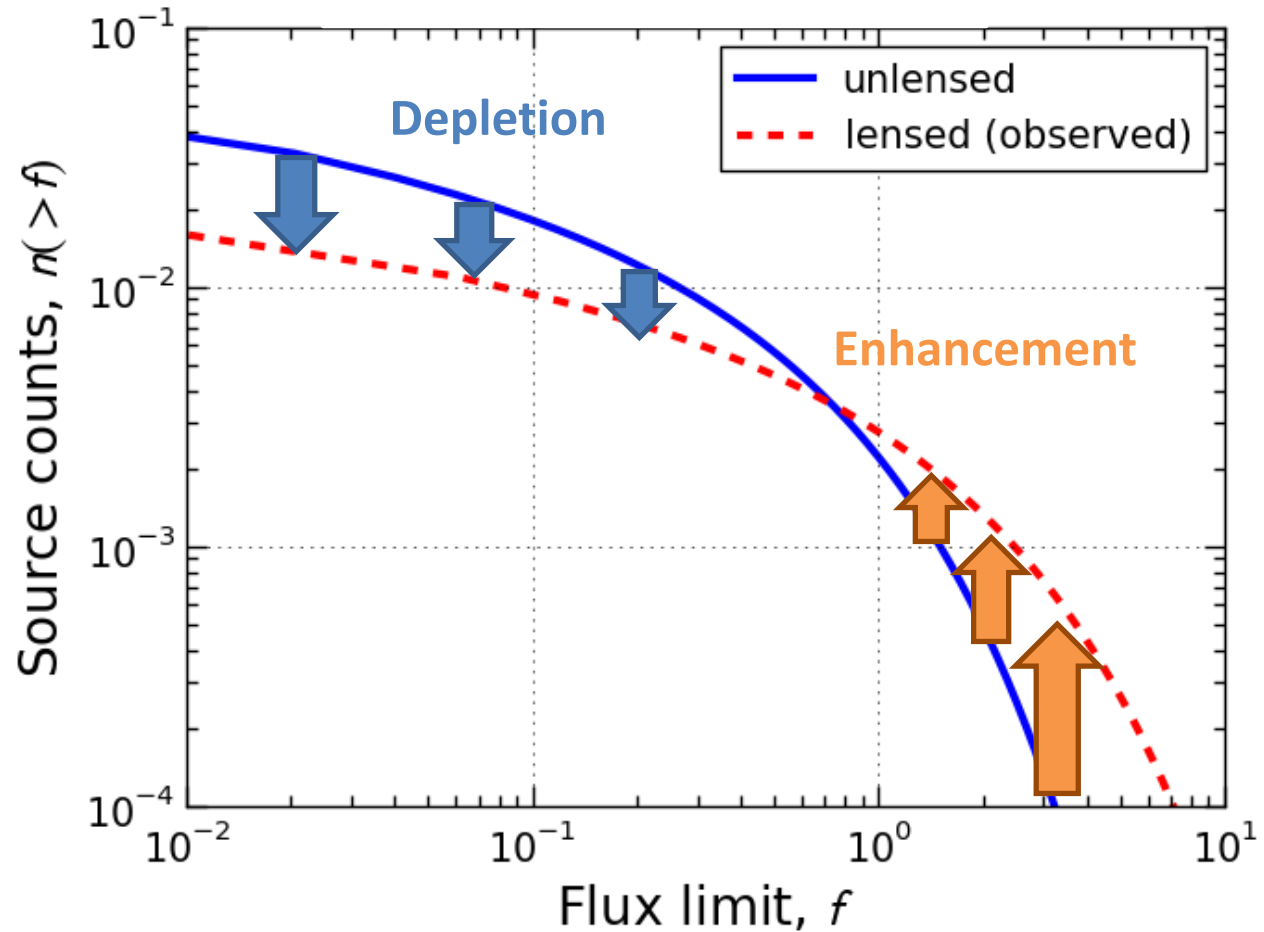


# Magnification bias effects

Flux-limited  
source counts:

$$n_{\text{obs}}(>f) = \mu^{-1} n(>\mu^{-1}f)$$

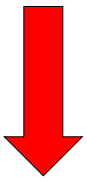
Broadhurst, Taylor &  
Peacock 95



Flux amplification



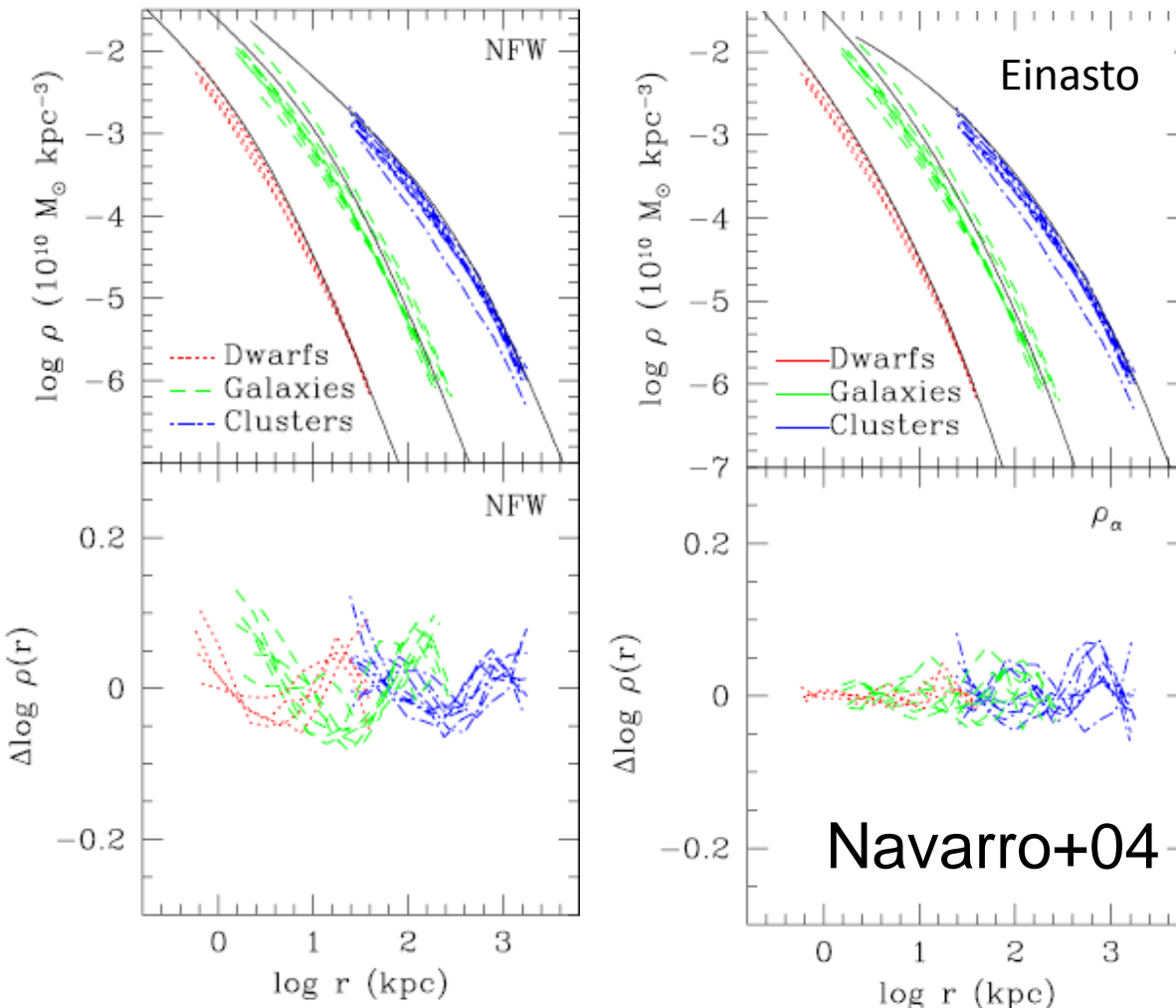
Geometric area  
distortion



$n/\mu$

# “Diversity” of halo density profiles

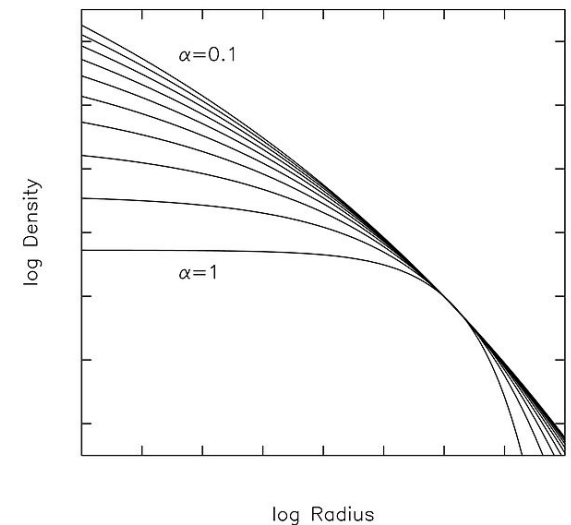
Mass profiles of DM halos are not strictly self-similar:



