







• EM electrodynamic potential

Lorentz gauge

$$\begin{array}{rcl} A' &=& A + \nabla \Lambda \\ \Phi' &=& \Phi - \frac{1}{c} \dot{\Lambda} \end{array} \Rightarrow \nabla \cdot A + \frac{\epsilon \mu}{c} \dot{\Phi} = 0 \end{array}$$

$$\nabla^2 A - \frac{\epsilon \mu}{c^2} \ddot{A} = -\frac{4\pi}{c} \mu J$$
$$\nabla^2 \Phi - \frac{\epsilon \mu}{c^2} \ddot{\Phi} = -\frac{4\pi}{\epsilon} \rho$$

• Green's function and retarded potentials

$$\nabla^2 \psi - \frac{1}{v^2} \ddot{\psi} = -f(x,t)$$

$$\psi(x,t) = \int_{-\infty}^{\infty} \Psi(x,\omega) e^{i\omega t} d\omega$$

$$f(x,t) = \int_{-\infty}^{\infty} F(x,\omega) e^{i\omega t} d\omega$$

$$\Rightarrow (\nabla^2 + k^2) \Psi(x,\omega) = -F(x,\omega)$$

$$(\nabla^2 + k^2)G(x, x') = -\delta(x - x') \Rightarrow \qquad \Psi(x, \omega) = \int F(x', \omega)G(x - x')d^3x'$$
$$F(x, \omega) = \int F(x', \omega)\delta(x - x')d^3x'$$

$$g(x,t,x',t') = \frac{1}{2\pi} \int G(x,x') e^{i\omega t} d\omega \Rightarrow (\nabla^2 - \frac{1}{\upsilon^2} \frac{\partial^2}{\partial t^2}) g(x,x',t,t') = -\delta(x-x')\delta(t-t')$$

spherical coord.
$$G(x,t) = \frac{1}{4\pi r} e^{\pm ikr} \Rightarrow g(x,t,x',t') = \frac{\delta(t' + \frac{|x-x'|}{v} - t)}{4\pi |x-x'|}$$

• Green's function and retarded potentials

$$\psi(x,t) = \frac{1}{4\pi} \int \frac{f(x',t - \frac{|x-x'|}{v})}{|x-x'|} d^3x'$$



• The Hertz dipole

$$\hat{J} = \hat{J}_0 e^{i\omega t} = \frac{I}{a} e^{i\omega t} \hat{z} \qquad \qquad A_\rho = 0, A_\phi = 0, A_z = \frac{\mu}{c} \frac{I\Delta l}{r} e^{-i(\omega t - \frac{2\pi}{\lambda}r)}$$

in cylindrical coord.

$$H_{\rho} = 0, H_{\phi} = -i\frac{I\Delta l}{2\lambda}\frac{\sin\theta}{r}\left[1 - \frac{1}{i\frac{2\pi}{\lambda}r}\right]e^{-i(\omega t - \frac{2\pi}{\lambda}r)}, H_{z} = 0$$

in spherical coord.

$$E_r = i \frac{I\Delta l}{2\lambda} \frac{2\cos\theta}{r} \left[\frac{1}{i\frac{2\pi}{\lambda}r} - \frac{1}{(i\frac{2\pi}{\lambda}r)^2} \right] e^{-i(\omega t - \frac{2\pi}{\lambda}r)}, E_\theta = -i \frac{I\Delta l}{2\lambda} \frac{2\cos\theta}{r} \left[1 - \frac{1}{i\frac{2\pi}{\lambda}r} + \frac{1}{(i\frac{2\pi}{\lambda}r)^2} \right] e^{-i(\omega t - \frac{2\pi}{\lambda}r)}, E_\phi = 0$$

far field

$$\begin{split} |\langle S \rangle| &= \frac{c}{4\pi} |Re(E \times H^*)| = \frac{c}{4\pi} (\frac{I\Delta l}{2\lambda})^2 \frac{\sin^2\theta}{r^2} \\ P &= \int |\langle S \rangle| r^2 \sin\theta d\theta d\phi = \frac{2c}{3} (\frac{I\Delta l}{2\lambda})^2 \end{split}$$

- The reciprocity theorem
 - Two (no ohmic loss) antennas A1,A2 and non-directional medium



 $U_2I_1 = U_1I_2$

• Thus, the properties of an antenna will be the same no matter if it is transmitting or receiving