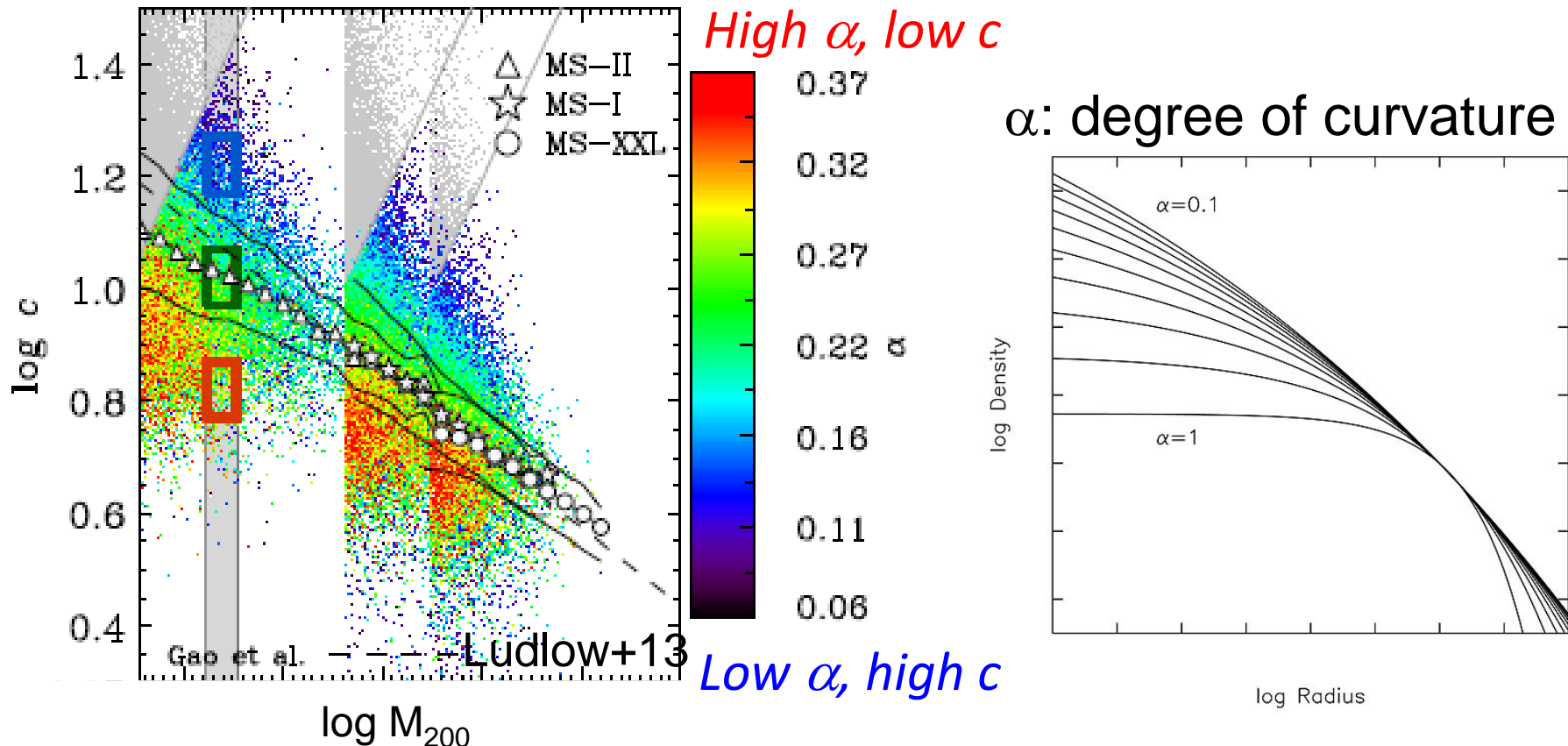
**Einasto profile ( $\rho_s, r_s, \alpha$ )**

$$\frac{d \ln \rho(r)}{d \ln r} = -2 \left( \frac{r}{r_s} \right)^{\alpha}$$

$\alpha$ : degree of curvature



# Intrinsic Scatter in $c(M)$ : Mass Assembly Histories (MAH)



*High  $\alpha$ , low  $c$*

*Low  $\alpha$ , high  $c$*

$\alpha$ : degree of curvature

- Scatter is due to another DoF ( $\alpha$ ), related to MAH (Ludlow+13)
- Larger or smaller values of  $\alpha$  correspond to halos that have been assembled more or less rapidly than the NFW curve
- Clusters with average  $c_{200}$  have the NFW-equivalent  $\alpha \sim 0.18$



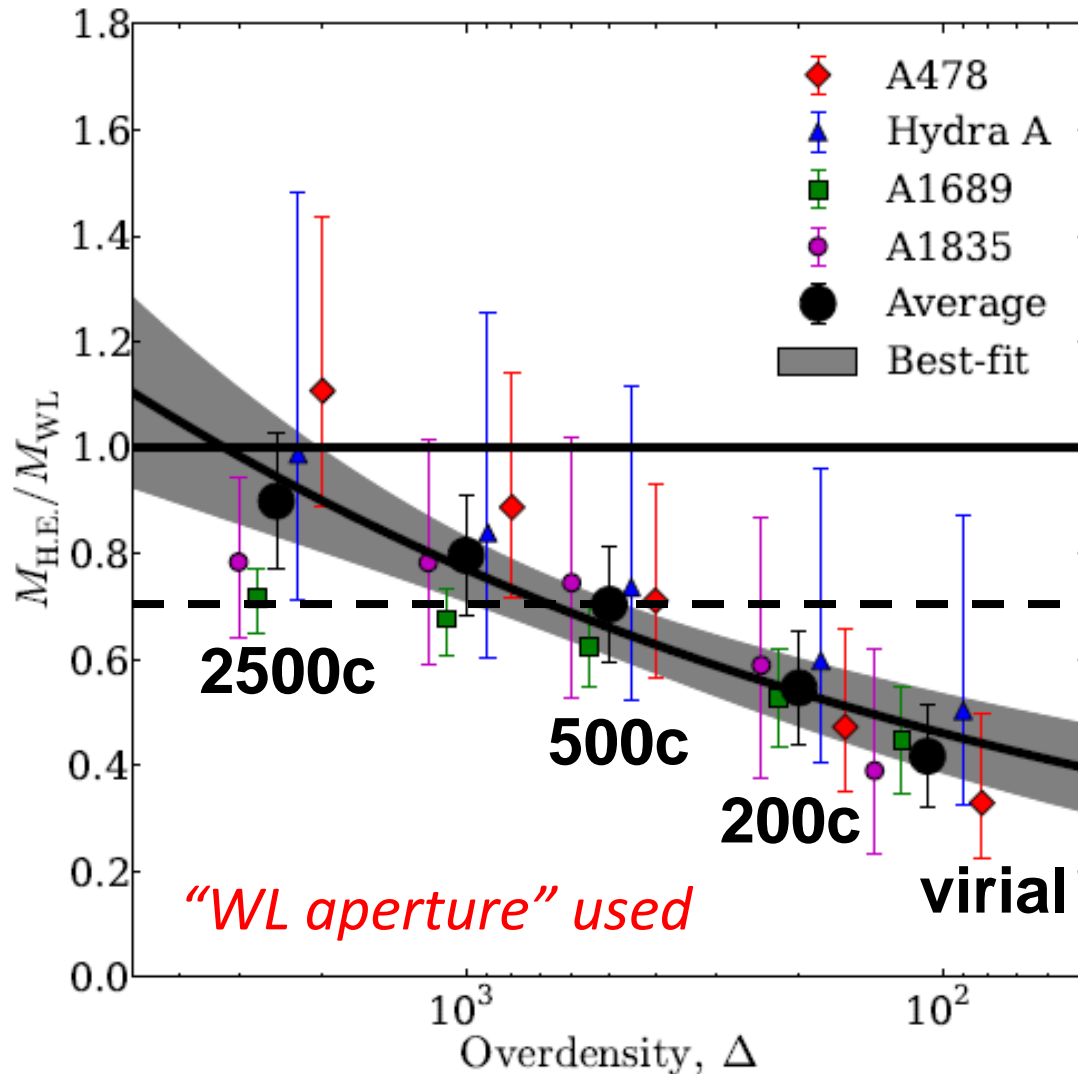
$$\begin{aligned}
ds^2 &= a^2(\eta)d\tilde{s}^2 = a^2\tilde{g}_{\mu\nu}dx^\mu dx^\nu \\
&= a^2 \left[ -(1+2\Psi)d\eta^2 + (1-2\Psi) \left\{ d\chi^2 + r^2(\chi)(d\theta^2 + \sin^2\theta d\phi^2) \right\} \right]
\end{aligned}$$

$$\delta k^\mu(\lambda) = -\frac{2}{r^2(\lambda)} \int_0^{\lambda_s} d\lambda' \partial^\mu \Psi(\lambda') / c^2 \quad (\mu = \theta, \phi).$$

$$\beta - \theta = \int_{\text{Observer}}^{\text{Source}} d\alpha = \alpha(\chi_s),$$

$$\alpha(\chi_s) = -\frac{2}{c^2} \int_0^{\lambda_s} d\lambda \frac{r(\lambda_s - \lambda)}{r(\lambda_s)} \nabla_{\perp} \Psi(x(\lambda)); \quad x(\lambda) = x^{(b)}(\lambda) + \delta x(\lambda)$$

# Suzaku-X HSE vs. Subaru WL

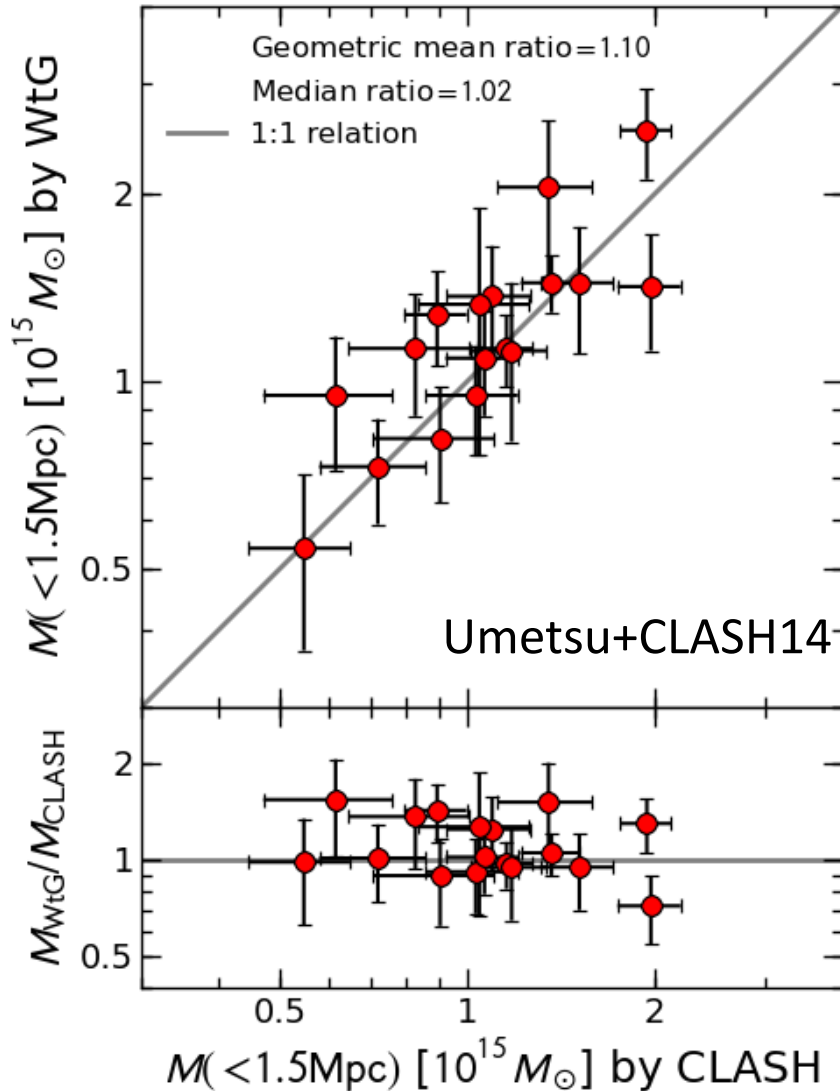


Independent *Suzaku*-HSE vs. Subaru-WL results, consistent with XMM-HSE vs. Subaru-WL of CLASH collaboration

Okabe, Umetsu et al.14, PASJ, in press (arXiv:1406.3451)



# Comparison with WtG @R=1.5Mpc



17 clusters in common (Subaru):

- **WtG**: shear-only (Applegate+14), NFW  $c_{200c}=4$  prior
- **CLASH**: shear + magnification, NFW log-uniform:  $0.1 < c_{200c} < 10$

**Un-weighted geometric mean mass ratio ( $\langle Y/X \rangle = 1/\langle X/Y \rangle$ )**

- $\langle M_{\text{WtG}}/M_{\text{CLASH}} \rangle = 1.10$
- Median ratio = 1.02

Systematic uncertainty in the overall mass calibration of 8% from shear-magnification consistency (Umetsu+14)

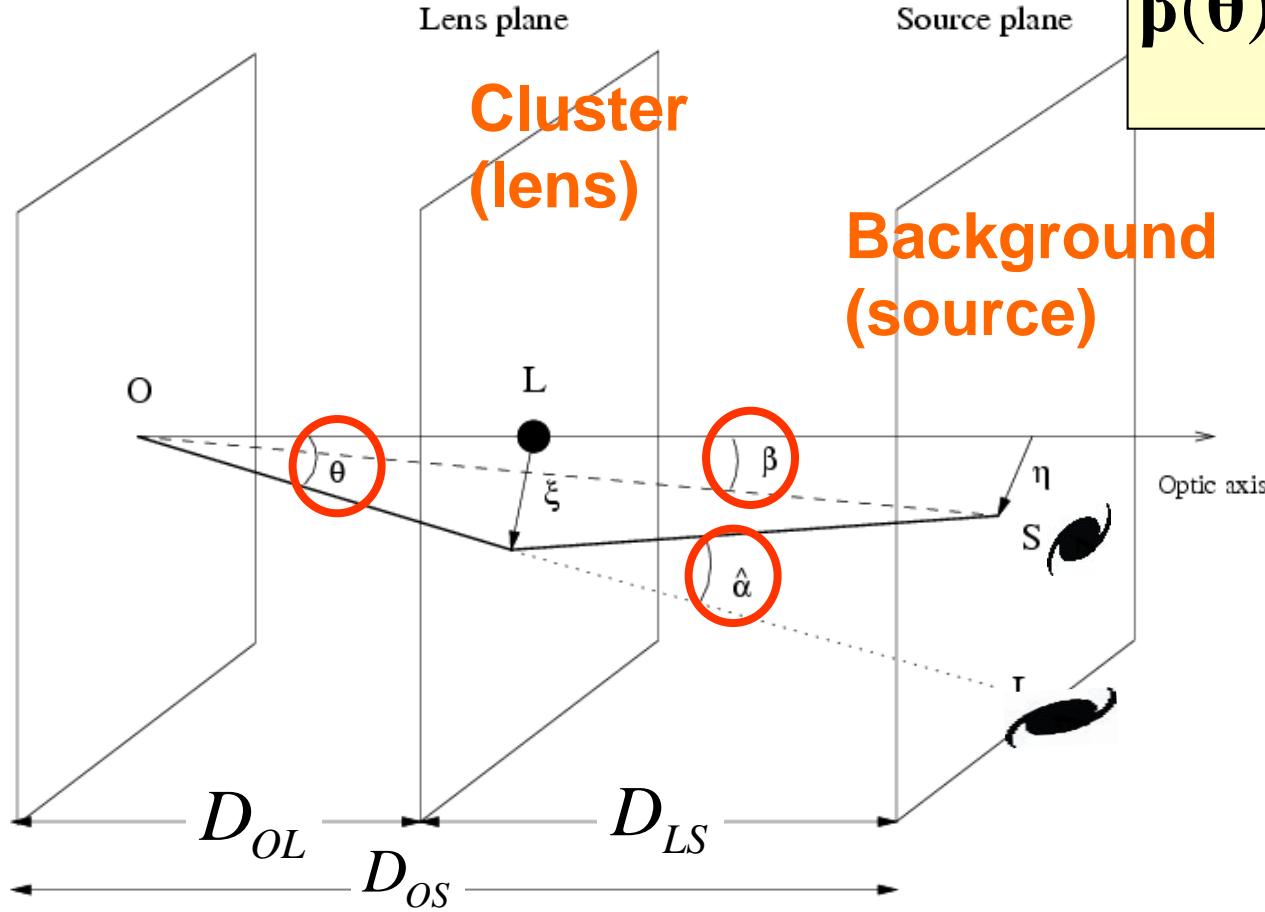
No mass dependent bias

# Cluster Lens Equation

## Cosmological lens equation + single/thin-lens approximations

$\beta$ : true (but unknown) source position

$\theta$ : apparent image position



$$\beta(\theta) - \theta = \frac{D_{LS}}{D_{OS}} \int \delta \hat{\alpha}(\theta)$$