- antenna parameters
 - power pattern

$$|\langle S \rangle| = P(\theta, \phi)$$
 $|\langle S \rangle| \equiv P(\theta, \phi) = S_0 sin^2 \theta$

c.f. Hertz dipole

$$P_n(\theta, \phi) = \frac{1}{P_{max}} P(\theta, \phi) \qquad P_n(\theta, \phi) = \sin^2 \theta$$
$$G(\theta, \phi) = \frac{4\pi P(\theta, \phi)}{\int \int P(\theta, \phi) d\Omega} \qquad G(\theta, \phi) = \frac{3}{2} \sin^2 \theta$$

 antenna parameters (in general, draw figure) beam solid angle

$$\Omega_A = \int P_n(\theta, \phi) d\Omega = \int \int P_n(\theta, \phi) \sin\theta d\theta d\phi$$

main beam solid angle

$$\Omega_{MB} = \int_{main\ lobe} P_n(\theta, \phi) d\Omega$$

main beam efficiency

$$\eta_B = \frac{\Omega_{MB}}{\Omega_A}$$

directivity

$$D = G_{max} = \frac{4\pi}{\Omega_A}$$

HPBW (half power beam width), FWHP(full width half power), FWHM (full width half maximum) BWFN (beam width between first nulls) EWMB (equivalent width of the main beam)

 antenna parameters geometric aperture

effective aperture

$$A_e = \frac{P_e}{|\langle S \rangle|}$$

 A_g

aperture efficiency

$$\eta_A = \frac{A_e}{A_g}$$

effective aperture and directivity

$$D = G_{max} = \frac{4\pi A_e}{\lambda^2}$$

Antenna Temperature

$$W = \frac{1}{2}A_e \int B_{\nu}(\theta, \phi) P_n(\theta, \phi) d\Omega \equiv kT_A$$
$$T_A(\theta, \phi) \equiv \frac{\int T_b(\theta', \phi') P_n(\theta' - \theta, \phi' - \phi) d\Omega'}{\int P(\theta', \phi') d\Omega'}$$

Solutions from the Wave Equation

$$A(x,t) = \frac{\mu}{c} \int \frac{J(x',t - \frac{|x - x'|}{v})}{|x - x'|} d^3x'$$

if
$$J(x,t) = J(x)e^{-i\omega t}$$

 $A(x,t) = A(x)e^{-i\omega t}$ $A(x) = \frac{\mu}{c} \int J(x) \frac{e^{ik|x-x'|}}{|x-x'|} d^3x'$

Fraunhofer approximation: far from the antenna and and small antenna size

$$|x - x'| \approx r - n \cdot x', \ r = |x|, r \gg n \cdot x'$$
$$A(x) = \frac{\mu}{c} \frac{e^{ikr}}{r} \int J(x')^{-ikn \cdot x'} d^3x'$$

• A 2D case with approximation

last time in spherical coord. far field

$$E_{r} = i \frac{I\Delta l}{2\lambda} \frac{2\cos\theta}{r} \left[\frac{1}{i\frac{2\pi}{\lambda}r} - \frac{1}{(i\frac{2\pi}{\lambda}r)^{2}} \right] e^{-i(\omega t - \frac{2\pi}{\lambda}r)}, E_{\theta} = -i \frac{I\Delta l}{2\lambda} \frac{2\cos\theta}{r} \left[1 \frac{1}{i\frac{2\pi}{\lambda}r} + \frac{1}{(i\frac{2\pi}{\lambda}r)^{2}} \right] e^{-i(\omega t - \frac{2\pi}{\lambda}r)}, E_{\phi} = 0$$

$$k(nx') = \frac{2\pi}{\lambda} (n \cdot x') \ge 1$$

$$dE = -\frac{i}{2} \lambda J_{0}g(x') \frac{F_{e}(n)}{|x - x'|} e^{-i(\omega t - k|x - x'|)} \frac{dx'}{d\lambda} \frac{dy'}{d\lambda}$$

$$E = -\frac{i}{2} \lambda J_{0} \frac{F_{e}(n)}{r} e^{-i(\omega t - kr)} \int \int g(x') e^{-ikn \cdot x'} \frac{dx'}{d\lambda} \frac{dy}{di}$$

$$= -i\lambda J_{0}\pi \frac{F_{e}(n)}{r} f(n) e^{-i(\omega t - kr)}$$

$$f(n) = \frac{1}{2\pi} \int \int g(x') e^{-ikn \cdot x'} \frac{dx'}{\lambda} \frac{dy'}{\lambda}$$

$$P_{n}(n) = \frac{P(n)}{P_{max}} = \frac{|E \cdot E^{*}|}{|E \cdot E^{*}|_{max}} = \frac{|f(n)|^{2}}{|f_{max}|^{2}}$$