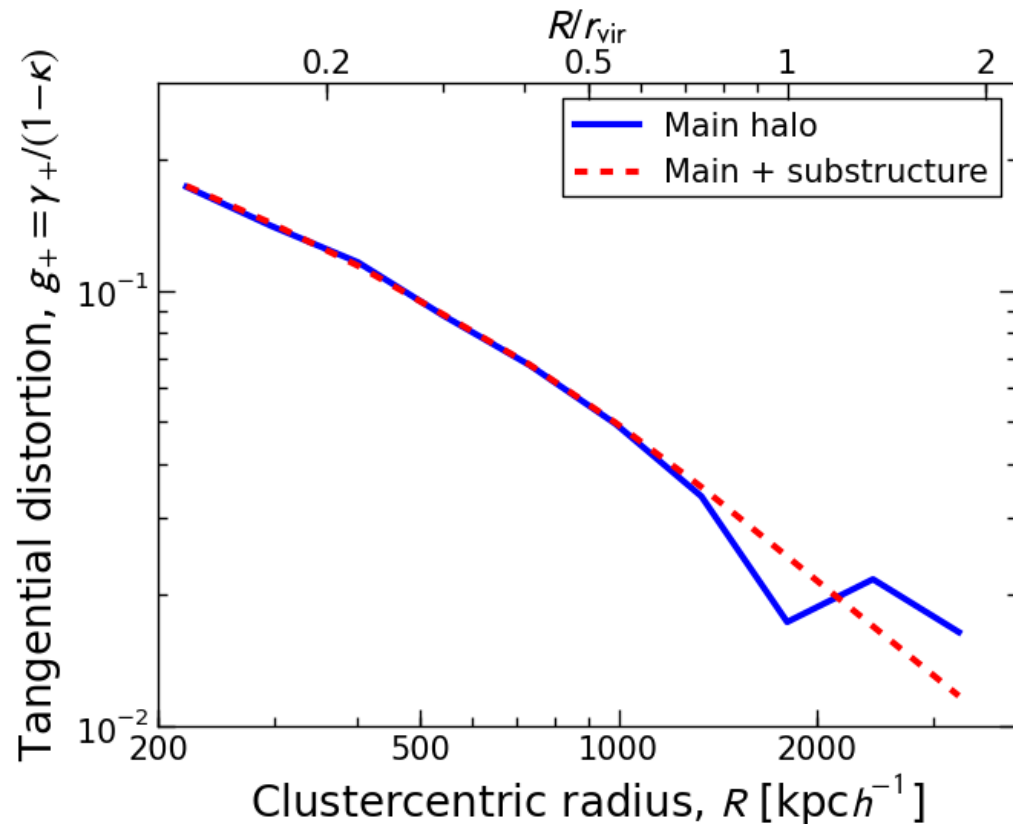
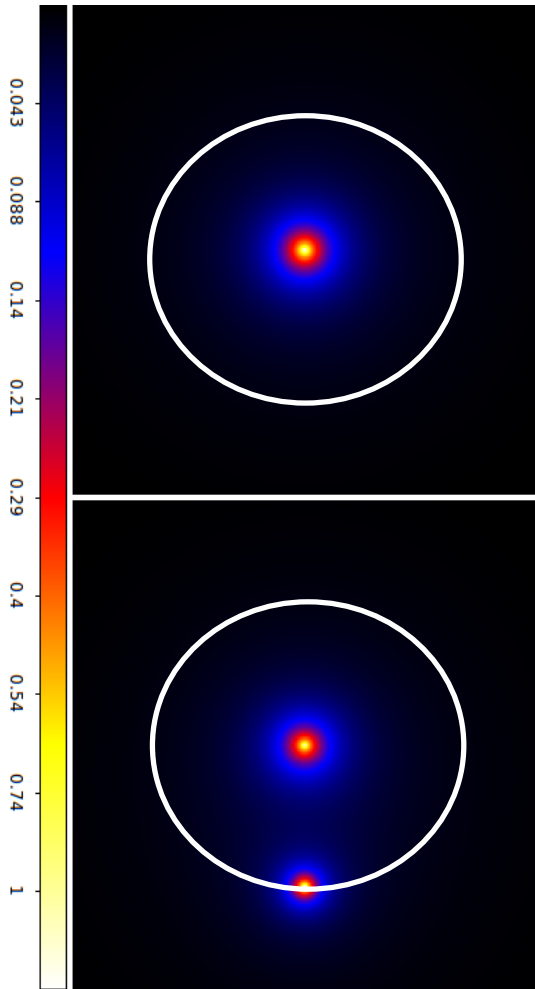
Angular diameter distances:

$$D_{OL}, D_{LS}, D_{OS} \sim O(c/H_0)$$

For a rigid derivation of cosmological lens eq., see, e.g., Futamase 95

# Non-local substructure effect

A substructure at  $R \sim r_{\text{vir}}$  of the main halo, modulating  $\Delta\Sigma(R) = \Sigma(< R) - \Sigma(R)$



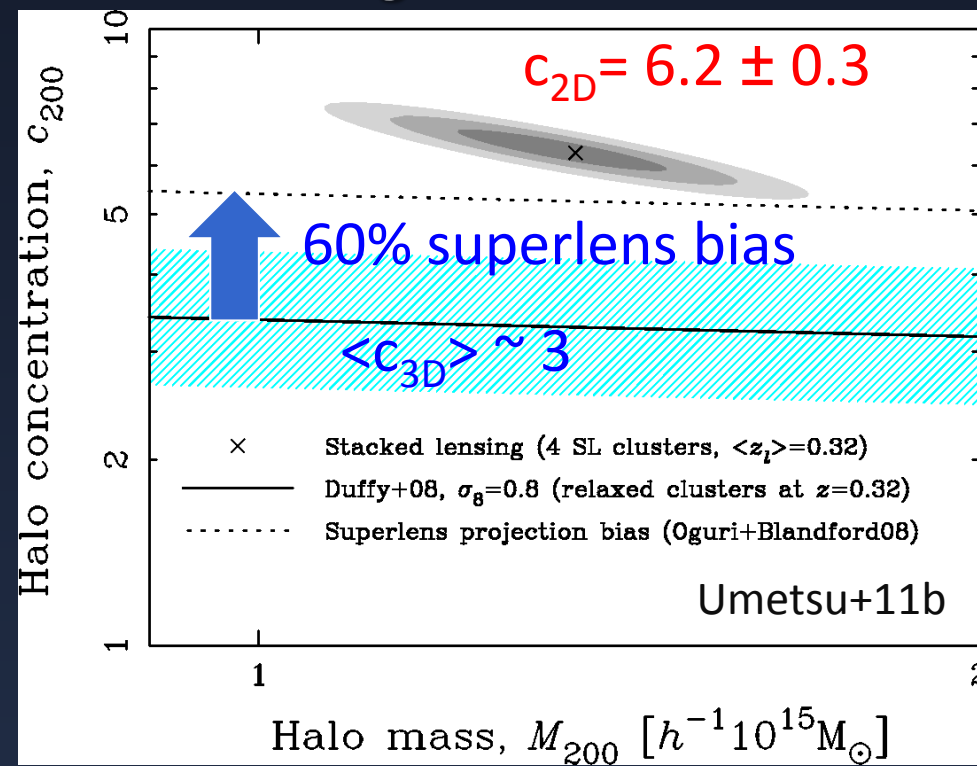
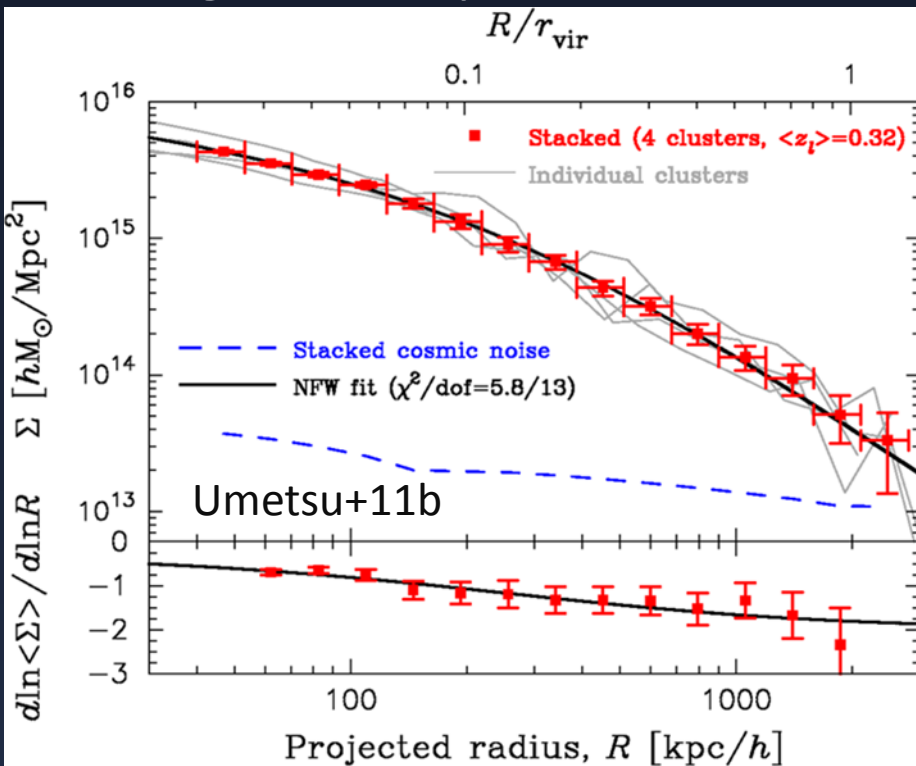
Known  $\sim 10\%$  negative bias in mass estimates from tangential-shear fitting, inherent to clusters sitting in substructured field (Rasia+12)



# CLASH Objectives & Motivation

$$c_{200c} := \frac{r_{200c}}{r_s}$$

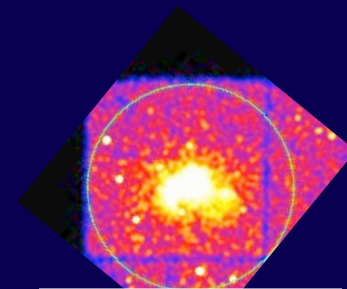
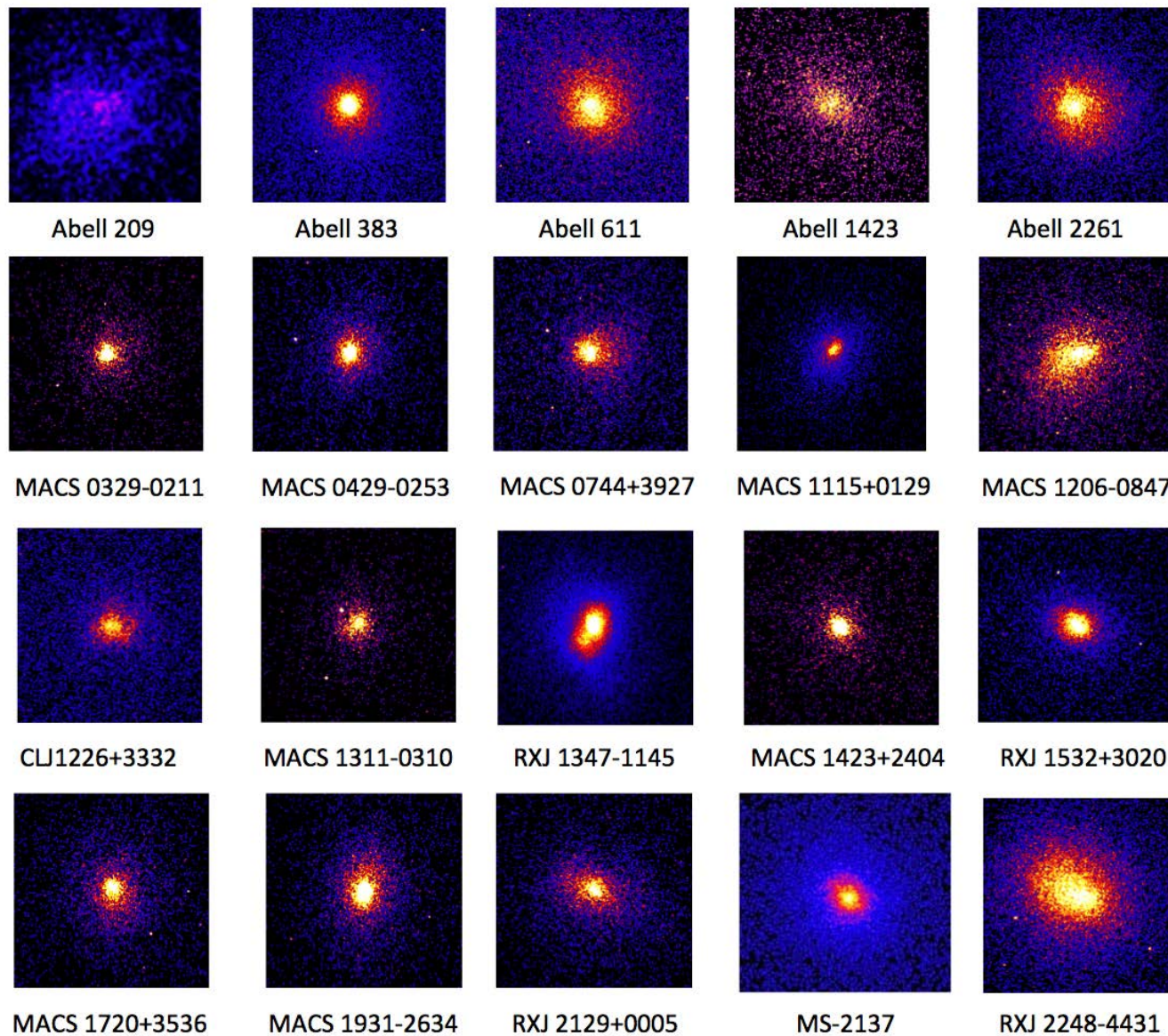
Before CLASH (2010), deep-multicolor Strong (*HST*) + Weak (Subaru) lensing data only available for a handful of **strong-lens clusters**



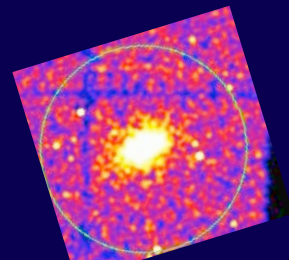
**Total mass profile shape:** consistent w CDM (self-similar universal profile)

**Degree of concentration:** maximum superlens correction not enough if  $\langle c_{\text{LCDM}} \rangle \sim 3$ ?

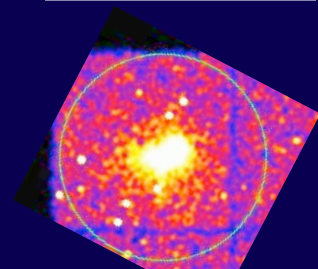
# X-ray observations with Chandra and XMM-Newton Satellites



MACS 0717+3745



RXJ 0647+7015



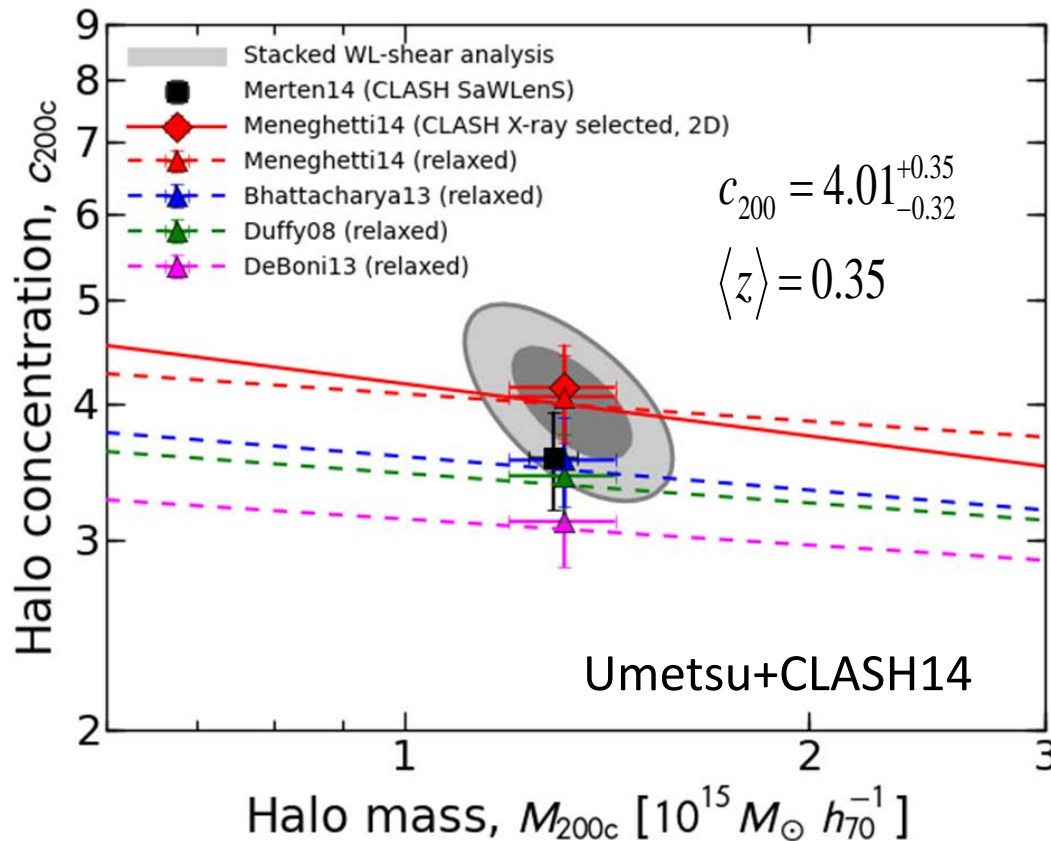
MACS 1149+2223

All have  
 $T_x > 5$  keV

X-ray images of 23 of the 25 CLASH clusters. 20 are selected to be “relaxed” clusters (based on their x-ray properties only). 5 are selected specifically because they are strongly lensing  $\theta_E > 30''$