• A special case (squire aperture)

$$g(x,y)=1 \text{ for } |x| \leq \frac{L_x}{2}, |y| \leq \frac{L_y}{2}$$
$$g(x,y)=0 \text{ elsewhere}$$

$$f(l,m) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{\infty}^{\infty} g(x',y') e^{-i\frac{2\pi}{\lambda}(lx'+my')} \frac{dx'}{\lambda} \frac{dy'}{\lambda}$$
$$= \frac{\sin(\pi lL_x/\lambda)}{\pi lL_x/\lambda} \frac{\sin(\pi mL_y/\lambda)}{\pi mL_y/\lambda}$$

$$P_n(l,m) = \left[\frac{\sin(\pi l L_x/\lambda)}{\pi l L_x/\lambda} \frac{\sin(\pi m L_y/\lambda)}{\pi m L_y/\lambda}\right]^2$$

BWFNx
$$=\frac{\lambda}{L_x}$$

HPBWx $=0.88\frac{\lambda}{L_x}$



• Another special case (circular aperture)

$$x = \lambda \rho \cos \phi$$

$$y = \lambda \rho \sin \phi$$

$$u = \sin \theta$$

$$f(u) = \frac{1}{2\pi} \int_0^{2\pi} \int_0^{\infty} g(\rho) e^{-2\pi u \rho \cos \phi} \rho d\rho d\phi$$

$$= \int_0^{\infty} g(\rho) J_0(2\pi u \rho) \rho d\rho$$

$$Pn(u) = \left[\frac{\int_0^{\infty} g(\rho) J_0(2\pi u \rho) \rho d\rho}{\int_0^{\infty} g(\rho) \rho d\rho}\right]^2$$

if uniform illumination

 $g(\rho) = 1 for \rho \le D/2\lambda$ $g(\rho) = 0 elsewhere$

$$Pn(u) = \left[\frac{2J_1(\pi u D/\lambda)}{\pi u D/\lambda}\right]^2$$
$$= \Lambda_1^2(\pi u D/\lambda)$$

BWFN =2.44
$$\frac{\lambda}{D}$$
 rad
HPBW =1.02 $\frac{\lambda}{D}$ rad

- surface accuracy phase errors
 - linear phase error axis tilt
 - quadratic phase error defocusing
 - random phase error reduction of directive gain

$$G = \frac{4\pi}{\lambda^2} \frac{|\int g_0(x) e^{-i[kn \cdot x - \delta(x)]} d^2 x|^2}{\int g_0^2(x) d^2 x}$$

• in the simplified case when phase error is small and the average phase error can be made to 0

$$\frac{G}{G_0} = 1 - \bar{\delta^2} \qquad \qquad \bar{\delta^2} = \frac{\int g_0(x)\delta^2(x)d^2x}{\int g_0(x)d^2x}$$

Practical Construction Consideration

- mostly mechanical properties
 - deformation gravity, thermal, wind, etc
 - surface accuracy measurements
 - fiducial mark surveying
 - holography (artificial/celestial source)
 - pointing/focusing mount
 - quatorial
 - simple design
 - only useful for small aperature
 - alt-az
 - complicated control
 - constant gravity angle
 - difficult zenith tracking
 - change of position angle
 - optical or radio measurements
 - feedleg blockage standing waves

Practical Construction Consideration

- overcome mechanical difficulties
 - backup structure and surface (wires/solid panel)
 - reduce gravity deformation
 - motion restricted Arecibo
 - rigid design
 - passive compensation (homologous deformation)
 - parabolic shape maintained
 - apex, focal point, axis changed
 - applicable for all radio wavelengths (OVRO...)
 - active compensation (GBT)
 - reduce thermal/wind deformation
 - shielding dorm (CSO,JCMT,ARO 10m/12m)
 - improved material (SMA, APEX)
 - feedleg blockage
 - off-axis design (GBT)

Back Ends

- continuum
- spectrometers
 - analog
 - Michelson interferometer
 - Fourier transform spectrometer (FTS)
 - Fabry-Perot spectrometer
 - Multichannel Filter Spectrometer (Filter Bank)
 - AOS
 - CTS
 - digital
 - digitization and FFT
 - (auto)correlator
 - advantages of (ditigal) correlator over analog systems
 - flexibility
 - stability





Homework 3

- 3.1 : Express the Faraday Rotation formula in more sensible units and check the typical values for this effect in a reference (Please cite).
- 3.2 : For the following telescopes at their operation frequencies, at which distances a plane wave approximation can be applied for sources:
 - I.Arecibo
 - 2.VLA
 - 3. SMA