# CLASH-WL vs. c-M relations



M14 (CLASH):  $\sigma_8 = 0.82$

Bhat13:  $\sigma_8 = 0.8$

Duffy08:  $\sigma_8 = 0.8$

DeBoni13:  $\sigma_8 = 0.78$

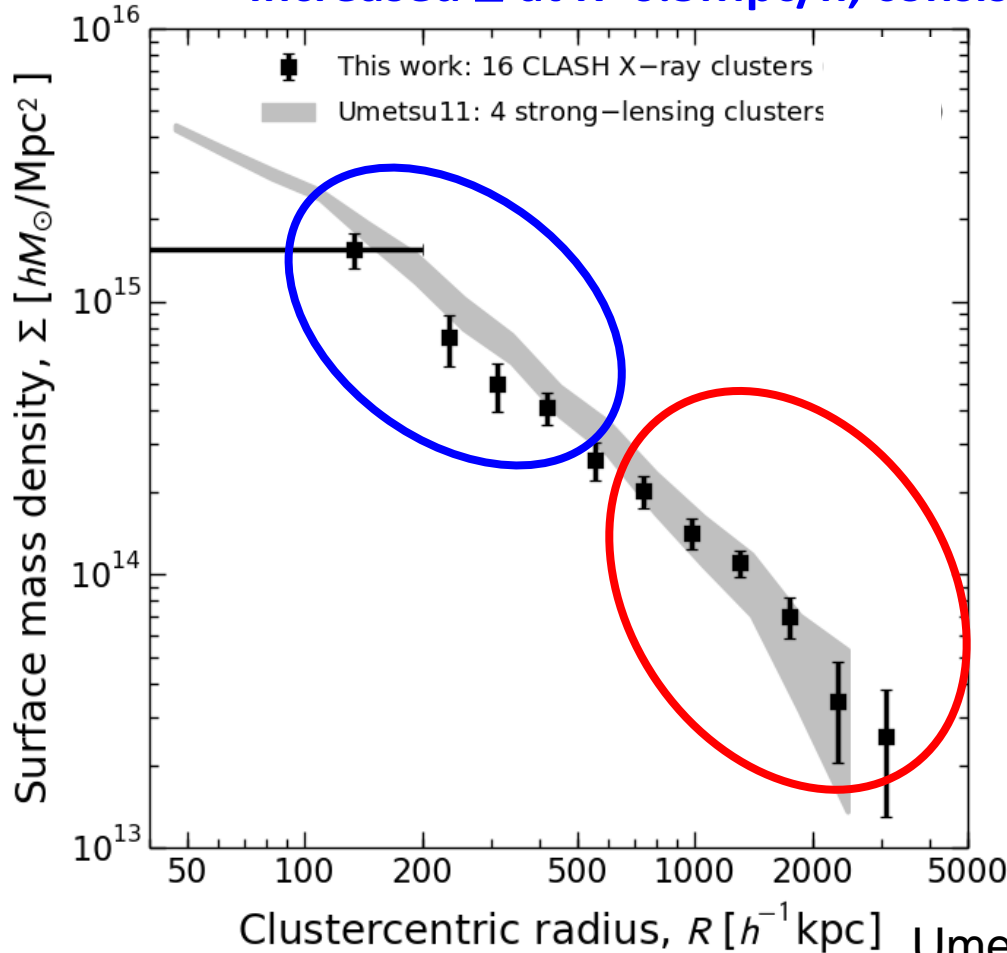
At low  $M_{200c}$ , X-ray selection picks up clusters with higher concentrations (Meneghetti+14)





# Comparison with pre-CLASH results

- $C_{200}$  vs  $\theta_E$  relation, consistent with triaxial CDM halos (Oguri+12)
- **Similar  $v$  (MAH), similar  $\Sigma$  in outskirts (Diemer & Kravtsov 14)**
- **Increased  $\Sigma$  at  $R < 0.5 \text{ Mpc}/h$ , consistent w orientation bias (Gao+12)**



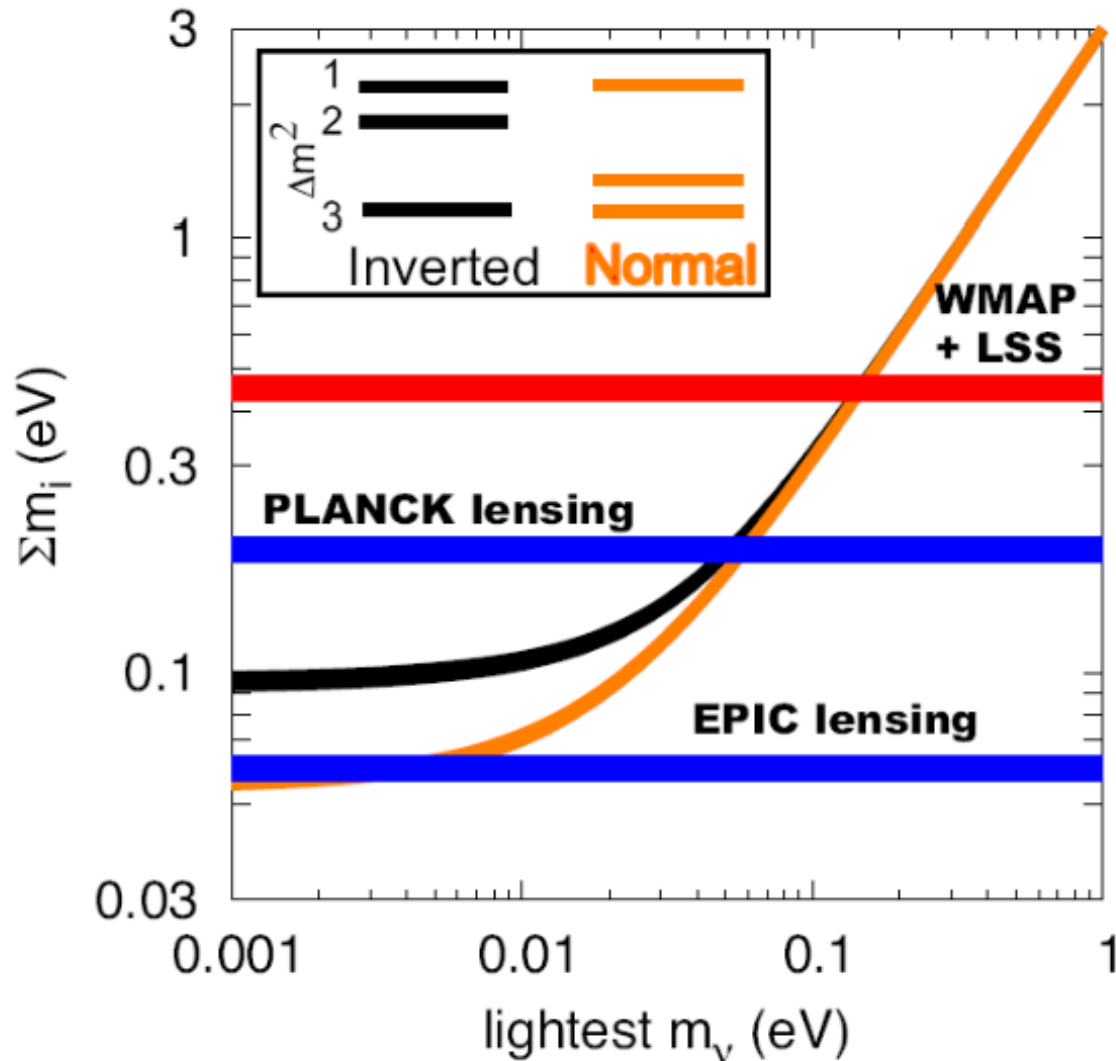
## CLASH X-ray-selected sample

- $M_{200} = 1.3e15 M_{\text{sun}}$
- $\underline{C_{200} = 4.0}$
- $\underline{\theta_E \sim 15'' (z_s=2)}$
- $\underline{v=3.8 (b_h \sim 9)}$

## Umetsu11b sample

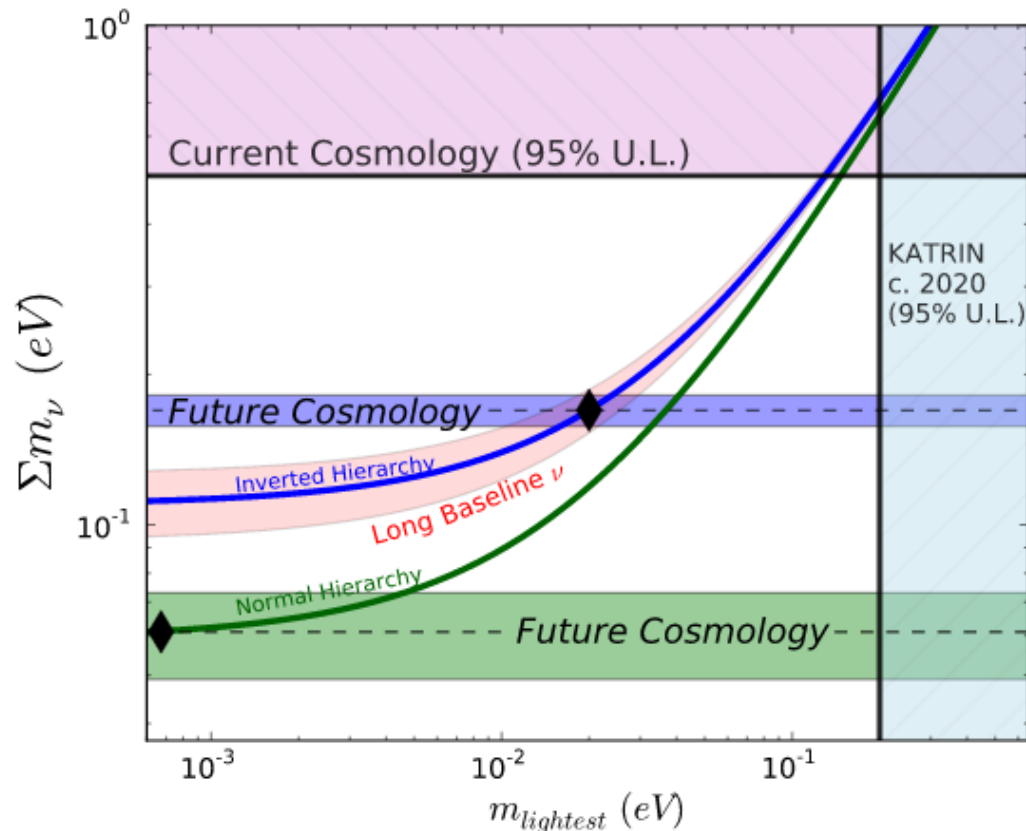
- $M_{200} = 1.7e15 M_{\text{sun}}$
- $\underline{C_{200} = 6.1}$
- $\underline{\theta_E \sim 36'' (z_s=2)}$
- $\underline{v=4.1 (b_h \sim 11)}$

# Neutrino Mass Hierarchy from Cosmology





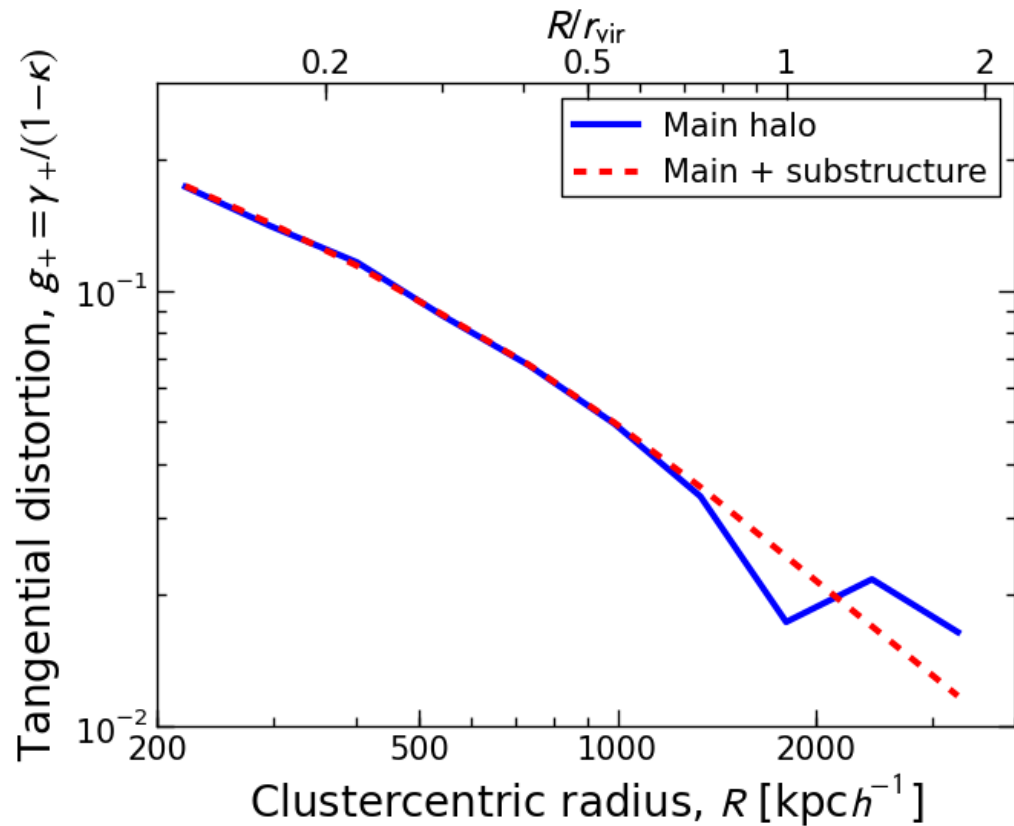
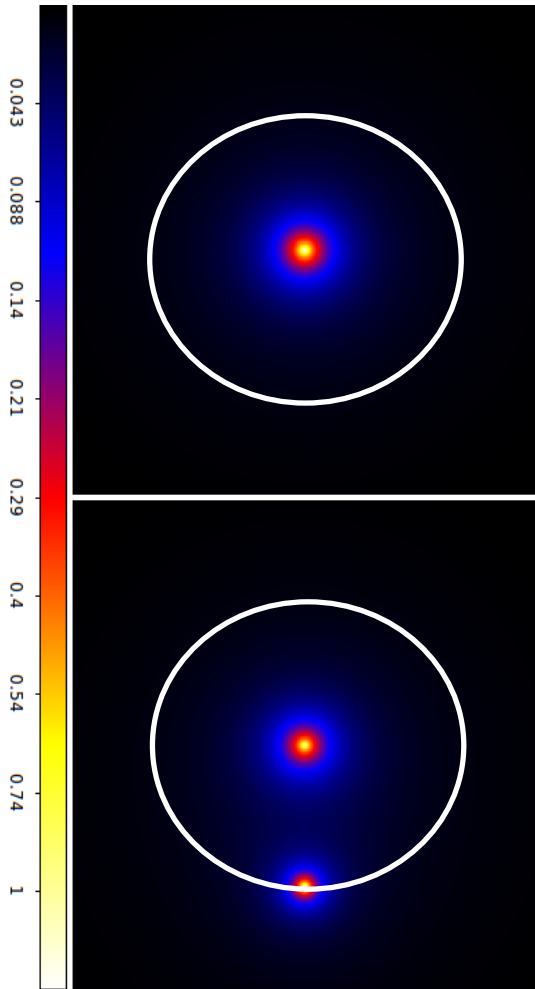
# Future Cosmological Constraints on Neutrino Hierarchy



**Figure 1-2.** Shown are the current constraints and forecast sensitivity of cosmology to the neutrino mass in relation to the neutrino mass hierarchy. In the case of an “inverted hierarchy,” with an example case marked as a diamond in the upper curve, future combined cosmological constraints would have a very high-significance detection, with  $1\sigma$  error shown as a grey band. In the case of a normal neutrino mass hierarchy with an example case marked as diamond on the lower curve, future cosmology would detect the lowest  $\Sigma m_\nu$  at a level of  $> 4\sigma$ .

# Shear: non-local substructure effect

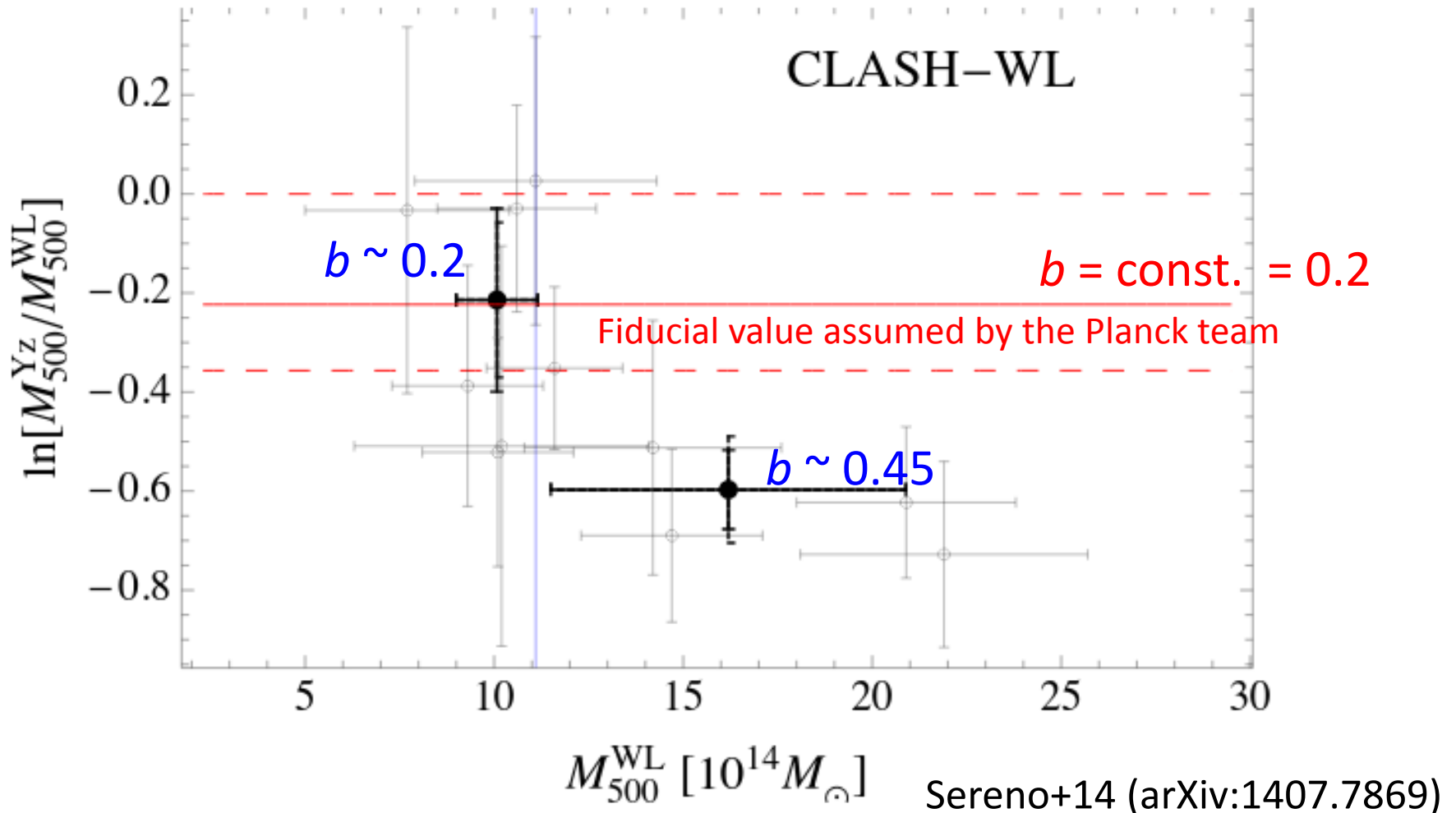
Substructure in outskirts of the main halo, modulating  $\Delta\Sigma(R) = \Sigma(< R) - \Sigma(R)$



Known 5-10% negative bias in mass estimates from tangential-shear fitting, inherent to rich clusters (Rasia+12)

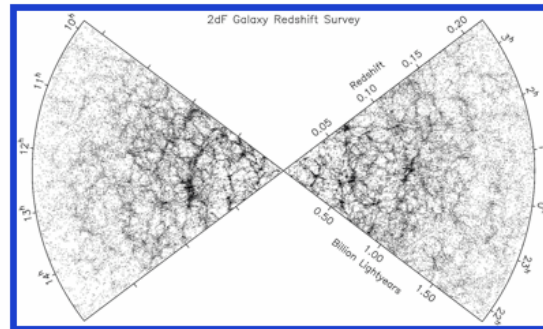
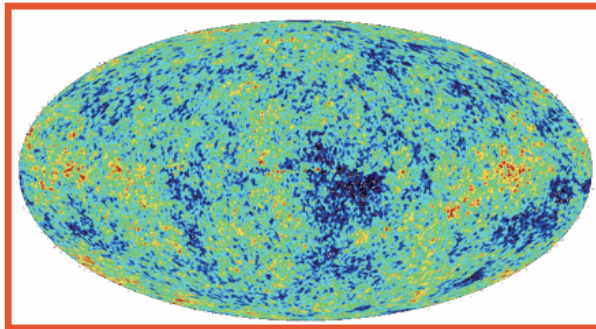
# Comparison with *Planck* Masses

Mass-dependent bias (20-45%) observed for *Planck* mass estimates



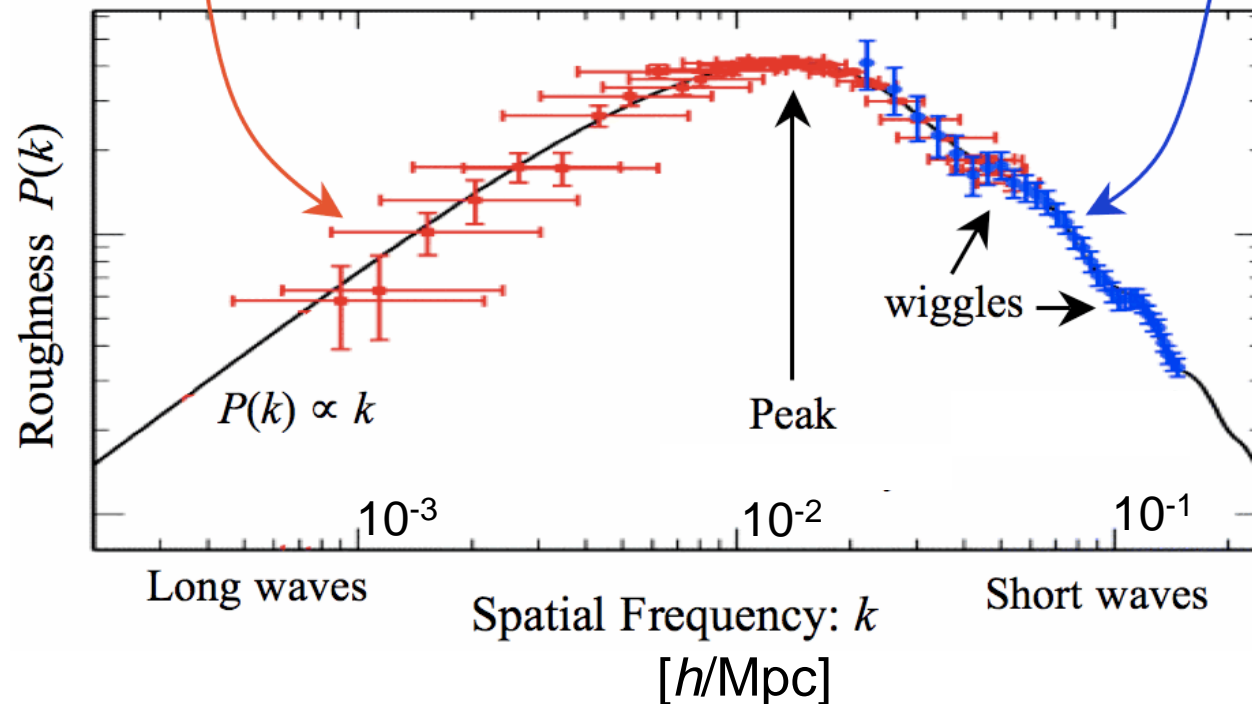
# $\Lambda$ CDM: Standard Structure Formation Paradigm

## Matter power-spectrum density, $P(k)$



$$\delta(\mathbf{x}) := \frac{\rho - \bar{\rho}}{\bar{\rho}} = \int \frac{d^3k}{(2\pi)^3} \tilde{\delta}(\mathbf{k}) e^{i\mathbf{k}\cdot\mathbf{x}}$$

$$\langle \tilde{\delta}(\mathbf{k}) \tilde{\delta}(\mathbf{k}') \rangle = (2\pi)^3 \delta_D^3(\mathbf{k} + \mathbf{k}') P(k)$$



**How about nonlinear small scales,  $\lambda < 10 \text{ Mpc}/h$ ?**